

Internal quantum reference frames and external frame independence

Philipp Höhn

Okinawa Institute of Science and Technology



Mathematical Physics Seminar, U Regensburg
29 January, 2026

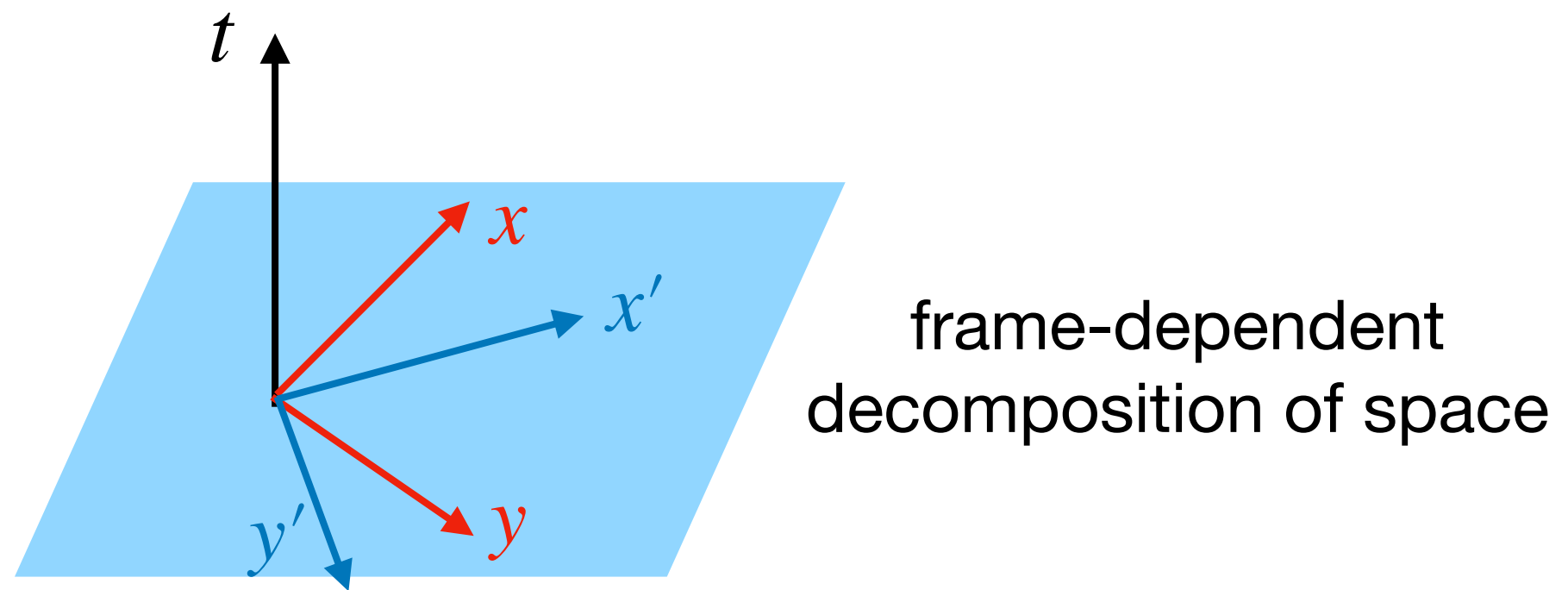
roughly based on:

PH, Mele, Kotecha 2308.09131
de la Hamette, Galley, PH, Loveridge, Müller 2110.13824
Araújo-Regado, PH, Sartini 2506.23459

Relativities

Galilean

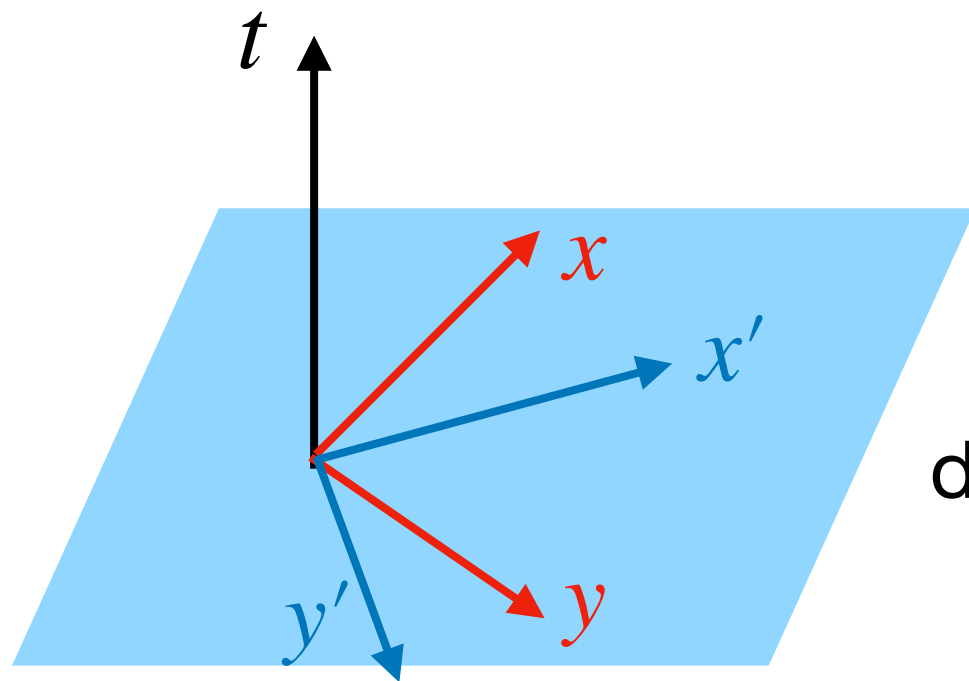
“All the laws of mechanics the same
in every inertial frame”



Relativities

Galilean

“All the laws of mechanics the same in every inertial frame”



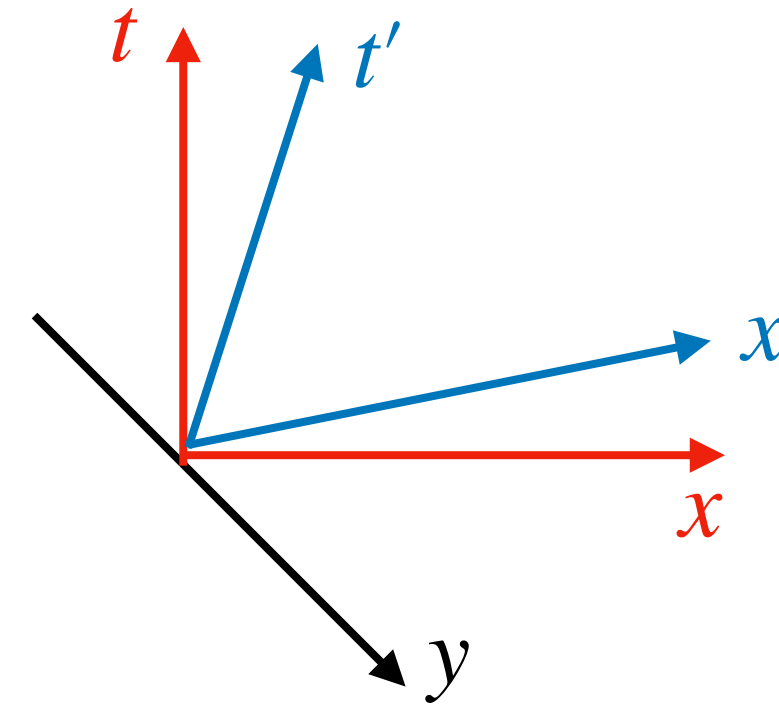
frame-dependent
decomposition of space

$$\frac{1}{c} \neq 0$$



Special

“All the laws of physics (exc. Newt. gravity) the same in every inertial frame”

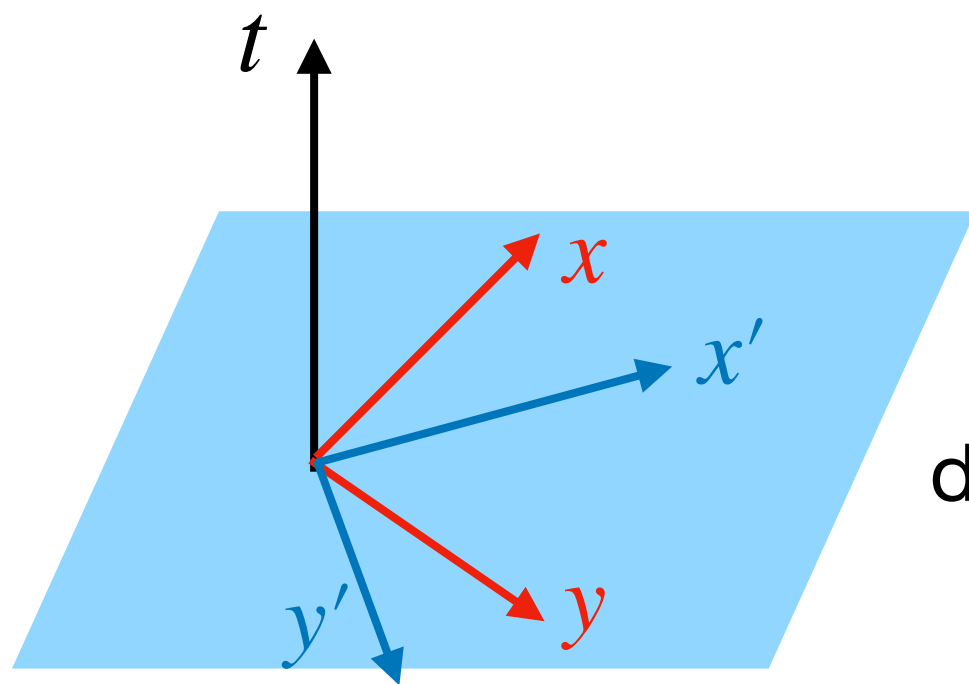


frame-dependent
decomposition of
spacetime into
space and time

Relativities

Galilean

“All the laws of mechanics the same in every inertial frame”



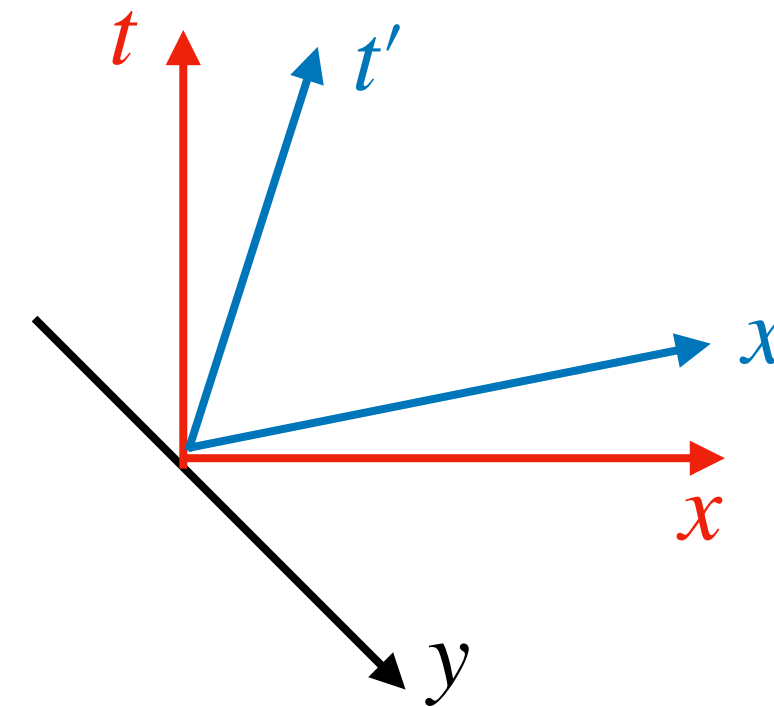
frame-dependent
decomposition of space

$$\frac{1}{c} \neq 0$$



Special

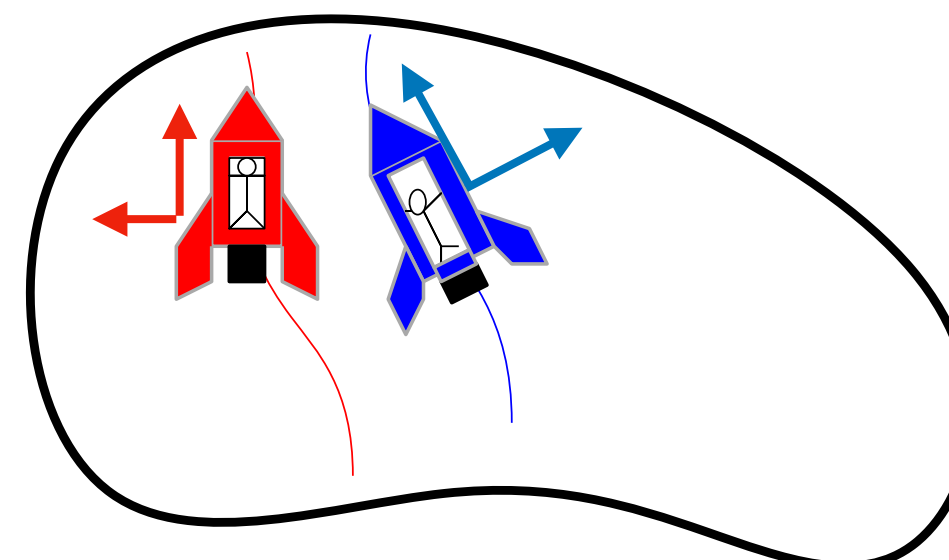
“All the laws of physics (exc. Newt. gravity) the same in every inertial frame”



frame-dependent
decomposition of
spacetime into
space and time

General

“All the laws of physics the same in every frame”



frame-dependent
local decomposition
of spacetime into
space and time

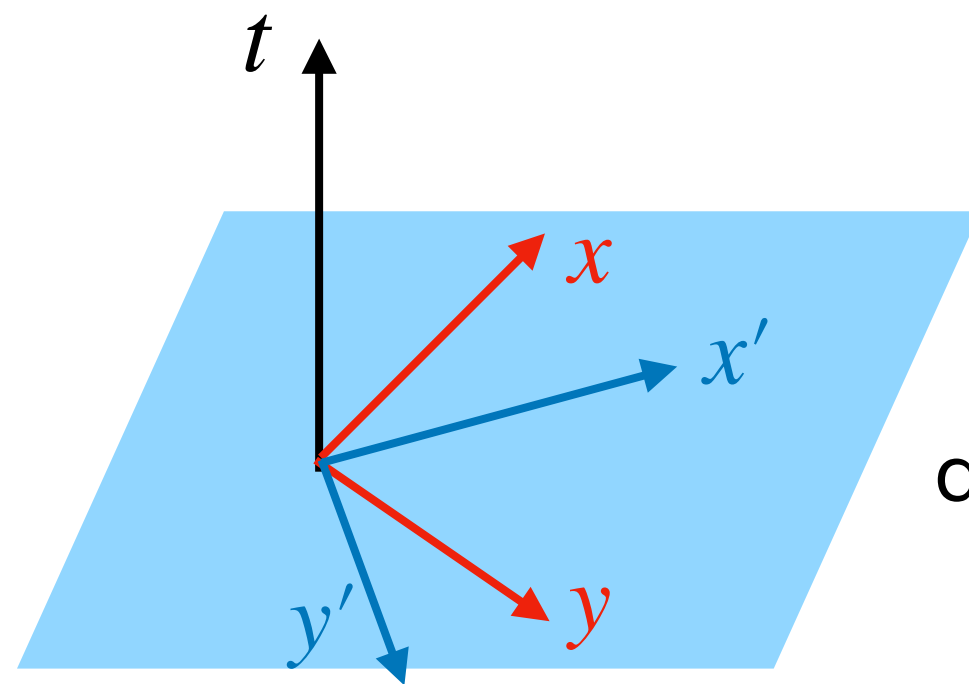
$$G \neq 0$$



Relativities

Galilean

“All the laws of mechanics the same in every inertial frame”



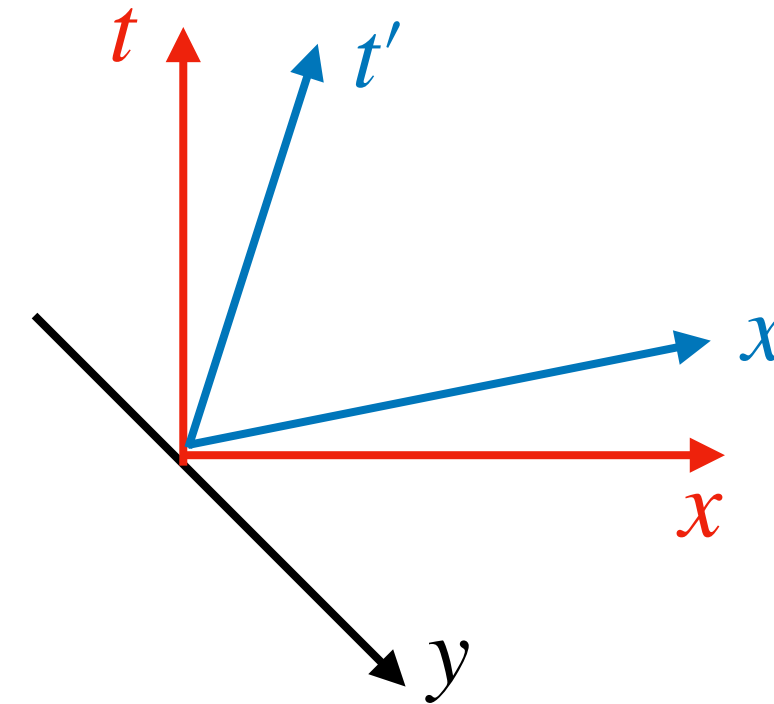
frame-dependent
decomposition of space

$$\frac{1}{c} \neq 0$$



Special

“All the laws of physics (exc. Newt. gravity) the same in every inertial frame”



frame-dependent
decomposition of
spacetime into
space and time

$$G \neq 0$$



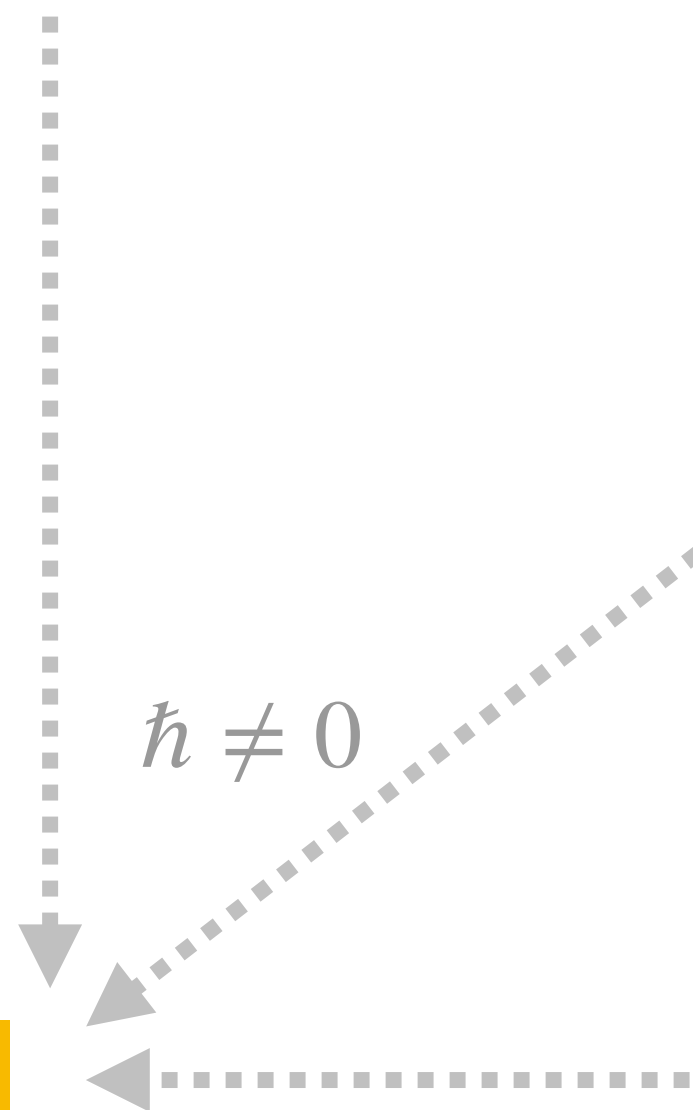
Quantum?

“All the ... laws of ... the same in every ... QRF”



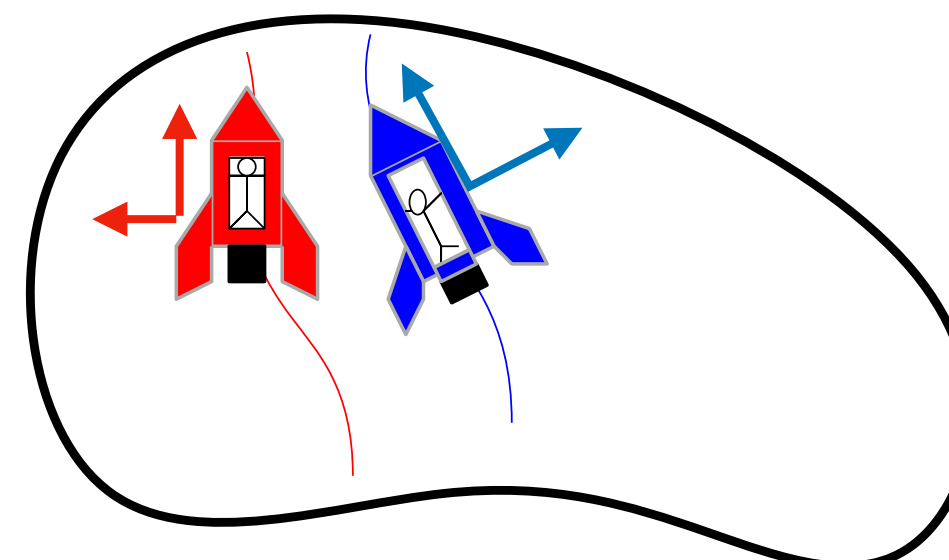
QRF-dependent
decomposition of

$$\hbar \neq 0$$



General

“All the laws of physics the same in every frame”



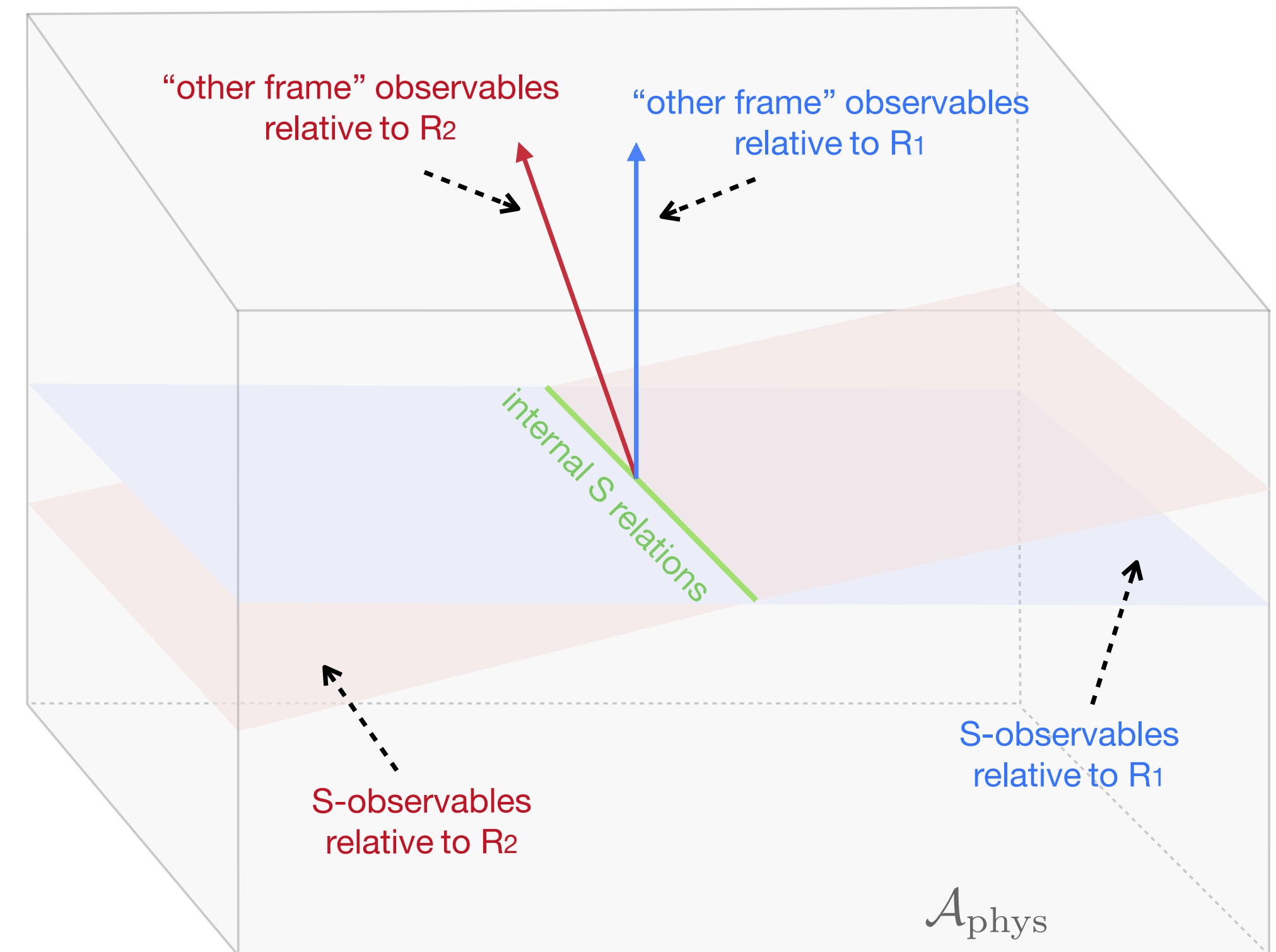
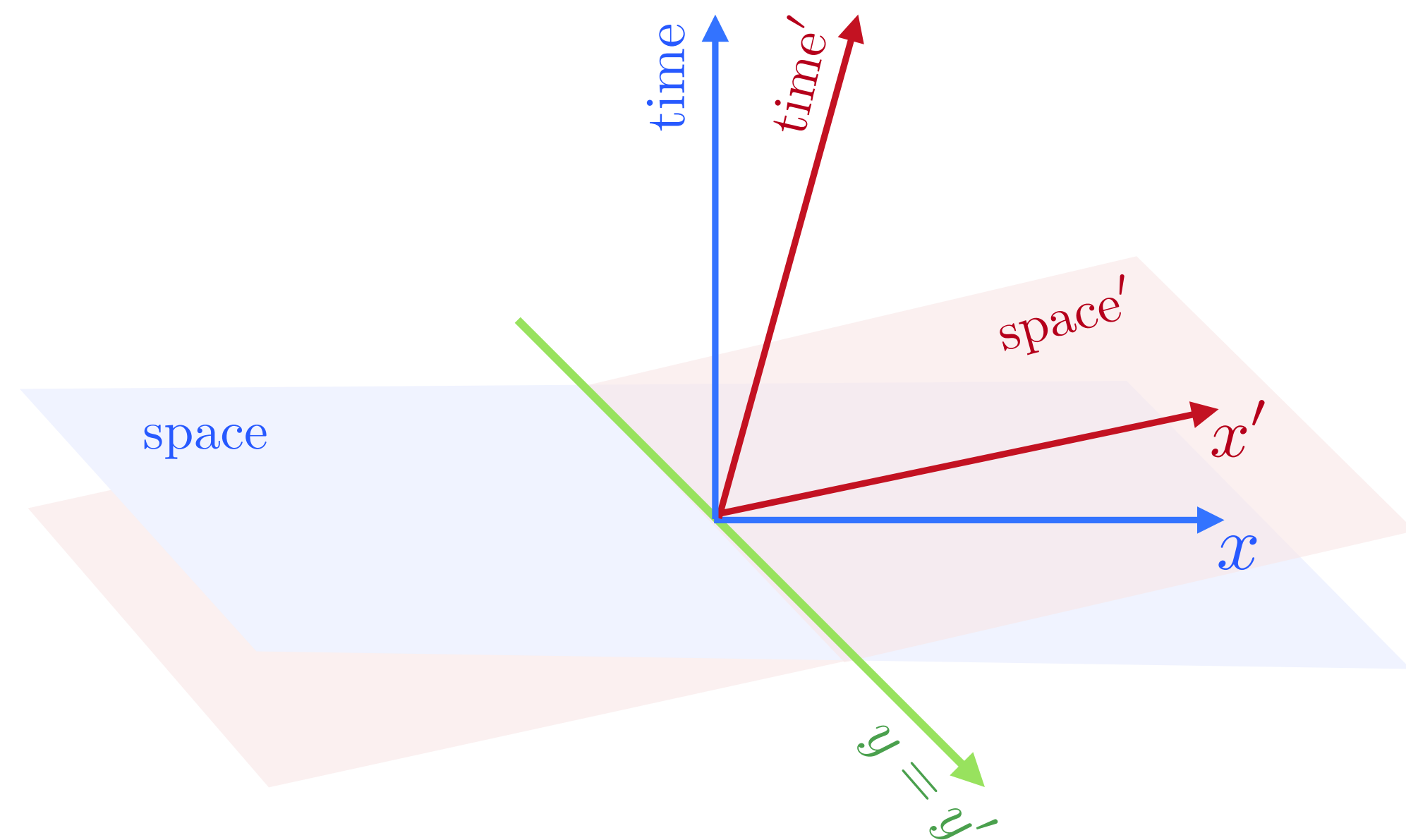
frame-dependent
local decomposition
of spacetime into
space and time

Quantum reference frames are internal subsystems and a universal tool for...

- dealing with symmetries in quantum systems
- describing a quantum system “from the inside”
(describing subsystems relative to one another)
- defining subsystem partitions in gauge systems
- completing the relational paradigm

applicable whenever an external reference isn't available/desired

Relativity of simultaneity vs. subsystem relativity



Why quantum reference frames?

- foundational interest:

- all frames in experiments physical
- no external frames (background indep.)
- quantum frame relations $\Leftrightarrow$ quantum spacetime structure?

- practical interest:

- needed for gauge-inv. description (or gauge fixing)
- act as regulator for QFT algebras and entropies
- definition of locality
- applications in QI
- ...

Menu

- Introduction to QRFs and notions of external frame independence
- subsystem relativity
- Physical and gauge QRFs in lattice gauge theories
+ entanglement entropies (maybe)

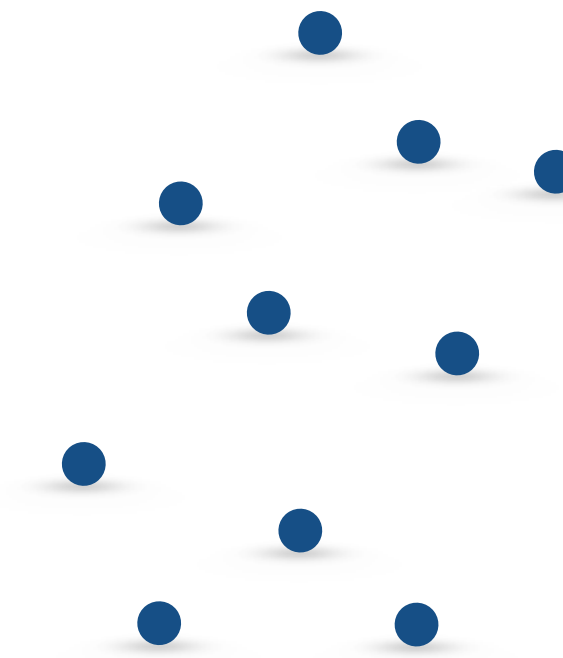
Internal QRFs and external frame independence

[de la Hamette, Galley, PH, Loveridge, Müller 2110.13824; PH, Kotecha, Mele 2308.09131;
PH, Smith, Lock 1912.00033; Vanrietvelde, PH, Giacomini, Castro-Ruiz 1809.00556,...]

RFs and symmetries

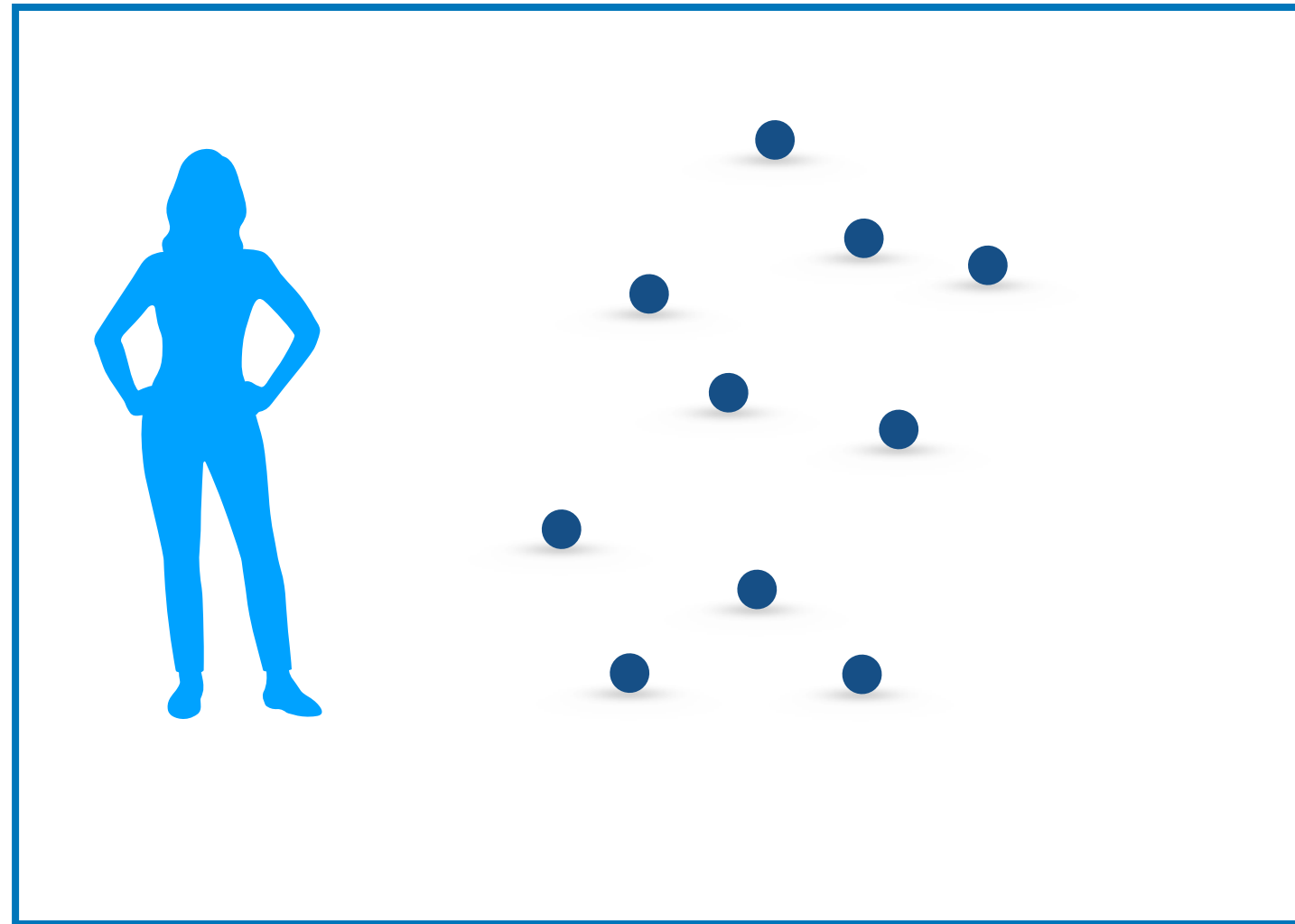
Premise:

System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation



RFs and symmetries

Premise:
System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation



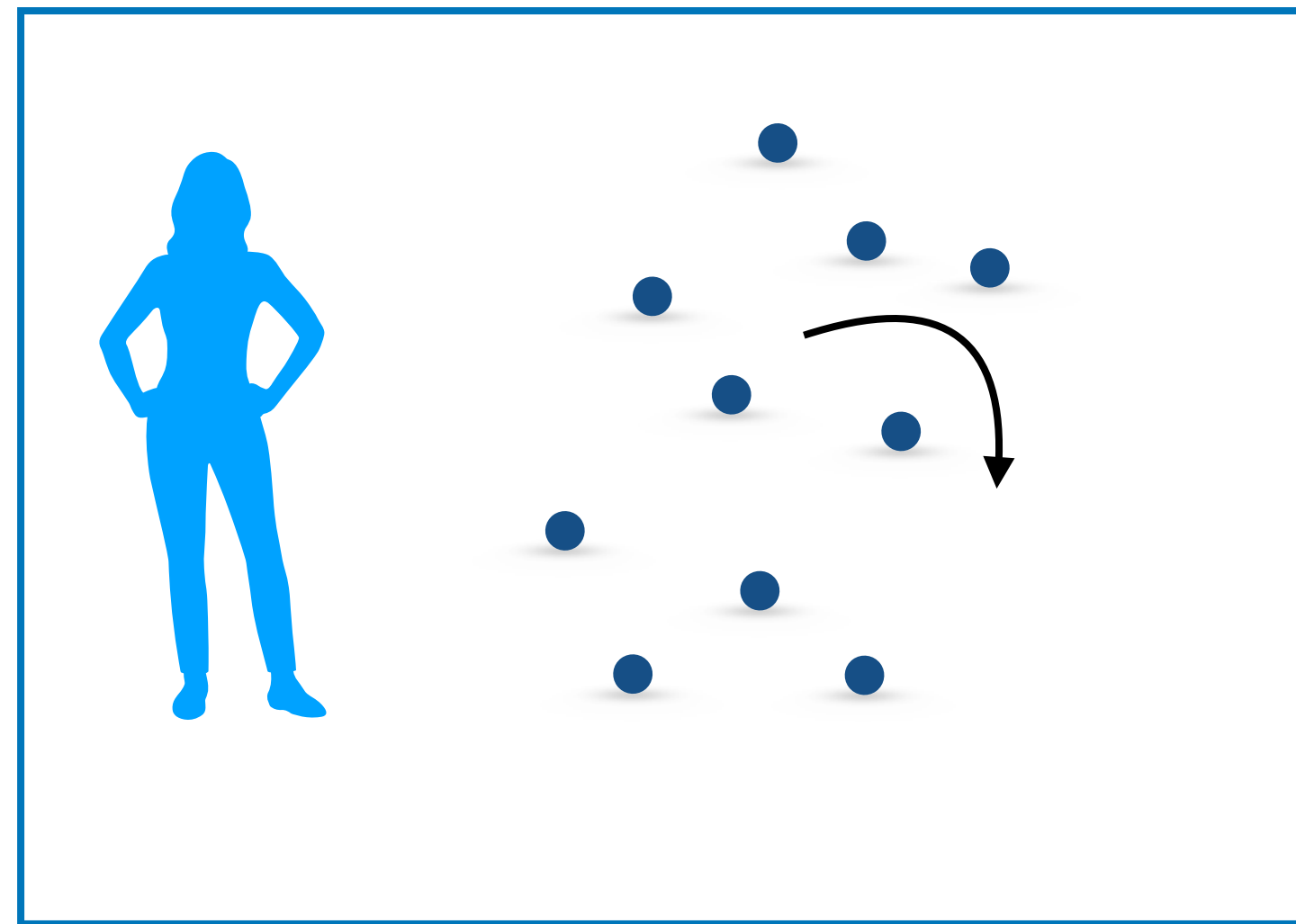
external frame

quantum information/foundations:
lab frame

gauge theory & gravity:
fictitious (or bdry)

RFs and symmetries

Premise:
System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation



external frame

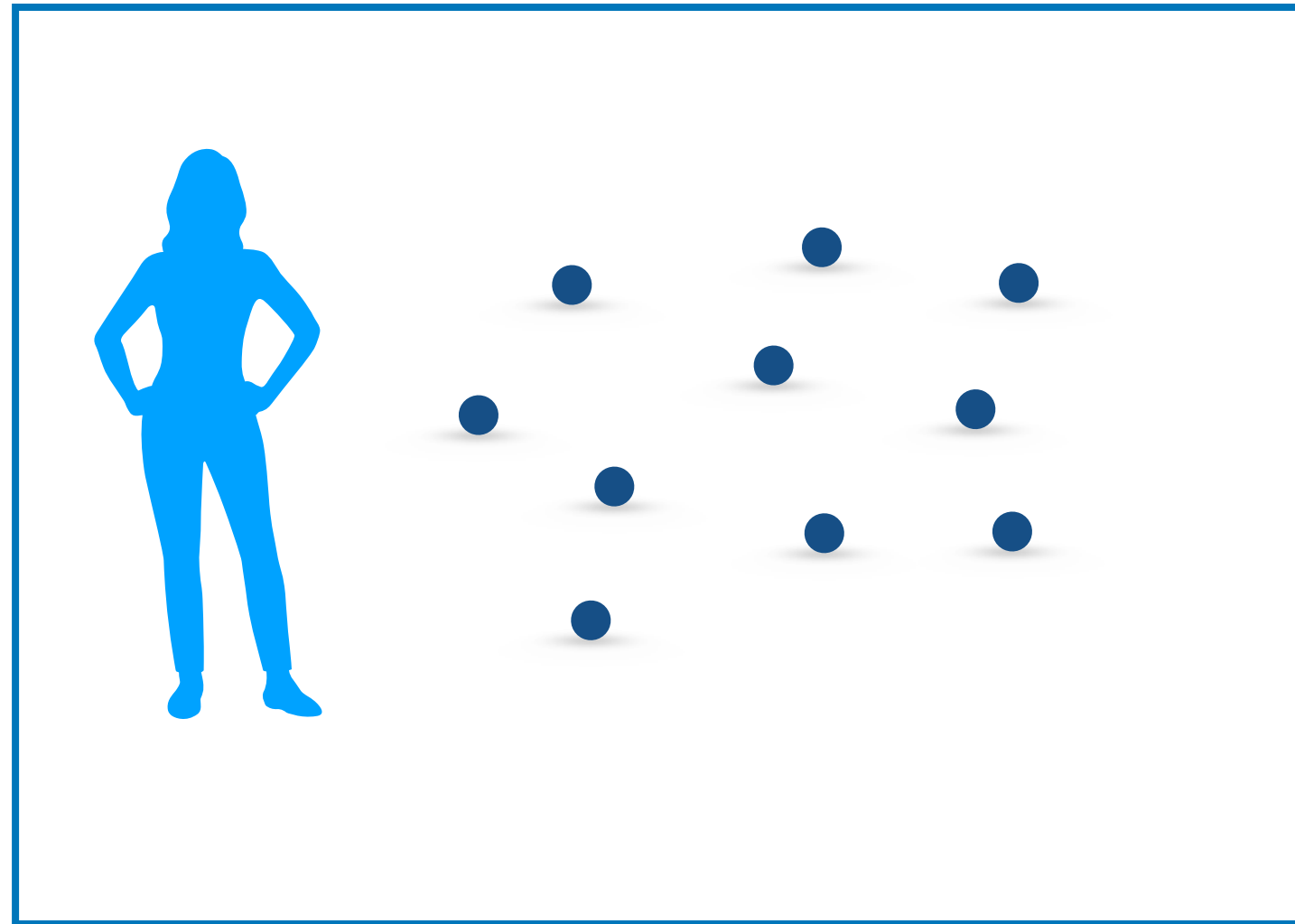
quantum information/foundations:
lab frame

gauge theory & gravity:
fictitious (or bdry)

G -transformations = ext. frame transf.

RFs and symmetries

Premise:
System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation



external frame

quantum information/foundations:
lab frame

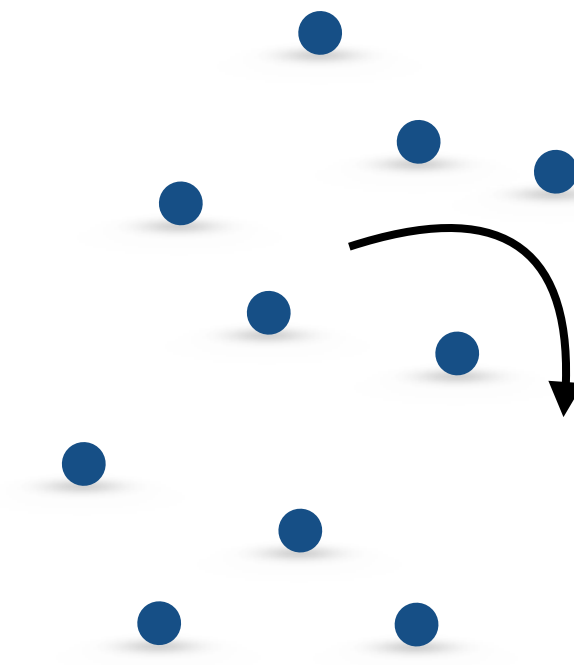
gauge theory & gravity:
fictitious (or bdry)

G -transformations = ext. frame transf.

RFs and symmetries

Premise:

System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation

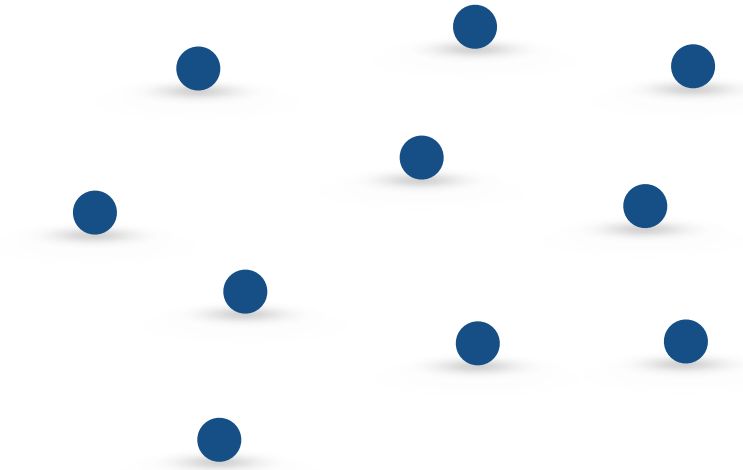


G -transformations = ext. frame transf.

RFs and symmetries

Premise:

System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation

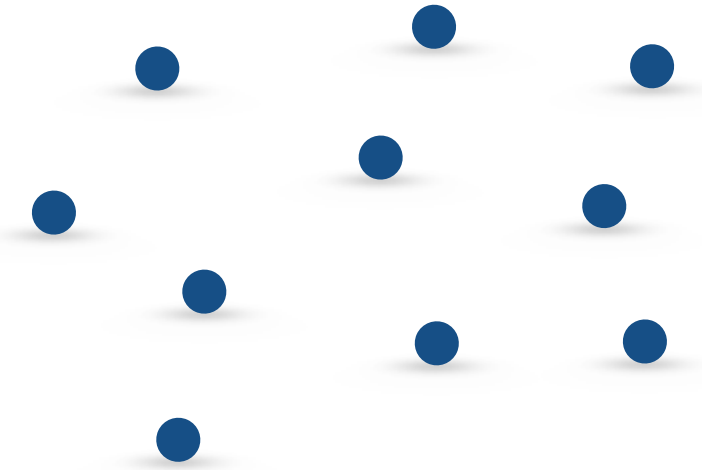


G -transformations = ext. frame transf.

RFs and symmetries

Premise:

System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation



internally indistinguishable

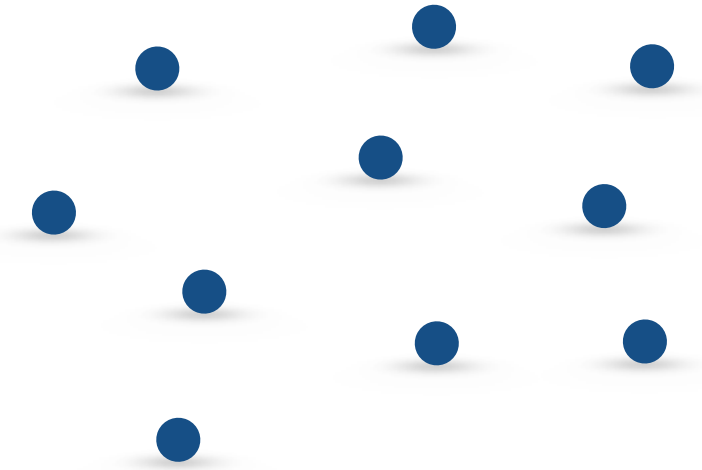
G -transformations = ext. frame transf.

$\Rightarrow$ want external frame independence, i.e.
 G -invariance

RFs and symmetries

Premise:

System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation



internally indistinguishable

quantum information/foundations:
change of ext. lab frame

G -transformations = ext. frame transf.

$\Rightarrow$ want external frame independence, i.e.
 G -invariance

gravity:
e.g. change of background coordinates (diffeo)
 $\Rightarrow$ change of fictitious background frame

RFs and symmetries

Premise:

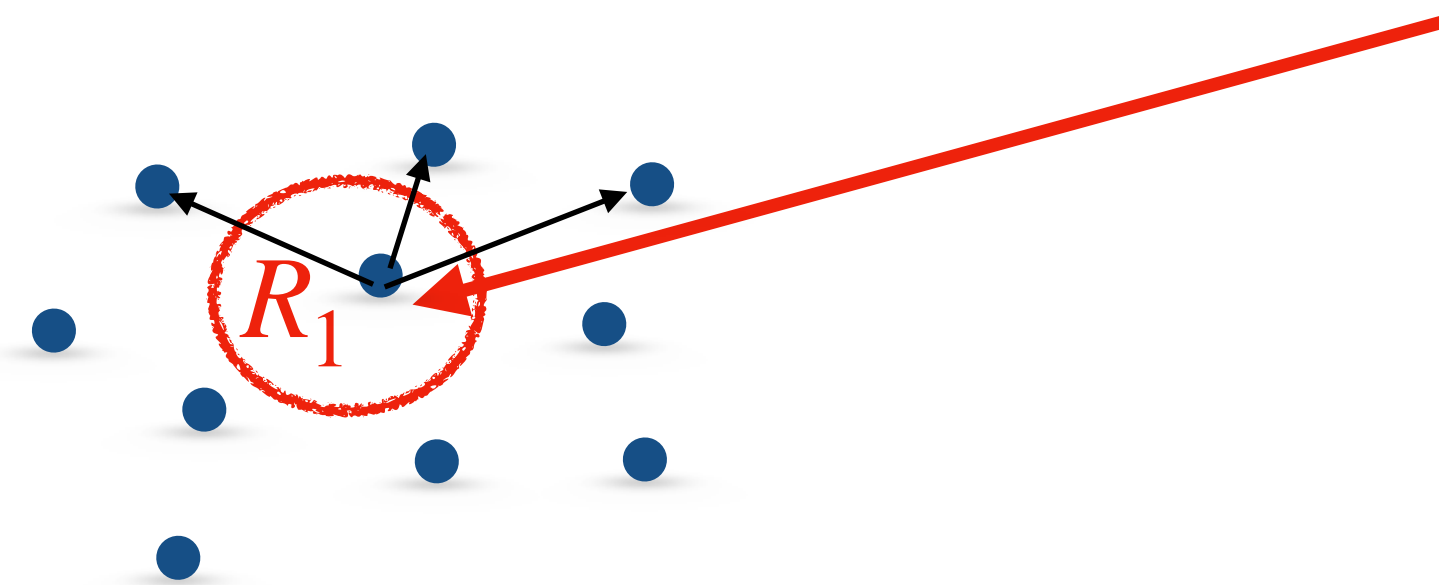
System $\tilde{S} = R_1 R_2 S$ subject to symmetry group G , s.t. states ρ and $g \cdot \rho$ are indistinguishable for all $g \in G$ when $\tilde{S}$ considered in isolation

quantum information/foundations:
change of ext. lab frame

gravity:

e.g. change of background coordinates (diffeo)
 $\Rightarrow$ change of fictitious background frame

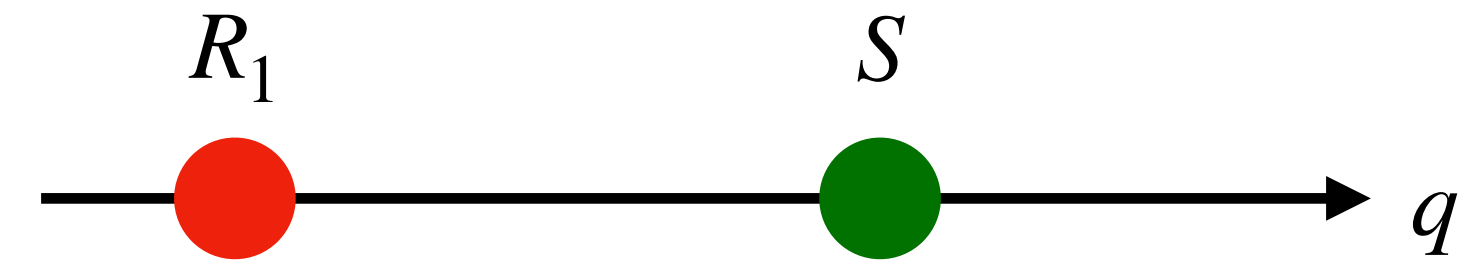
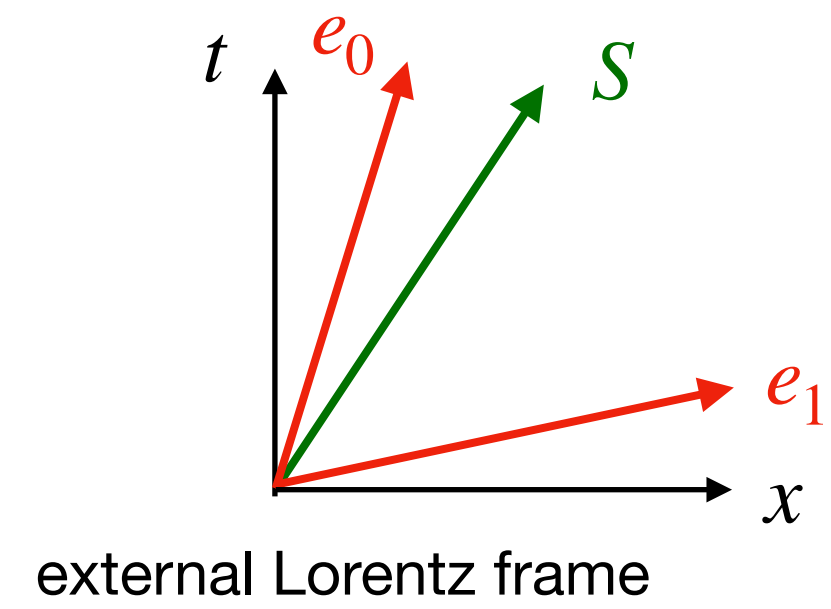
Describe $\tilde{S}$ relative to internal reference subsystem R_1



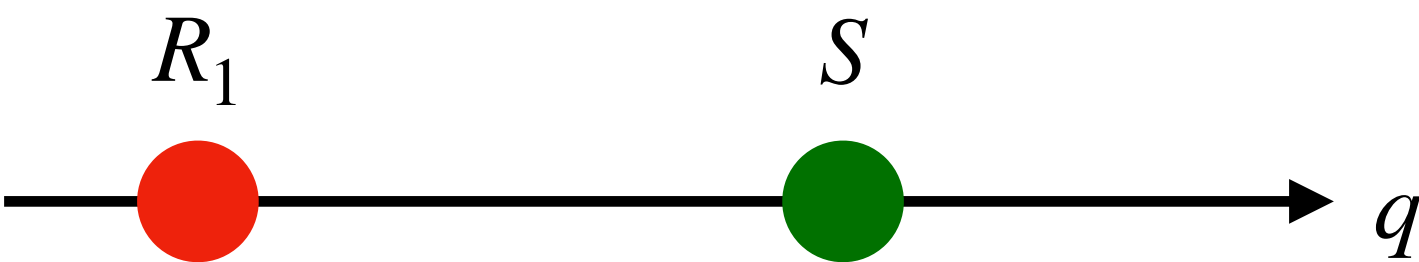
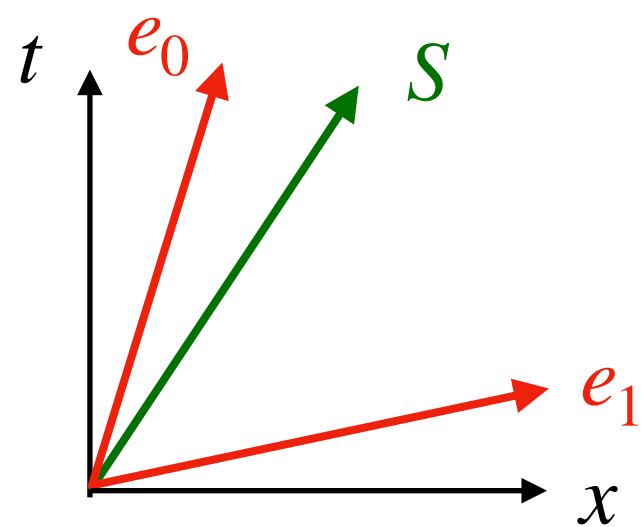
G -transformations = ext. frame transf.

$\Rightarrow$ want external frame independence, i.e.
 G -invariance

internal frames in SR vs. QT



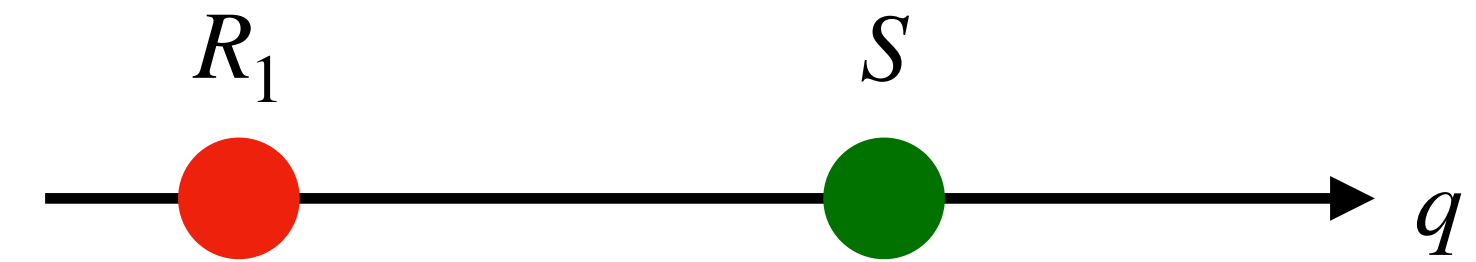
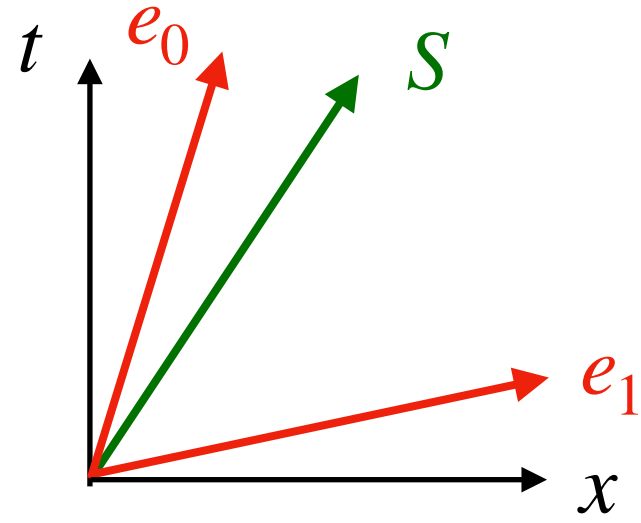
internal frames in SR vs. QT



internal frame
orientations

tetrad $e_a^\mu \in \text{SO}_+(3,1)$
(because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)

internal frames in SR vs. QT



**internal frame
orientations**

tetrad $e_a^\mu \in \text{SO}_+(3,1)$
(because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)

covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$

in practice, coh. states $|g\rangle$, $g \in G$

Example: quantum clocks

[PH, Smith, Lock '19]

clock states for non-deg. continuous spectrum clock Hamiltonian H_C (gen. of $\mathbb{R}$)

$$|t\rangle = \frac{1}{\sqrt{2\pi}} \int_{\text{Spec}(H_C)} d\epsilon e^{-it\epsilon} |\epsilon\rangle$$

clock orientations $t \in \mathbb{R}$ (group-valued)

covariant:

$$U_C(t) |t'\rangle = |t + t'\rangle$$

$$U_C(t) = e^{-itH_C}$$

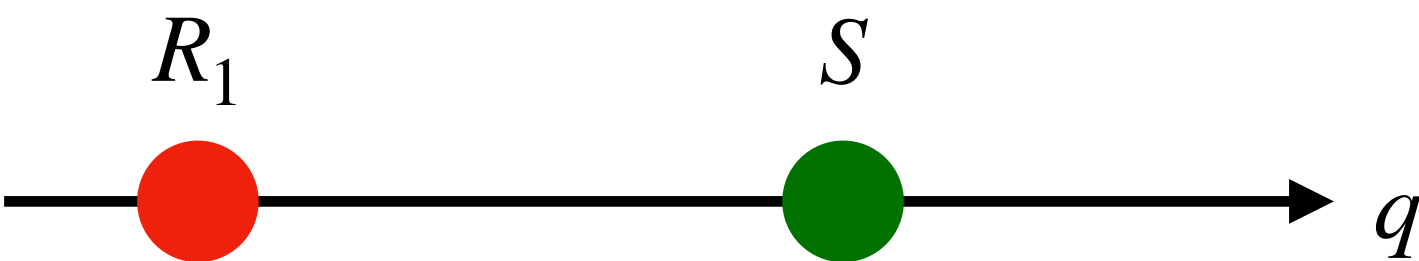
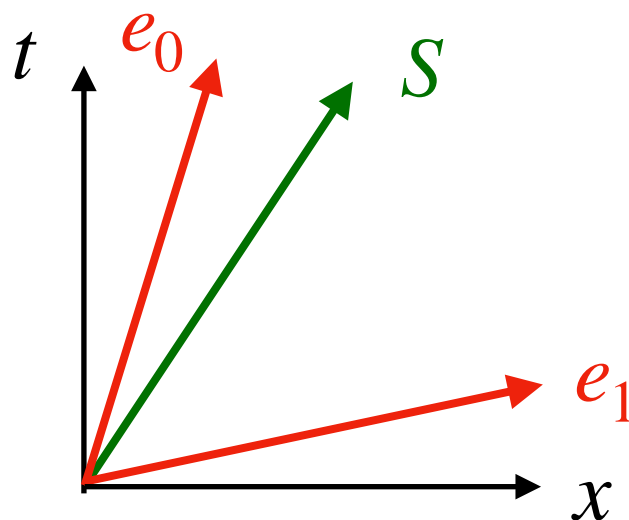
POVM elements

$$E_C(X) = \int_{X \subset \mathbb{R}} dt |t\rangle \langle t|$$

normalized $E_C(\mathbb{R}) = I$ (well-defined probabilities)

fuzzy clock states $\langle t | t' \rangle \approx \delta(t - t')$ (unless **ideal clock** with $\text{Spec}(H_C) = \mathbb{R}$)

internal frames in SR vs. QT



internal frame
orientations

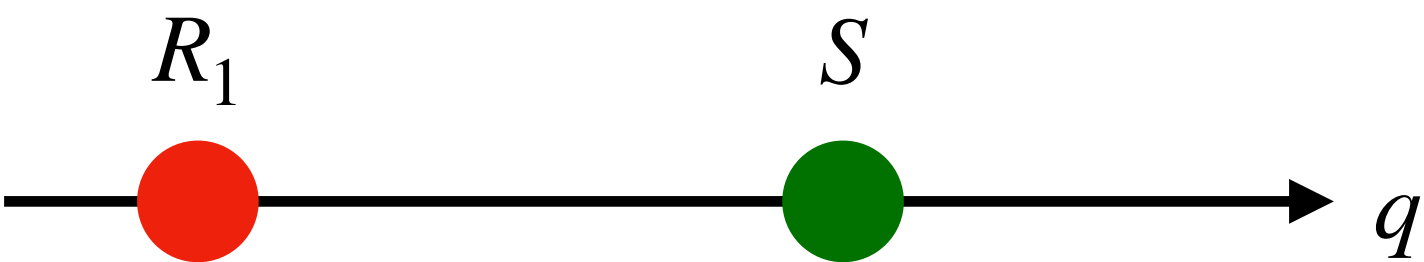
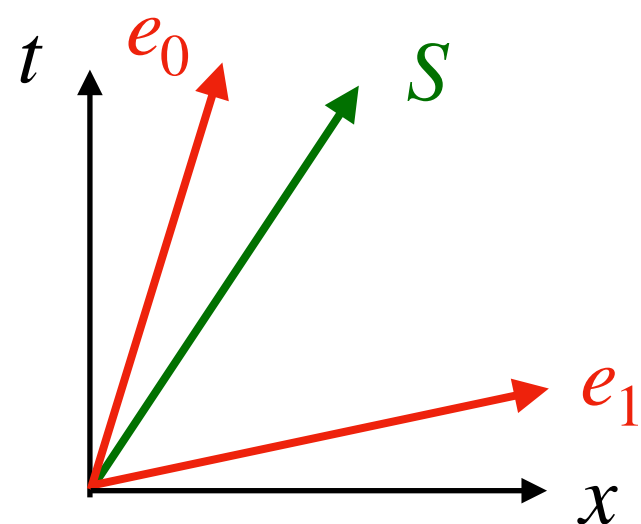
tetrad $e_a^\mu \in \text{SO}_+(3,1)$
(because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)

covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$
in practice, coh. states $|g\rangle$, $g \in G$

externally
distinguishable
configurations

$\{e_a^\mu, S^\mu\}$

internal frames in SR vs. QT



**internal frame
orientations**

tetrad $e_a^\mu \in \text{SO}_+(3,1)$
(because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)

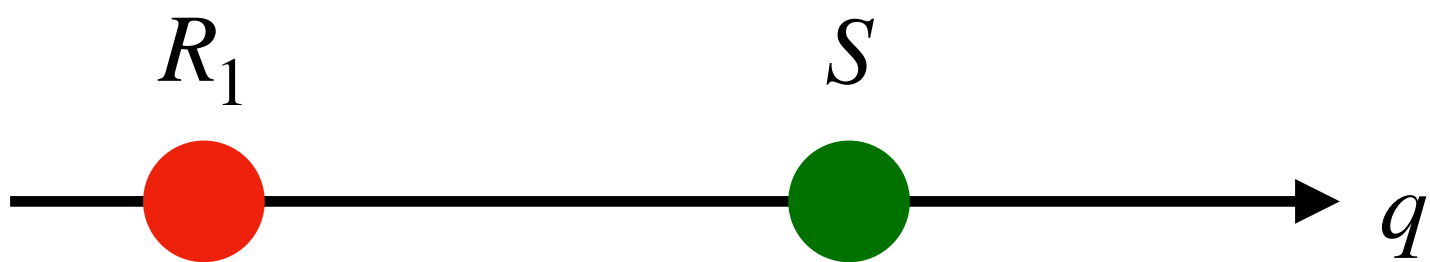
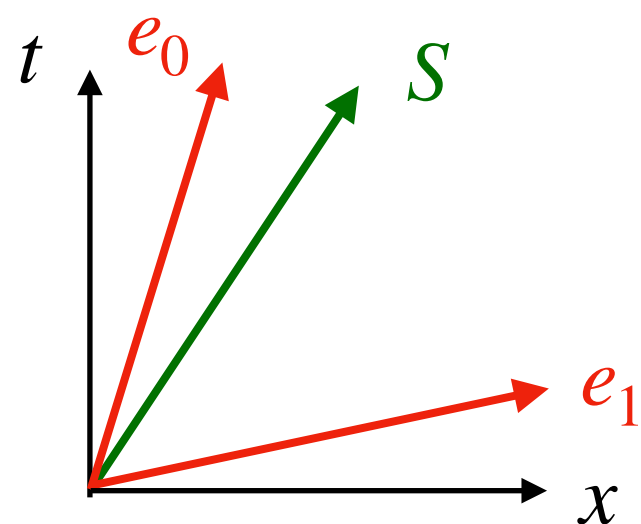
covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$
in practice, coh. states $|g\rangle, \quad g \in G$

**externally
distinguishable
configurations**

$\{e_a^\mu, S^\mu\}$

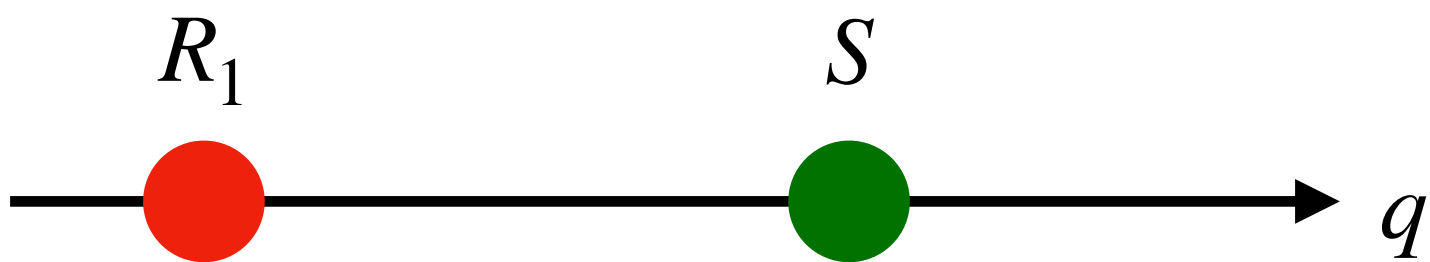
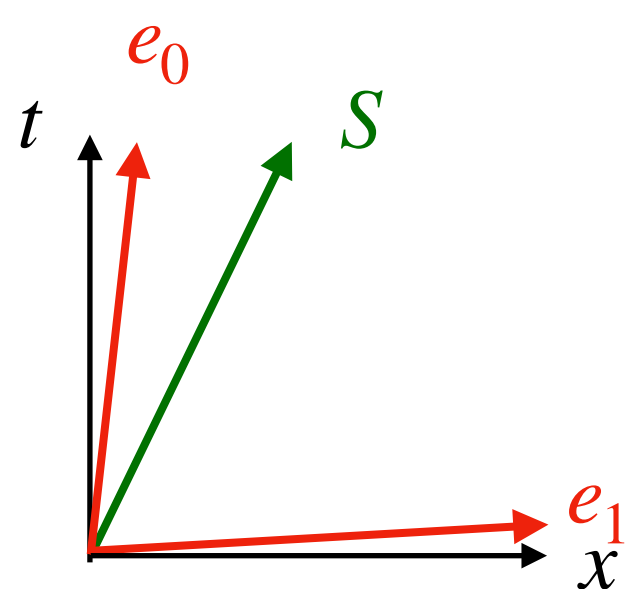
kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$
kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$

internal frames in SR vs. QT



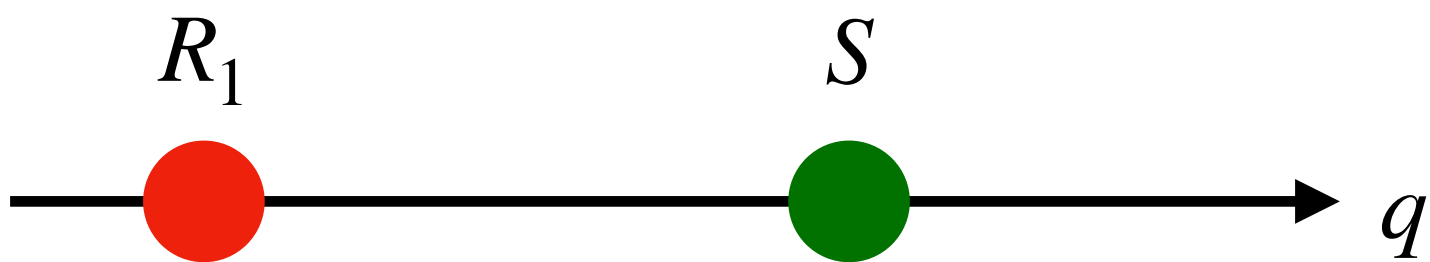
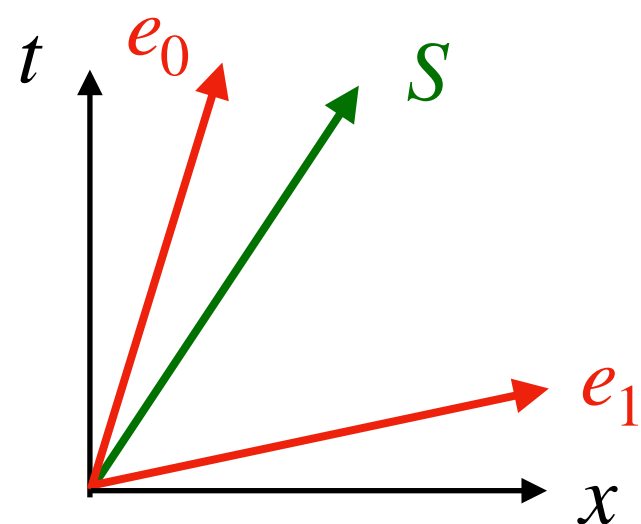
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	

internal frames in SR vs. QT



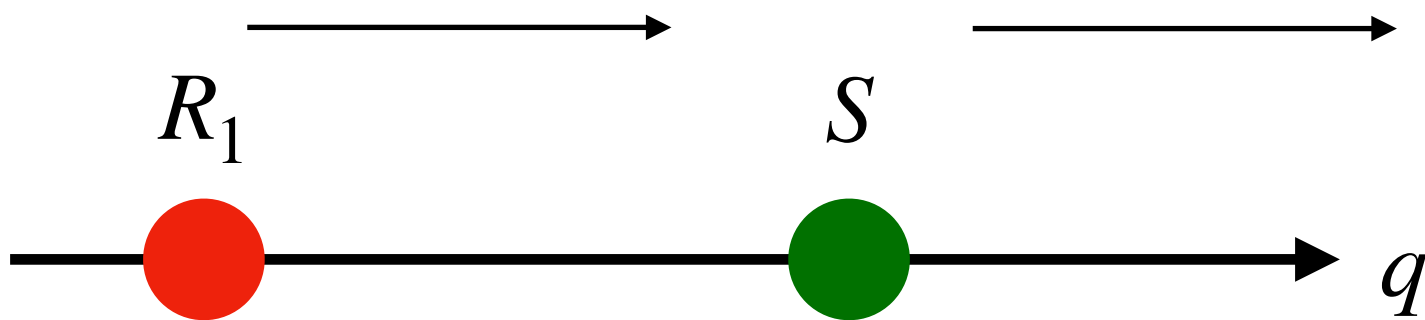
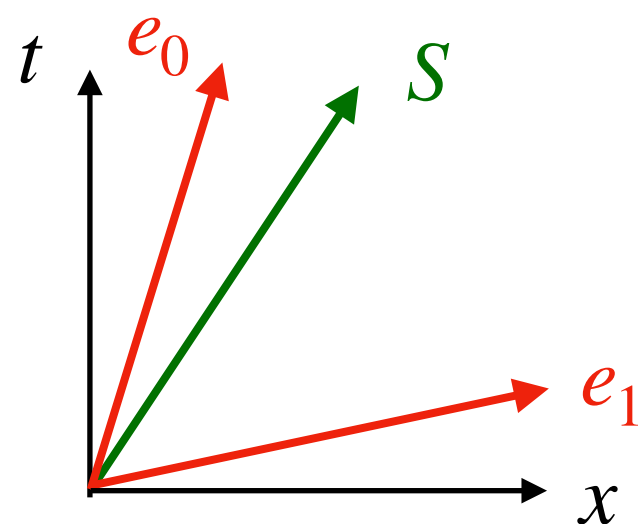
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	

internal frames in SR vs. QT



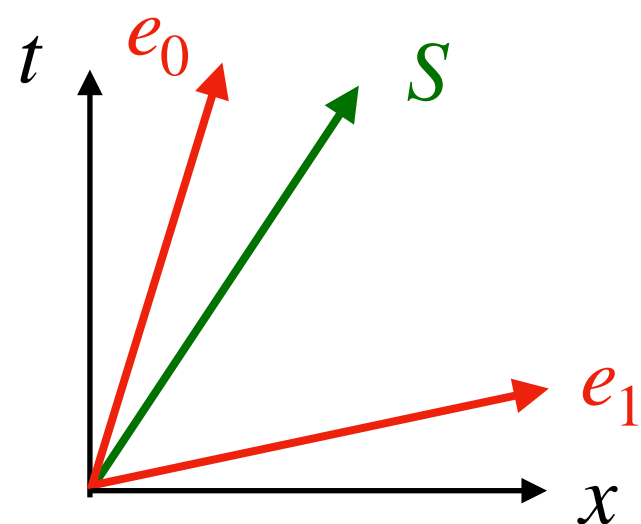
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G

internal frames in SR vs. QT



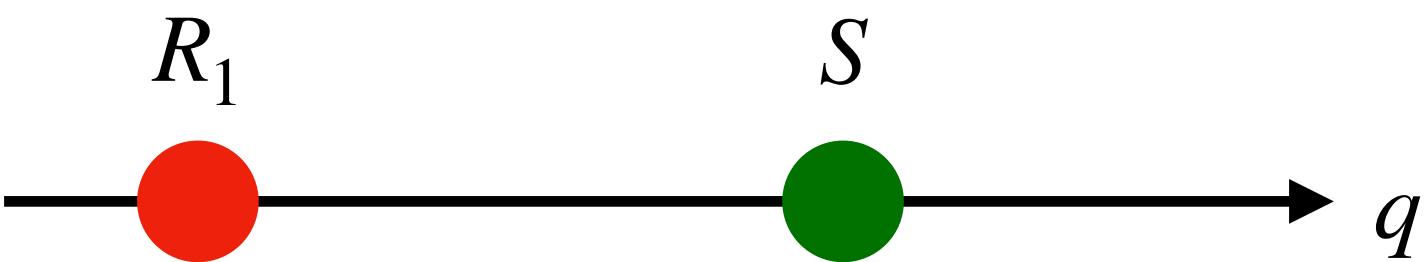
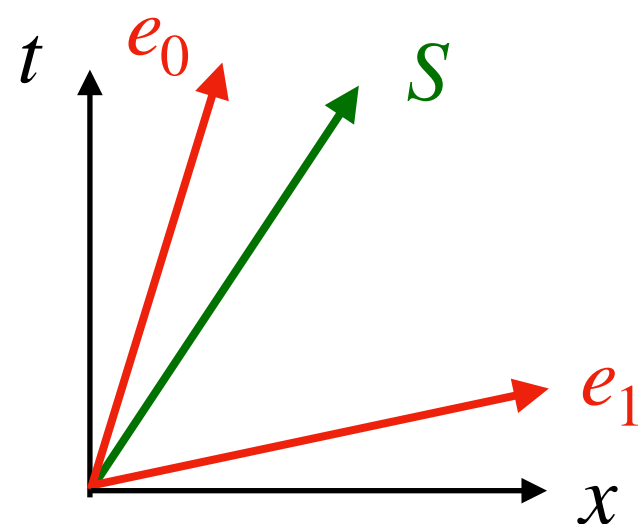
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G

internal frames in SR vs. QT



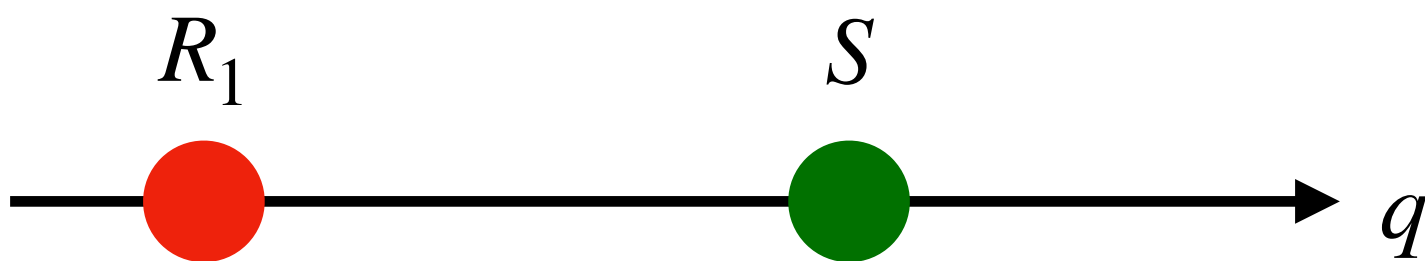
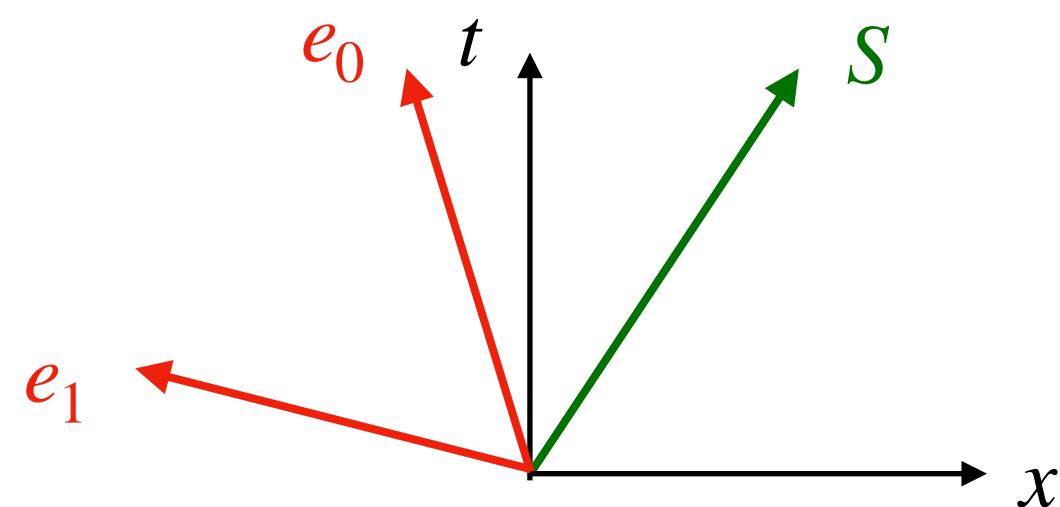
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G

internal frames in SR vs. QT



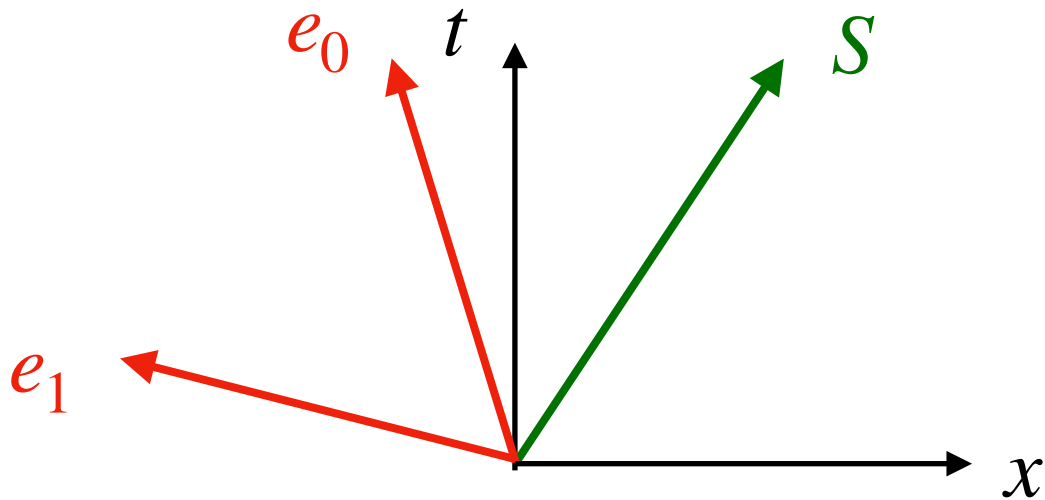
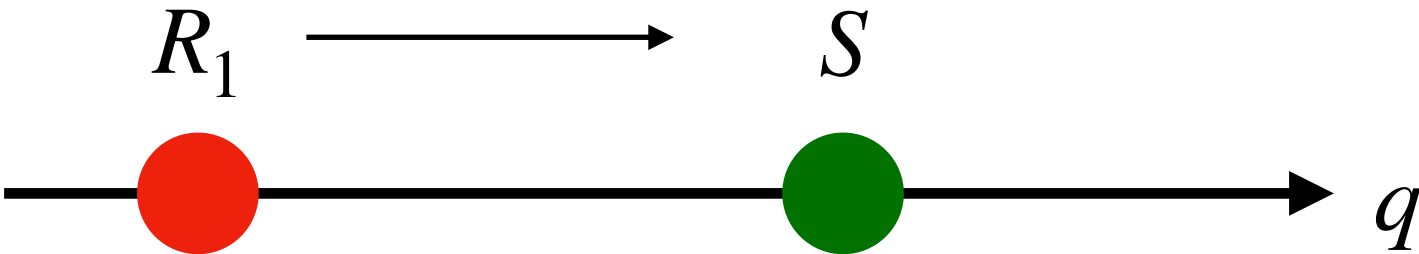
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G
reorientations (acts only on the frame)	$\Lambda_a{}^b \in \text{SO}_+(3,1)$	

internal frames in SR vs. QT

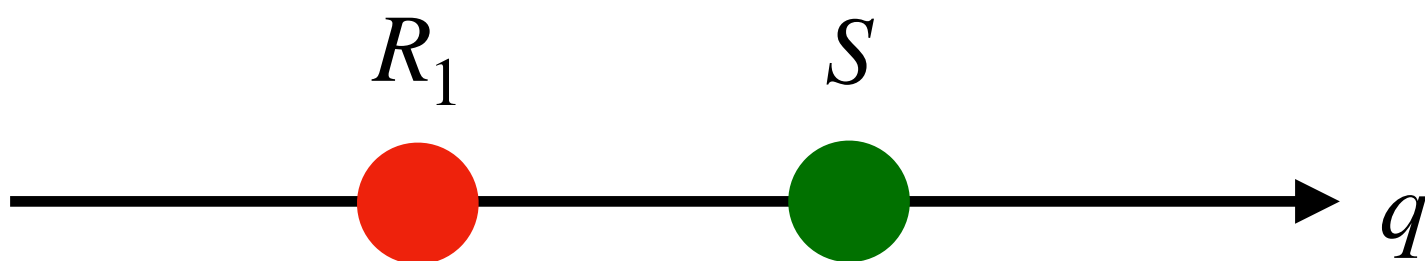
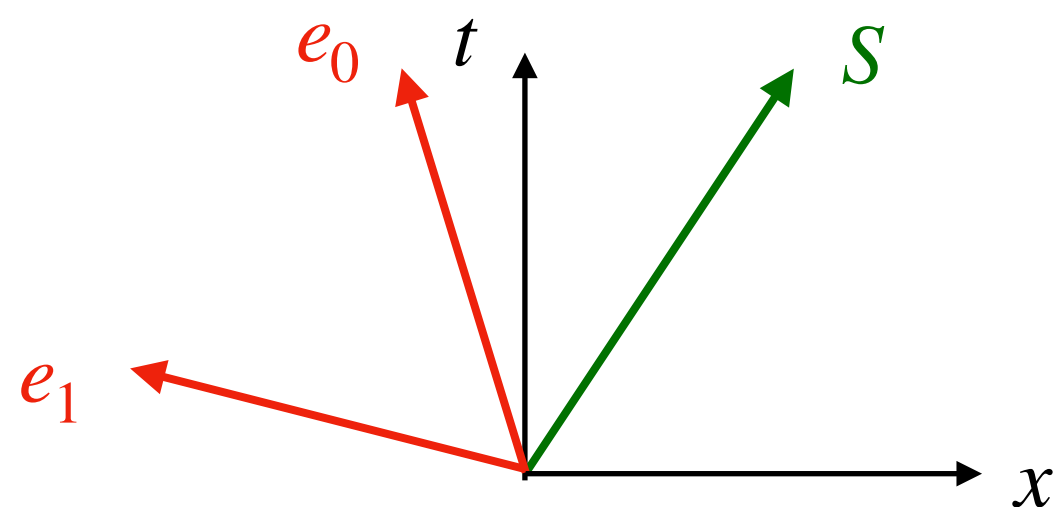


internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G
reorientations (acts only on the frame)	$\Lambda_a{}^b \in \text{SO}_+(3,1)$	

internal frames in SR vs. QT

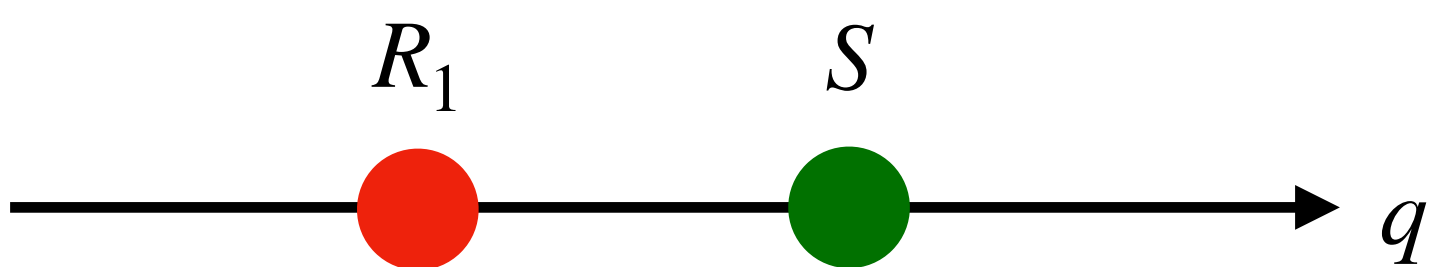
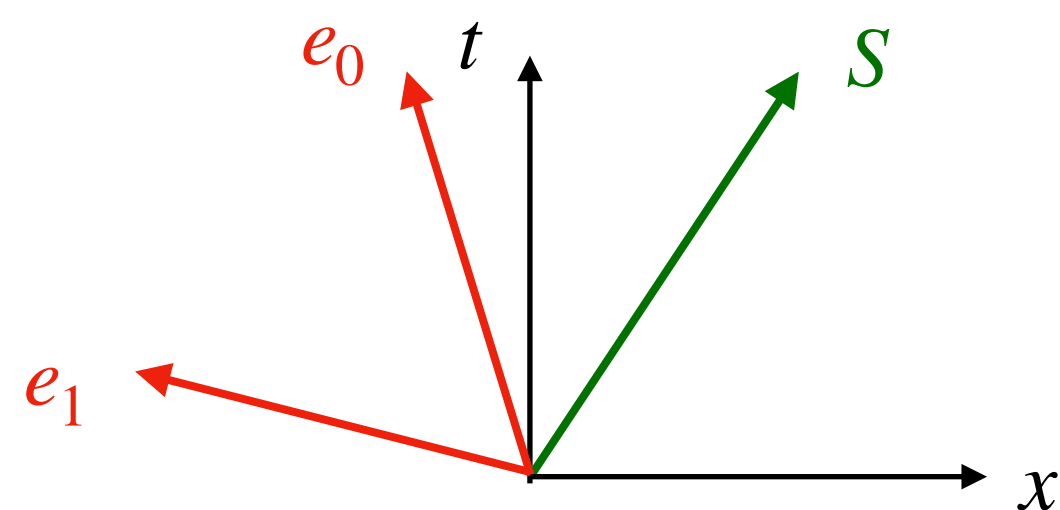
		
internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle, \quad g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G
reorientations (acts only on the frame)	$\Lambda_a{}^b \in \text{SO}_+(3,1)$	$V_R^g \otimes I_S$ s.t. $V_R^g g'\rangle = g'g^{-1}\rangle$

internal frames in SR vs. QT



internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $g\rangle, \quad g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G
reorientations (acts only on the frame)	$\Lambda_a{}^b \in \text{SO}_+(3,1)$	$V_R^g \otimes I_S$ s.t. $V_R^g g'\rangle = g'g^{-1}\rangle$

internal frames in SR vs. QT

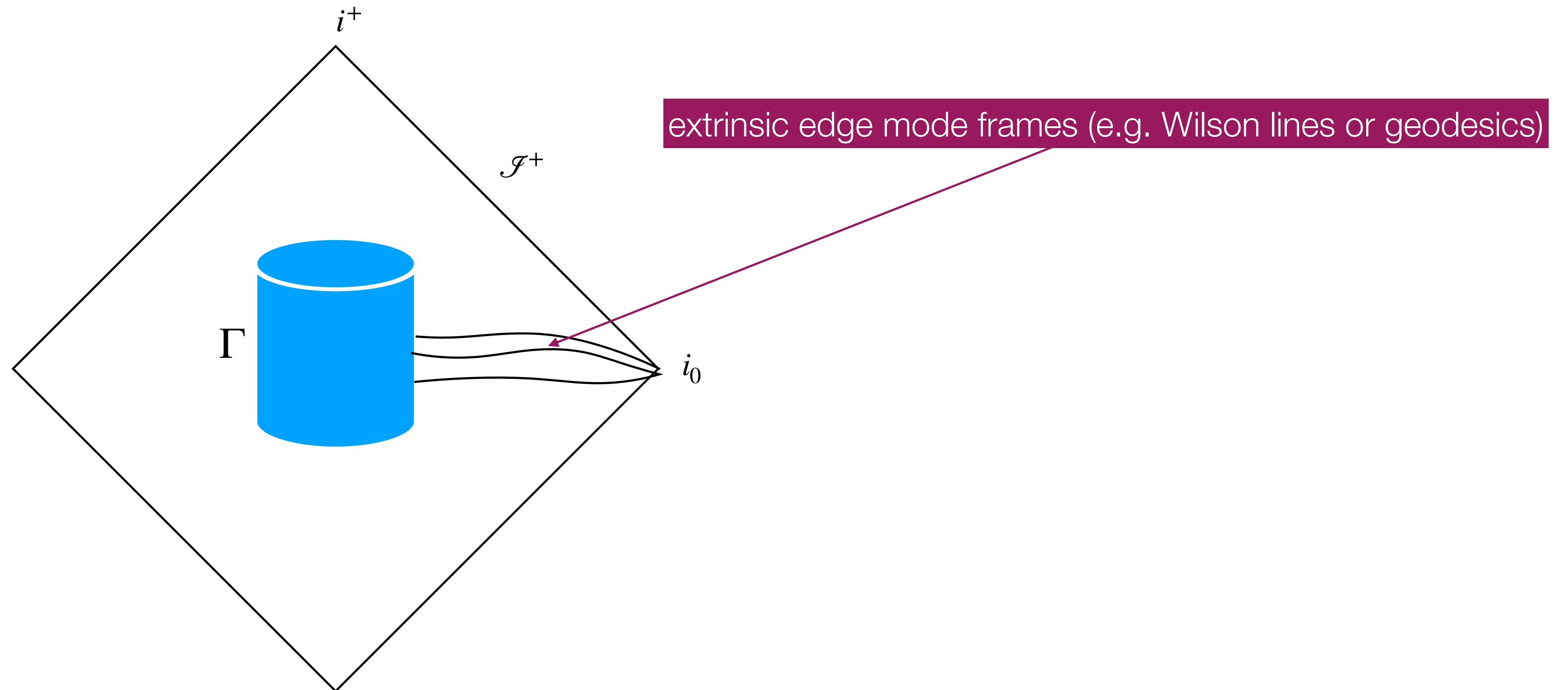


internal frame orientations	tetrad $e_a^\mu \in \text{SO}_+(3,1)$ (because $\eta_{\mu\nu} e_a^\mu e_b^\nu = \eta_{ab}$)	covariant POVM $E(g \cdot X) = U_R^g E(X \subset G) (U_R^g)^\dagger$ in practice, coh. states $ g\rangle$, $g \in G$
externally distinguishable configurations	$\{e_a^\mu, S^\mu\}$	kin. Hilbert space $\mathcal{H}_{\text{kin}} = \mathcal{H}_R \otimes \mathcal{H}_S$ kin. algebras $\mathcal{A}_{\text{kin}} = \mathcal{A}_R \otimes \mathcal{A}_S$
external frame/gauge transformations (acts on all subsystems)	$\Lambda^\mu{}_\nu \in \text{SO}_+(3,1)$	$U_R^g \otimes U_S^g$ unitary product rep of G
reorientations (acts only on the frame) commute with gauge transf.	$\Lambda_a{}^b \in \text{SO}_+(3,1)$	$V_R^g \otimes I_S$ s.t. $V_R^g g'\rangle = g'g^{-1}\rangle$

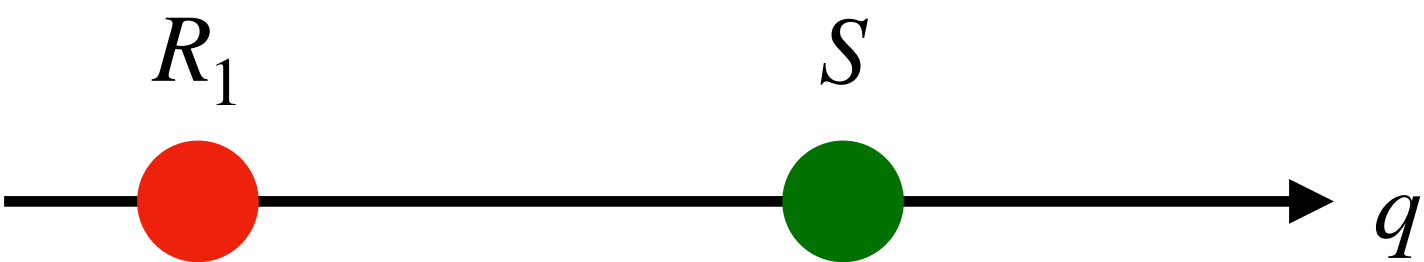
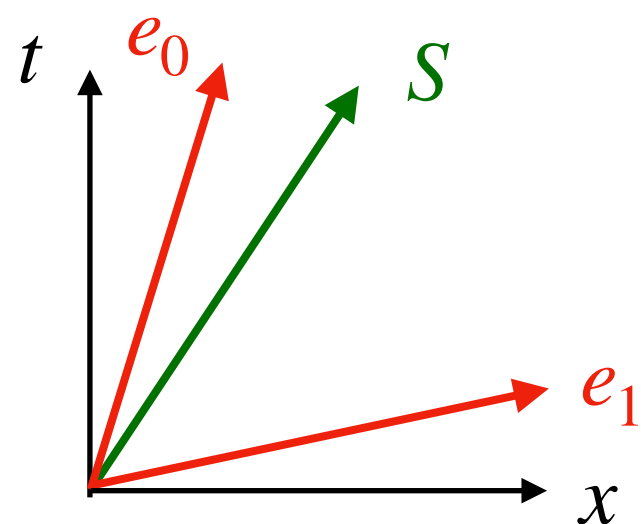
Reorientations are physical

- **relational time evolution** [PH, Smith, Lock '19; De Vuyst, Eccles, PH, Kirklin '24]
- **corner symmetries in gauge theory and gravity**

[Carrozza, PH '21; Carrozza, Eccles, PH '22; Goncalo-Araujo, PH, Sartini, Tomova '24]



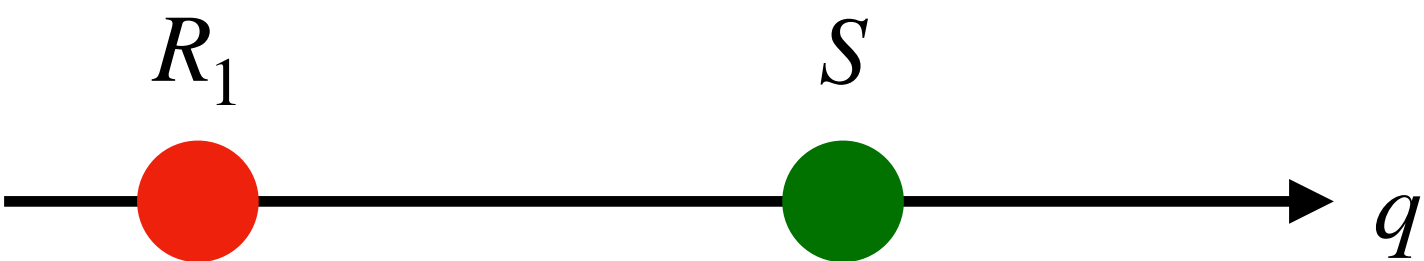
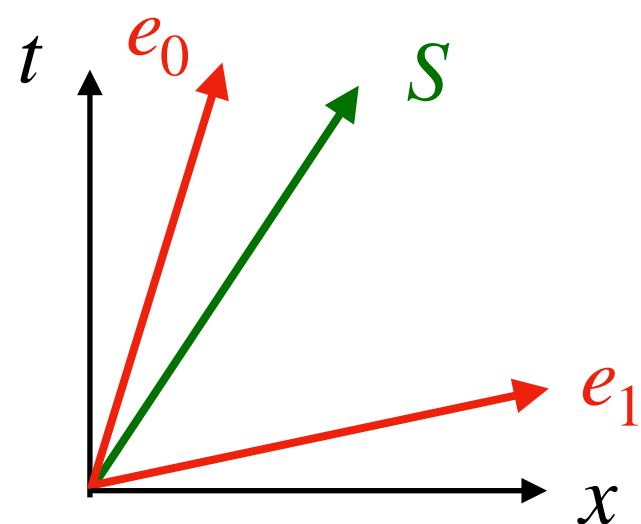
internal frames in SR vs. QT



**internally
distinguishable
configurations**
(extern frame indep.)

$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$
relational/dressed observables

internal frames in SR vs. QT



**internally
distinguishable
configurations**
(extern frame indep.)

$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$
relational/dressed observables

external frame indep. observables and states?

External frame independent observables

minimal condition for observables:

$$O \in \mathcal{A}_{\text{kin}} \text{ s.t. } [O, U_{RS}^g] = 0, \forall g \in G$$

External frame independent observables

minimal condition for observables:

$$O \in \mathcal{A}_{\text{kin}} \text{ s.t. } [O, U_{RS}^g] = 0, \forall g \in G$$

for simplicity assume G compact, then via **incoherent** group average (“ G -twirl”):

$\mathcal{G} : \mathcal{A}_{\text{kin}} \rightarrow \mathcal{A}_{\text{kin}}^G$ orthogonal projector onto G -inv. subalgebra

$$\mathcal{G}(\bullet) = \int_G dg U_{RS}^g(\bullet)(U_{RS}^g)^\dagger$$

External frame independent observables

minimal condition for observables:

$$O \in \mathcal{A}_{\text{kin}} \text{ s.t. } [O, U_{RS}^g] = 0, \forall g \in G$$

for simplicity assume G compact, then via **incoherent** group average (“ G -twirl”):

$$\mathcal{G} : \mathcal{A}_{\text{kin}} \rightarrow \mathcal{A}_{\text{kin}}^G \quad \text{orthogonal projector onto } G\text{-inv. subalgebra}$$

$$\mathcal{G}(\cdot) = \int_G dg \, U_{RS}^g(\cdot) (U_{RS}^g)^\dagger$$

$$\mathcal{A}_{\text{kin}} = \left\langle \boxed{O_{|R}^g(f_S)}, \boxed{V_R(g')} \mid f_S \in \mathcal{A}_S, g, g' \in G \right\rangle$$

(crossed product)

relational observables

QRF reorientations

External frame independent observables

minimal condition for observables:

$$O \in \mathcal{A}_{\text{kin}} \text{ s.t. } [O, U_{RS}^g] = 0, \forall g \in G$$

for simplicity assume G compact, then via **incoherent** group average (“ G -twirl”):

$$\mathcal{G} : \mathcal{A}_{\text{kin}} \rightarrow \mathcal{A}_{\text{kin}}^G \quad \text{orthogonal projector onto } G\text{-inv. subalgebra}$$

$$\mathcal{G}(\cdot) = \int_G dg U_{RS}^g(\cdot)(U_{RS}^g)^\dagger$$

$$\mathcal{A}_{\text{kin}} = \left\langle \boxed{O_{|R}^g(f_S)}, \boxed{V_R(g')} \mid f_S \in \mathcal{A}_S, g, g' \in G \right\rangle$$

(crossed product)

relational observables

QRF reorientations

$$O_{|R}^g(f_S) := \int_G dg' U_{RS}^{g'}(|g\rangle\langle g| \otimes f_S)(U_{RS}^{g'})^\dagger$$

(f_S conditional on QRF in orientation g)

Examples: relational observables

1. relational dynamics in gravity: some Hamiltonian constraint with clock $R = C$ $C = H_C + H_S$

$$O_{|C}^\tau(f_S) = \int dt e^{it(H_C+H_S)} (|\tau\rangle\langle\tau| \otimes f_S) e^{-it(H_C+H_S)} \quad \text{measures } f_S \text{ when clock reads } \tau$$

e.g. for **ideal clock** ($\text{Spec}(H_C) = \mathbb{R}$) have $O_C^0(f_S) = e^{-iTH_S} f_S e^{iTH_S}$

Examples: relational observables

1. relational dynamics in gravity: some Hamiltonian constraint with clock $R = C$

$$C = H_C + H_S$$

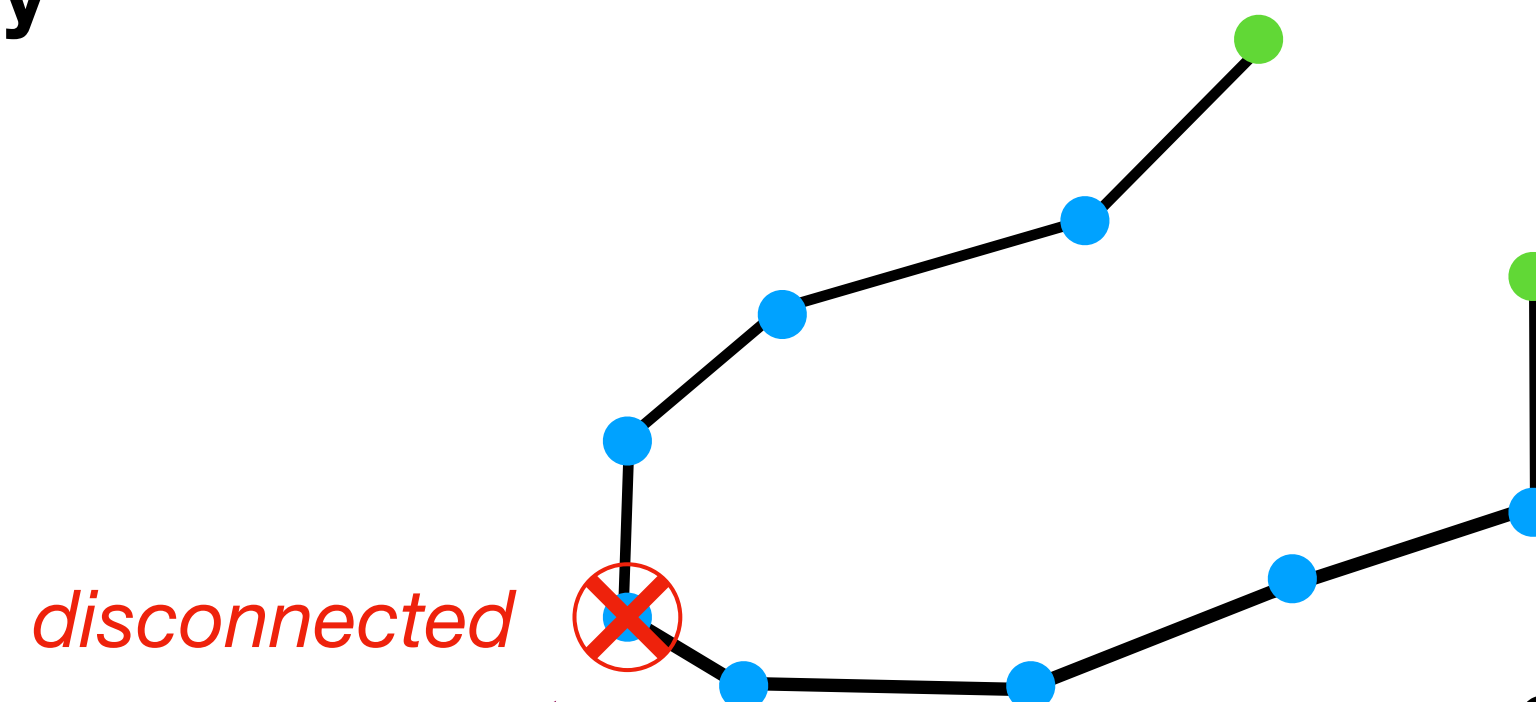
$$O_{|C}^\tau(f_S) = \int dt e^{it(H_C+H_S)} (|\tau\rangle\langle\tau| \otimes f_S) e^{-it(H_C+H_S)}$$

measures f_S when clock reads τ

e.g. for **ideal clock** ($\text{Spec}(H_C) = \mathbb{R}$) have

$$O_C^0(f_S) = e^{-iTH_S} f_S e^{iTH_S}$$

2. lattice gauge theory



2 Wilson lines define G -valued QRF R at green nodes

some root vertex in the lattice

Examples: relational observables

1. relational dynamics in gravity: some Hamiltonian constraint with clock $R = C$

$$C = H_C + H_S$$

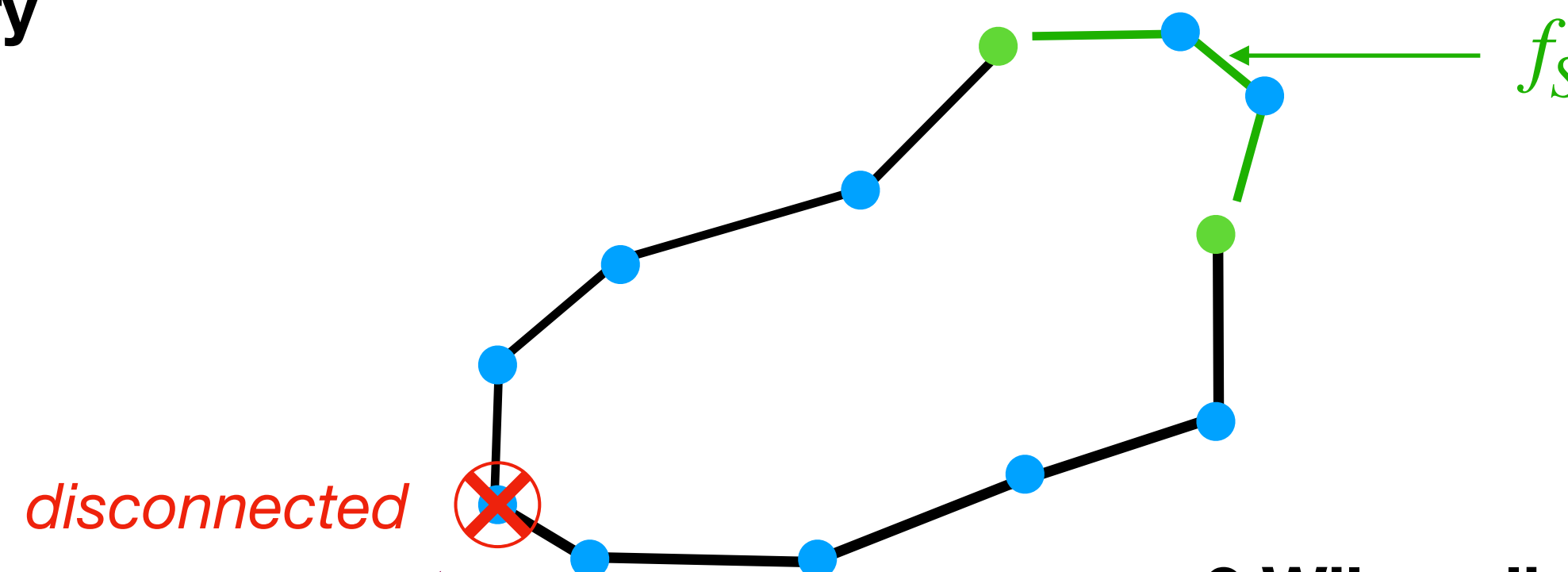
$$O_{|C}^\tau(f_S) = \int dt e^{it(H_C+H_S)} (|\tau\rangle\langle\tau| \otimes f_S) e^{-it(H_C+H_S)}$$

measures f_S when clock reads τ

e.g. for **ideal clock** ($\text{Spec}(H_C) = \mathbb{R}$) have

$$O_C^0(f_S) = e^{-iTH_S} f_S e^{iTH_S}$$

2. lattice gauge theory



2 Wilson lines define G -valued QRF R at green nodes

Examples: relational observables

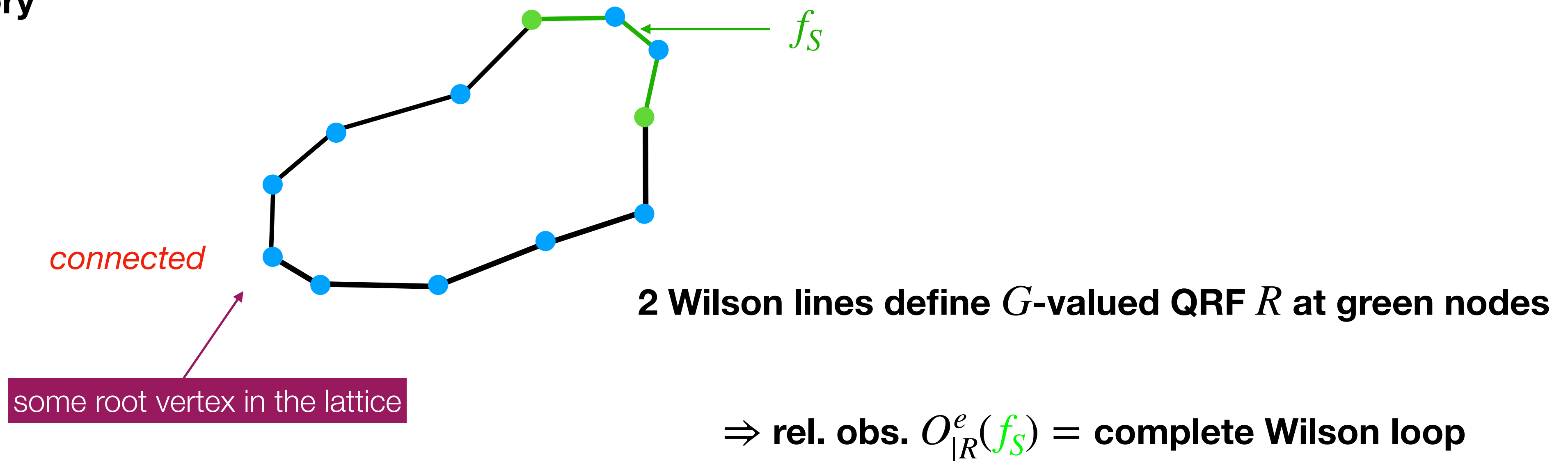
1. relational dynamics in gravity: some Hamiltonian constraint with clock $R = C$ $C = H_C + H_S$

$$O_{|C}^\tau(f_S) = \int dt e^{it(H_C+H_S)} (|\tau\rangle\langle\tau| \otimes f_S) e^{-it(H_C+H_S)} \quad \text{measures } f_S \text{ when clock reads } \tau$$

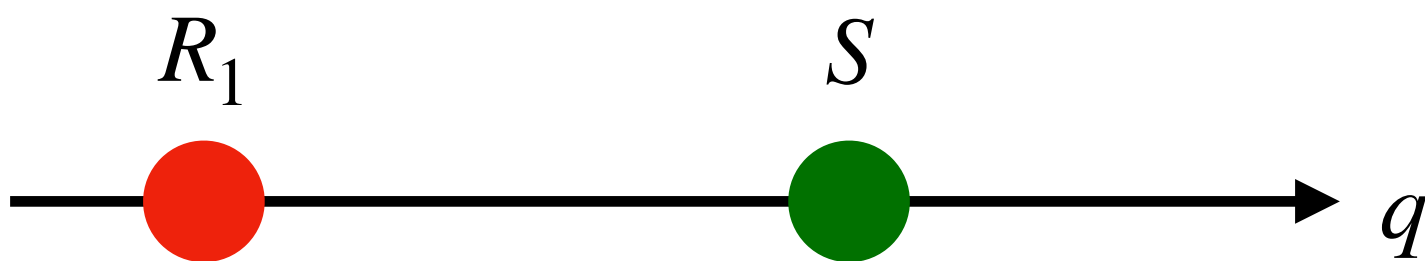
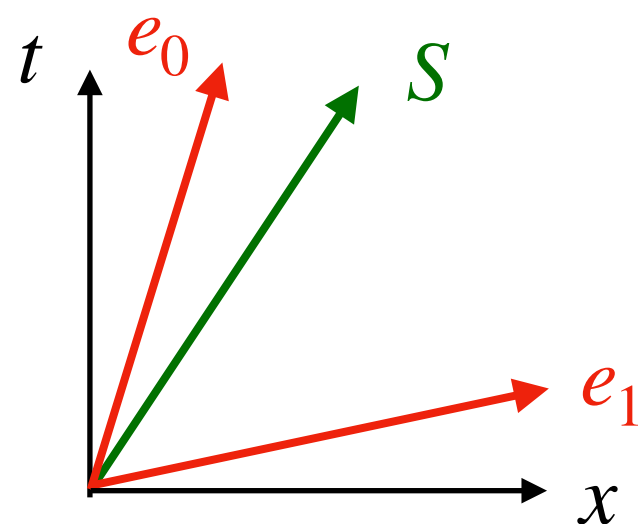
e.g. for **ideal clock** ($\text{Spec}(H_C) = \mathbb{R}$) have

$$O_C^0(f_S) = e^{-iTH_S} f_S e^{iTH_S}$$

2. lattice gauge theory



internal frames in SR vs. QT



**internally
distinguishable
configurations**
(extern frame indep.)

$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$
relational/dressed observables

observables: $O \in \mathcal{A}_{\text{kin}}$ s.t. $[O, U_{RS}^g] = 0, \forall g \in G$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g')

External frame independent observables

minimal condition for observables:

$$O \in \mathcal{A}_{\text{kin}} \text{ s.t. } [O, U_{RS}^g] = 0, \forall g \in G$$

for simplicity assume G compact, then via **incoherent** group average (“ G -twirl”):

$$\mathcal{G} : \mathcal{A}_{\text{kin}} \rightarrow \mathcal{A}_{\text{kin}}^G \quad \text{orthogonal projector onto } G\text{-inv. subalgebra}$$

$$\mathcal{G}(\cdot) = \int_G dg U_{RS}^g(\cdot)(U_{RS}^g)^\dagger$$

$$\mathcal{A}_{\text{kin}} = \left\langle \boxed{O_{|R}^g(f_S)}, \boxed{V_R(g')} \mid f_S \in \mathcal{A}_S, g, g' \in G \right\rangle$$

(crossed product)

relational observables

QRF reorientations

$$O_{|R}^g(f_S) := \int_G dg' U_{RS}^{g'}(|g\rangle\langle g| \otimes f_S)(U_{RS}^{g'})^\dagger$$

(f_S conditional on QRF in orientation g)

What states do they act on?

Two notions of external frame independence

states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

Two notions of external frame independence

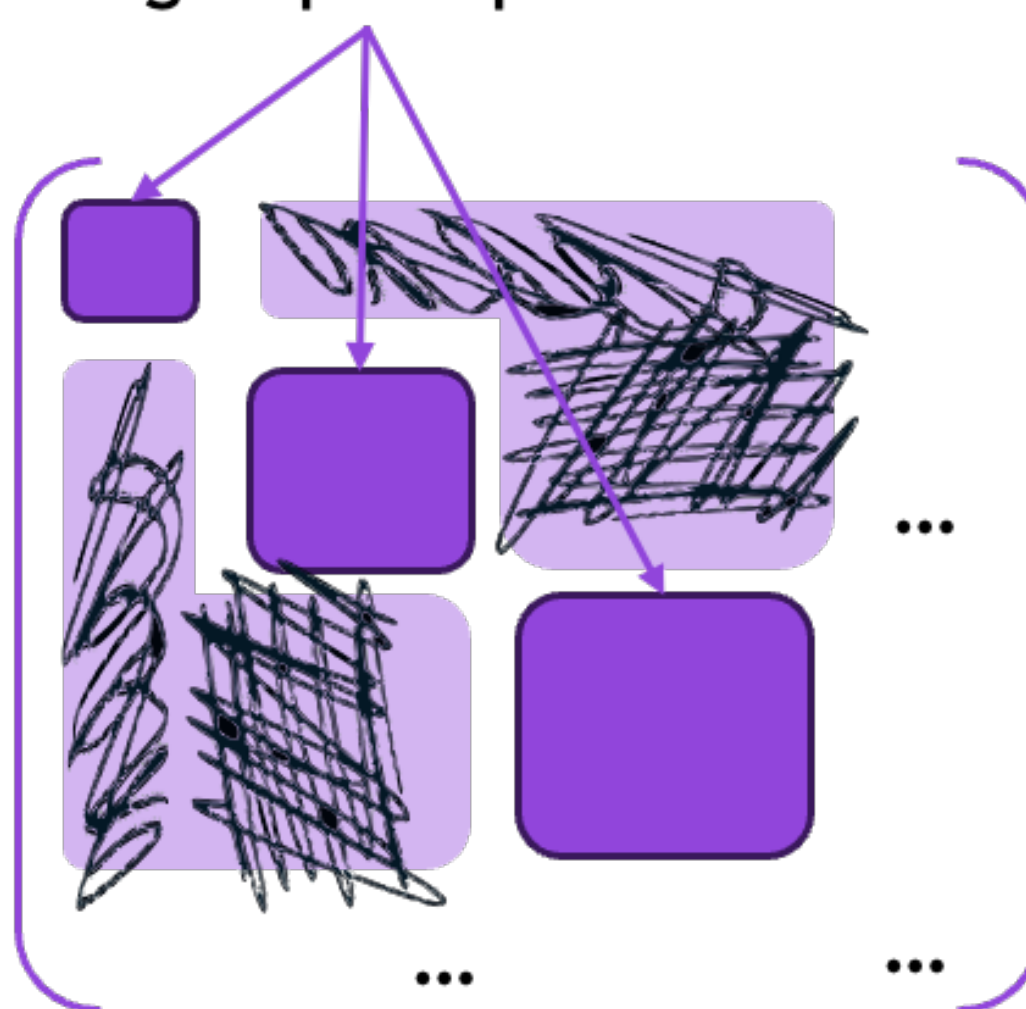
states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

image of the **incoherent group average**

$$\rho_{\text{inv}} = \mathcal{G}(\rho), \quad \rho \in \mathcal{S}(\mathcal{H}_{\text{kin}})$$

$$\mathcal{G}(\rho) = \int_G dg U(g) \rho U(g)^\dagger =$$


group irreps sectors

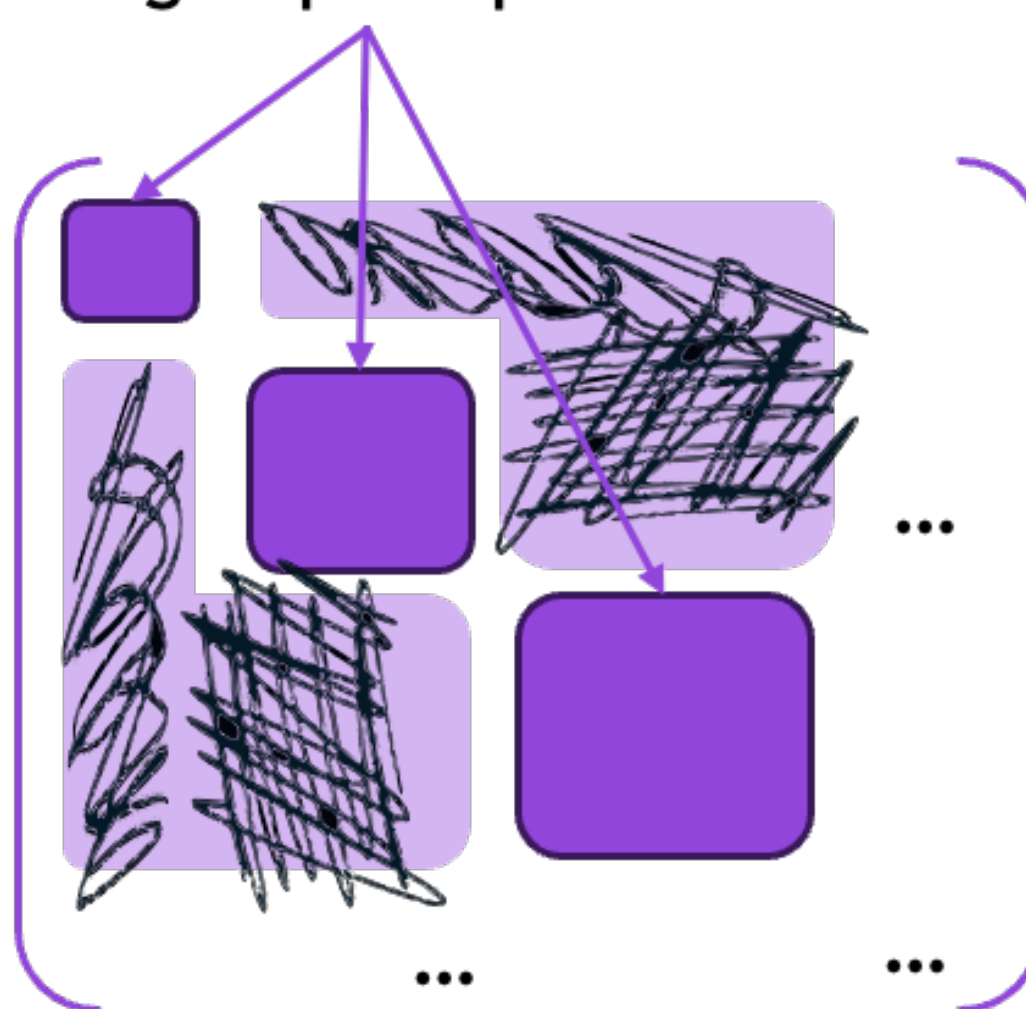
Two notions of external frame independence

states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

image of the **incoherent group average**

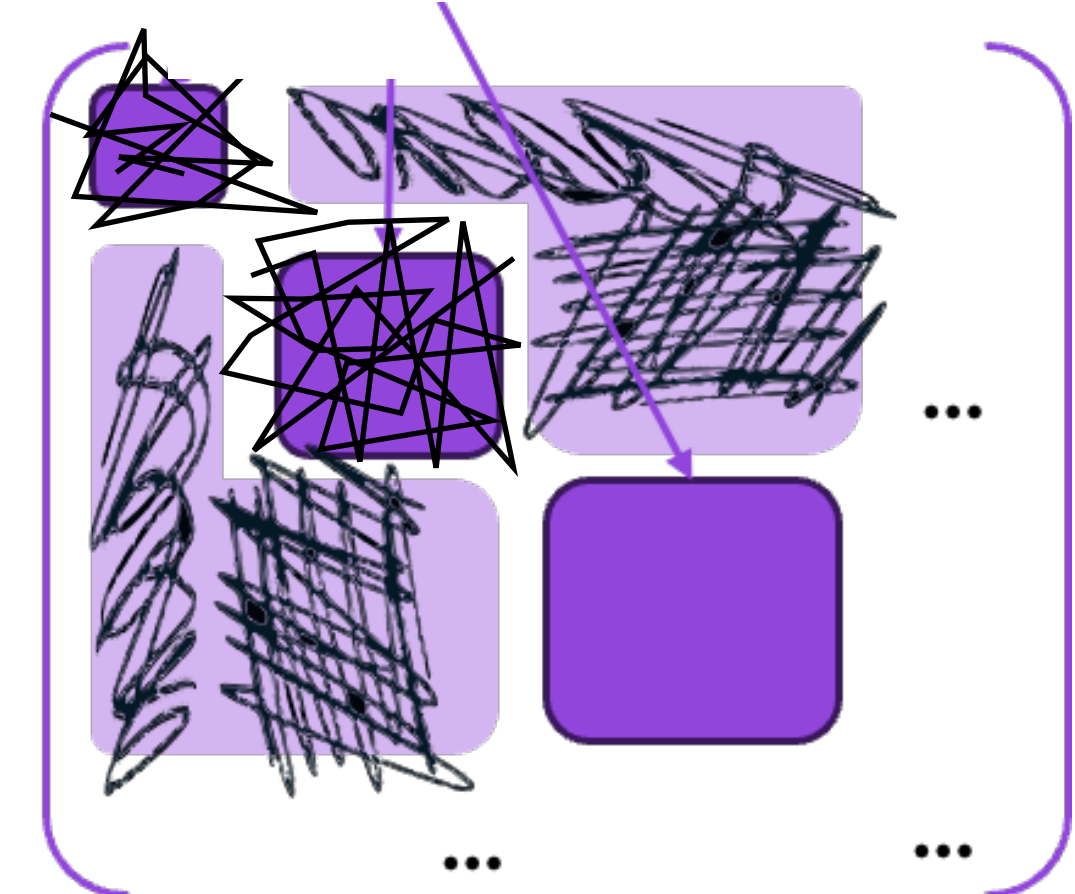
$$\rho_{\text{inv}} = \mathcal{G}(\rho), \quad \rho \in \mathcal{S}(\mathcal{H}_{\text{kin}})$$

$$\mathcal{G}(\rho) = \int_G dg U(g) \rho U(g)^\dagger =$$


strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

image of the **coherent group average**

$$|\psi\rangle = \Pi_{\text{pn}} |\psi\rangle, \quad |\psi\rangle \in \mathcal{H}_{\text{kin}} \quad \Pi_{\text{pn}} = \int_G dg U_{RS}^g$$

$$\Pi_{\text{pn}} \rho \Pi_{\text{pn}} =$$


Two notions of external frame independence

states can be G -invariant in two different ways:

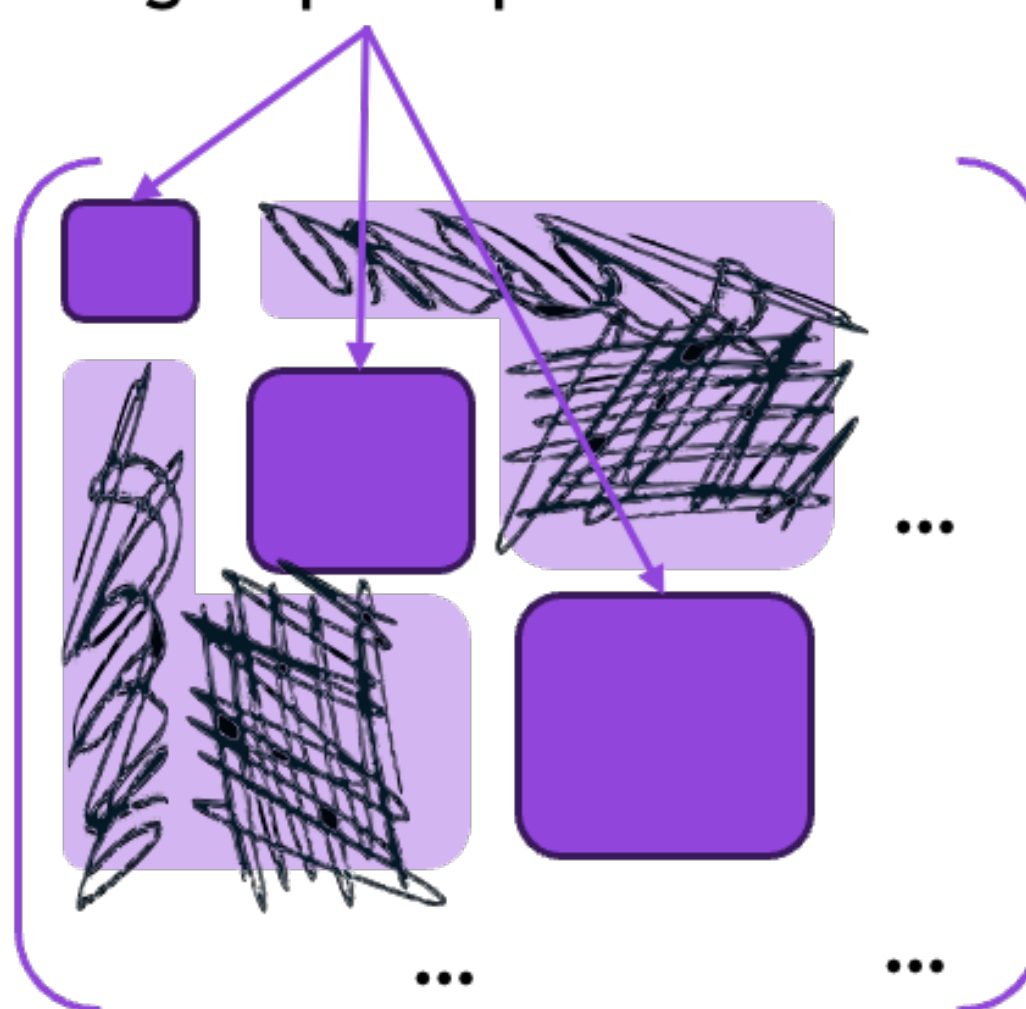
weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

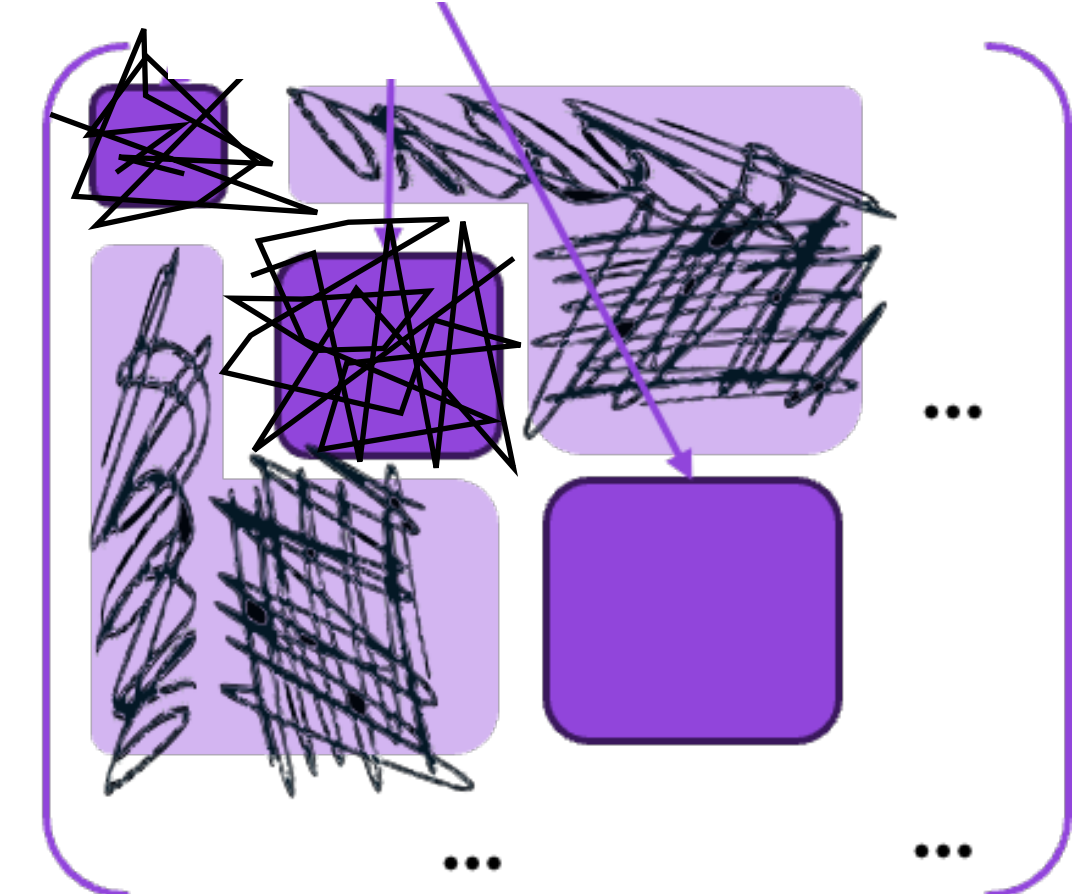
result of **one** equivalence relation:

$$|\psi\rangle \sim |\psi'\rangle \Leftrightarrow |\psi'\rangle = U_{RS}^g |\psi\rangle \text{ for some } g \in G \quad (*)$$

(“observer can’t distinguish between external frame transformed states”)

$$\mathcal{G}(\rho) = \int_G dg U(g) \rho U(g)^\dagger =$$


The diagram shows a 2x2 grid of purple squares representing different sectors of group irreps. The top-left square is solid purple. The top-right square is purple with diagonal hatching. The bottom-left square is purple with vertical hatching. The bottom-right square is solid purple. Purple arrows point from the text 'group irreps sectors' to each of the four squares. Ellipses (...) are placed to the right and below the grid, indicating more sectors.

$$\Pi_{\text{pn}} \rho \Pi_{\text{pn}} =$$


The diagram shows a 2x2 grid of purple squares representing charge zero sectors. The top-left square is purple with a complex, dense scribbled pattern. The top-right square is purple with diagonal hatching. The bottom-left square is purple with vertical hatching. The bottom-right square is solid purple. Purple arrows point from the text 'charge zero' to each of the four squares. Ellipses (...) are placed to the right and below the grid, indicating more sectors.

Two notions of external frame independence

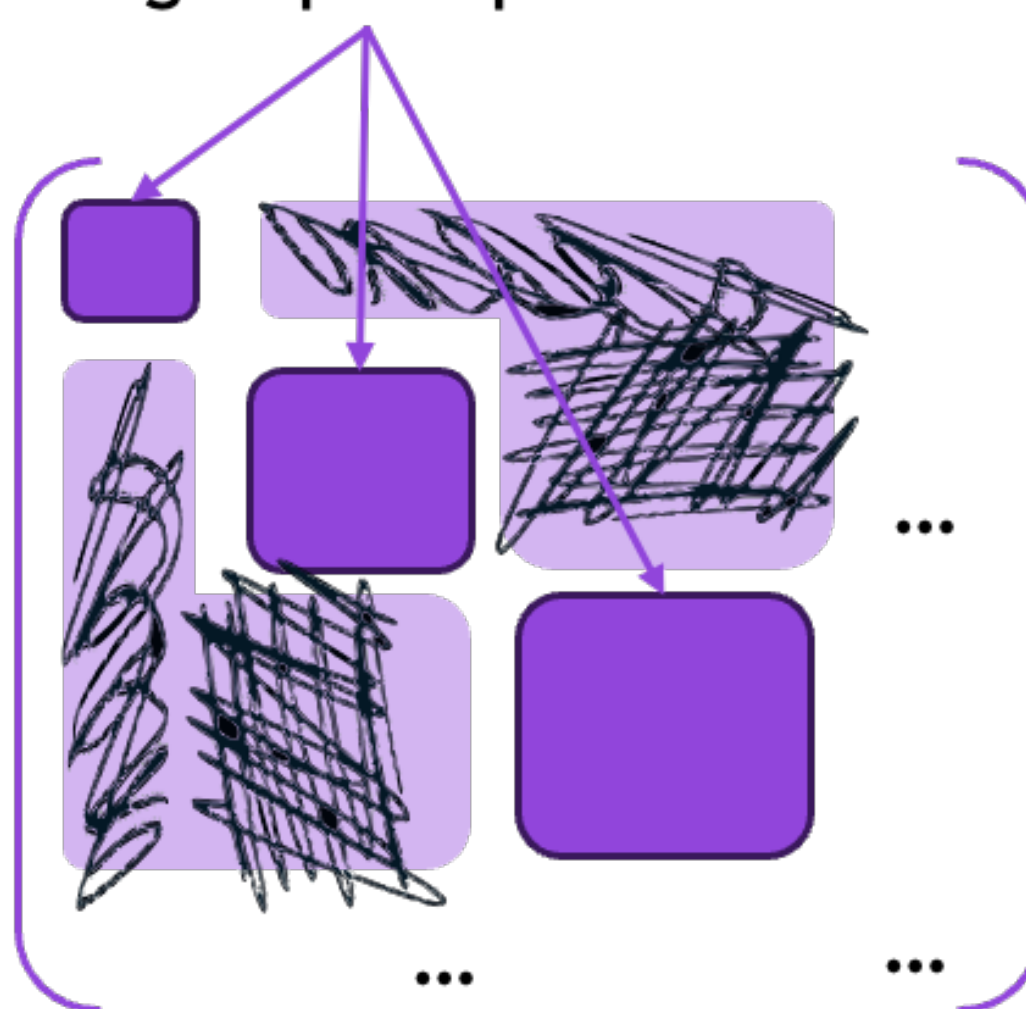
states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

result of **one** equivalence relation:

$$|\psi\rangle \sim |\psi'\rangle \Leftrightarrow |\psi'\rangle = U_{RS}^g |\psi\rangle \text{ for some } g \in G \quad (*)$$

(“observer can’t distinguish between external frame transformed states”)

$$\mathcal{G}(\rho) = \int_G dg U(g) \rho U(g)^\dagger =$$


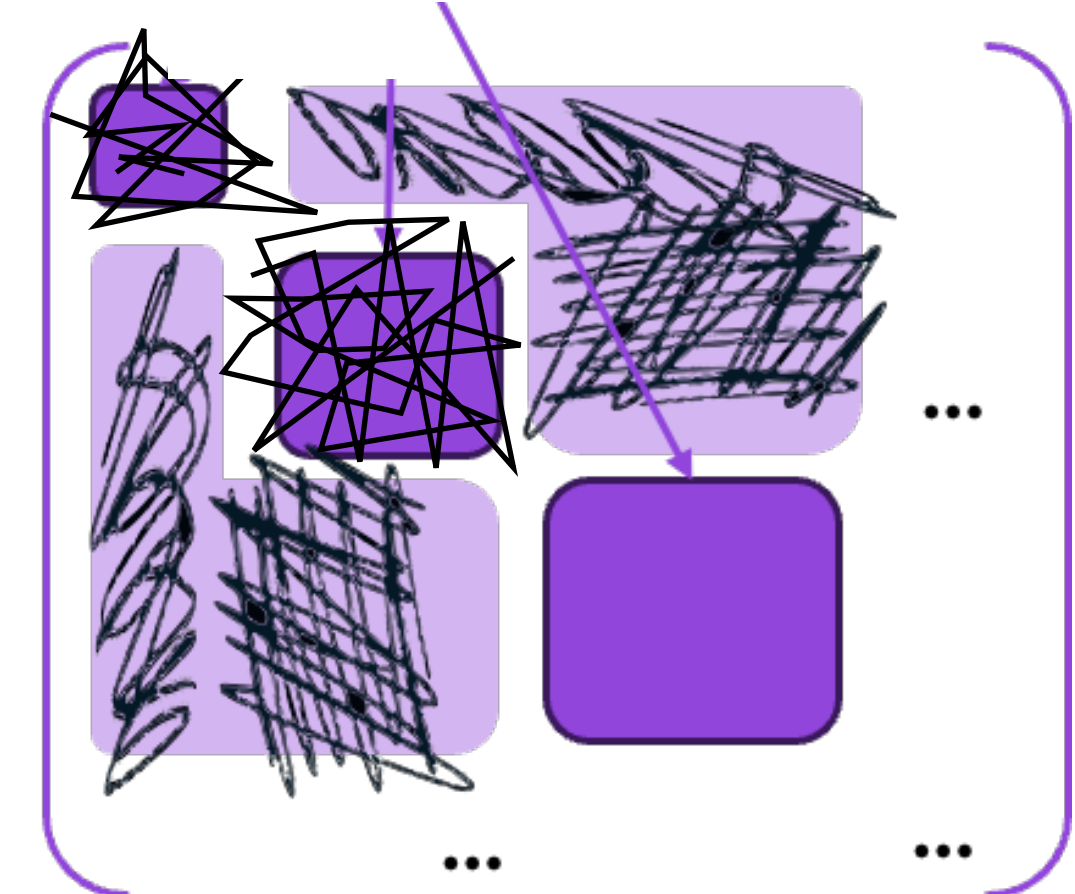
The diagram shows a large purple bracketed area containing several smaller purple squares. Some squares are solid purple, while others are filled with black scribbles. Arrows point from the text 'group irreps sectors' to the solid purple squares. Ellipses (...) are used to indicate that there are more squares and sectors than shown.

strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

result of **two** equivalence relations:

$$(*) + |\psi\rangle - U_{RS}^g |\psi\rangle \hat{\sim} 0, g \in G$$

(“observer can also not distinguish their difference from zero”)

$$\Pi_{\text{pn}} \rho \Pi_{\text{pn}} =$$


The diagram shows a large purple bracketed area containing several smaller purple squares. Some squares are solid purple, while others are filled with black scribbles. Arrows point from the text 'charge zero' to the solid purple squares. Ellipses (...) are used to indicate that there are more squares and sectors than shown.

Two notions of external frame independence

states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

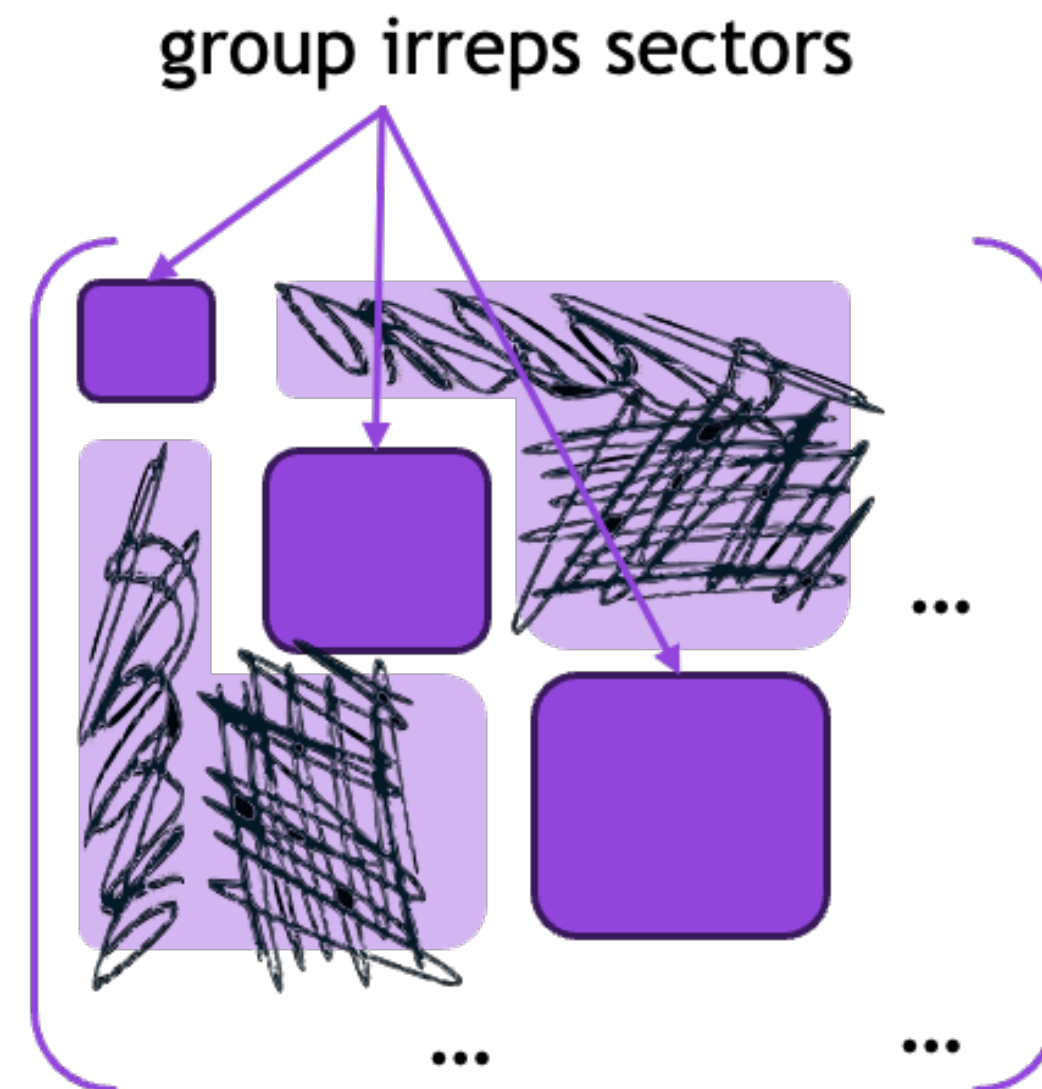
result of **one** equivalence relation:

$$|\psi\rangle \sim |\psi'\rangle \Leftrightarrow |\psi'\rangle = U_{RS}^g |\psi\rangle \text{ for some } g \in G \quad (*)$$

(“observer can’t distinguish between external frame transformed states”)

operational consequence:

CAN distinguish superpositions of external frame transformed states from their corresp. weighted mixtures

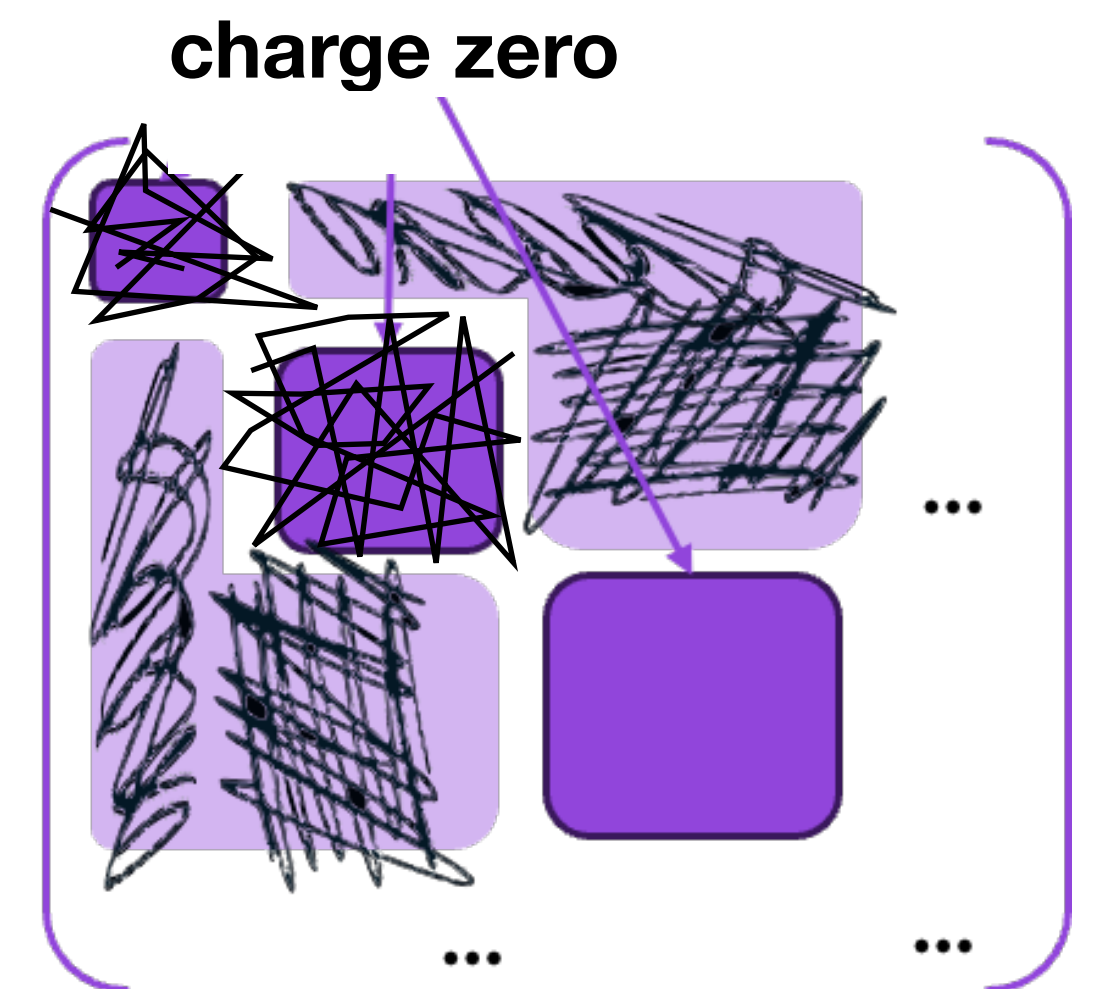


strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

result of **two** equivalence relations:

$$(*) + |\psi\rangle - U_{RS}^g |\psi\rangle \hat{\sim} 0, g \in G$$

(“observer can also not distinguish their difference from zero”)



Two notions of external frame independence

states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

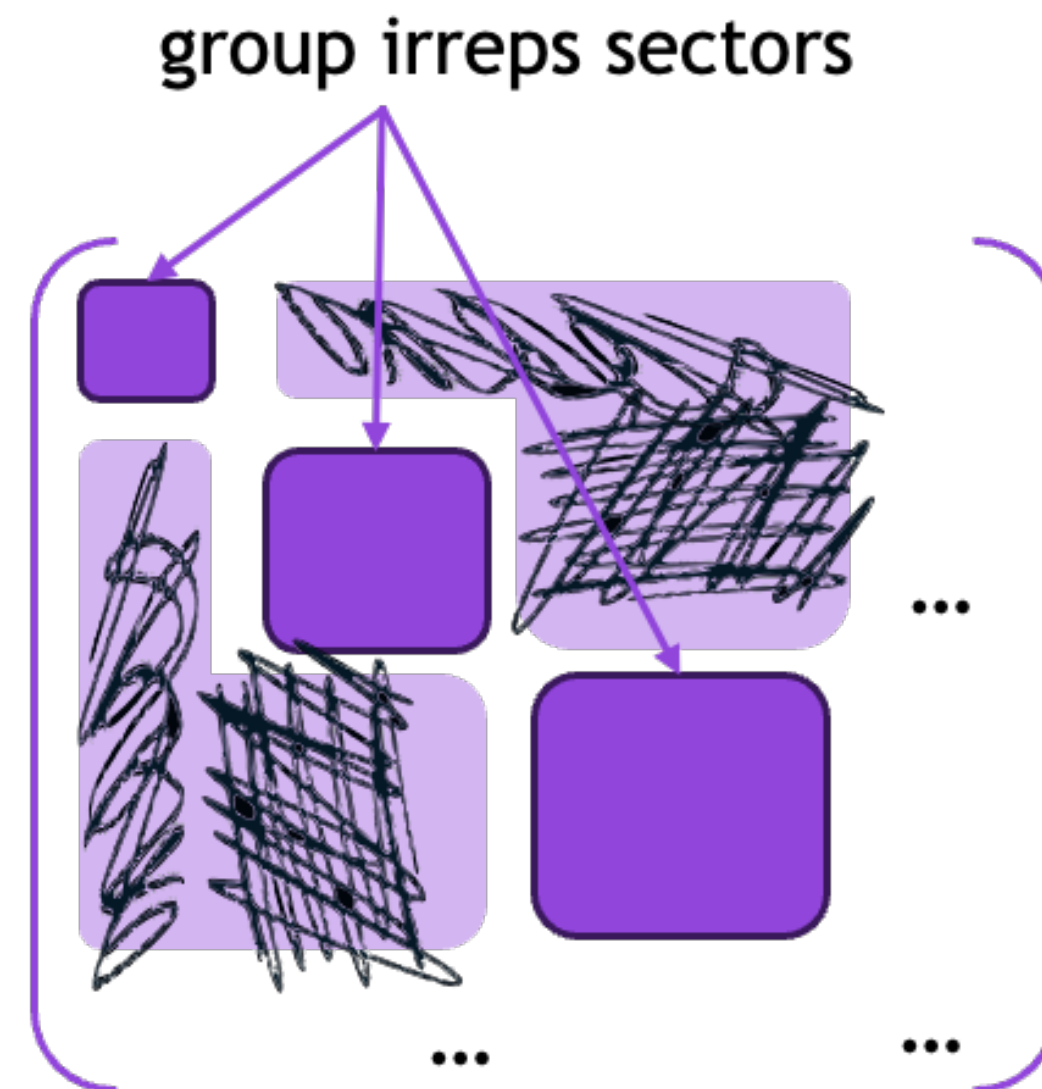
result of **one** equivalence relation:

$$|\psi\rangle \sim |\psi'\rangle \Leftrightarrow |\psi'\rangle = U_{RS}^g |\psi\rangle \text{ for some } g \in G \quad (*)$$

(“observer can’t distinguish between external frame transformed states”)

operational consequence:

CAN distinguish superpositions of external frame transformed states from their corresp. weighted mixtures



strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

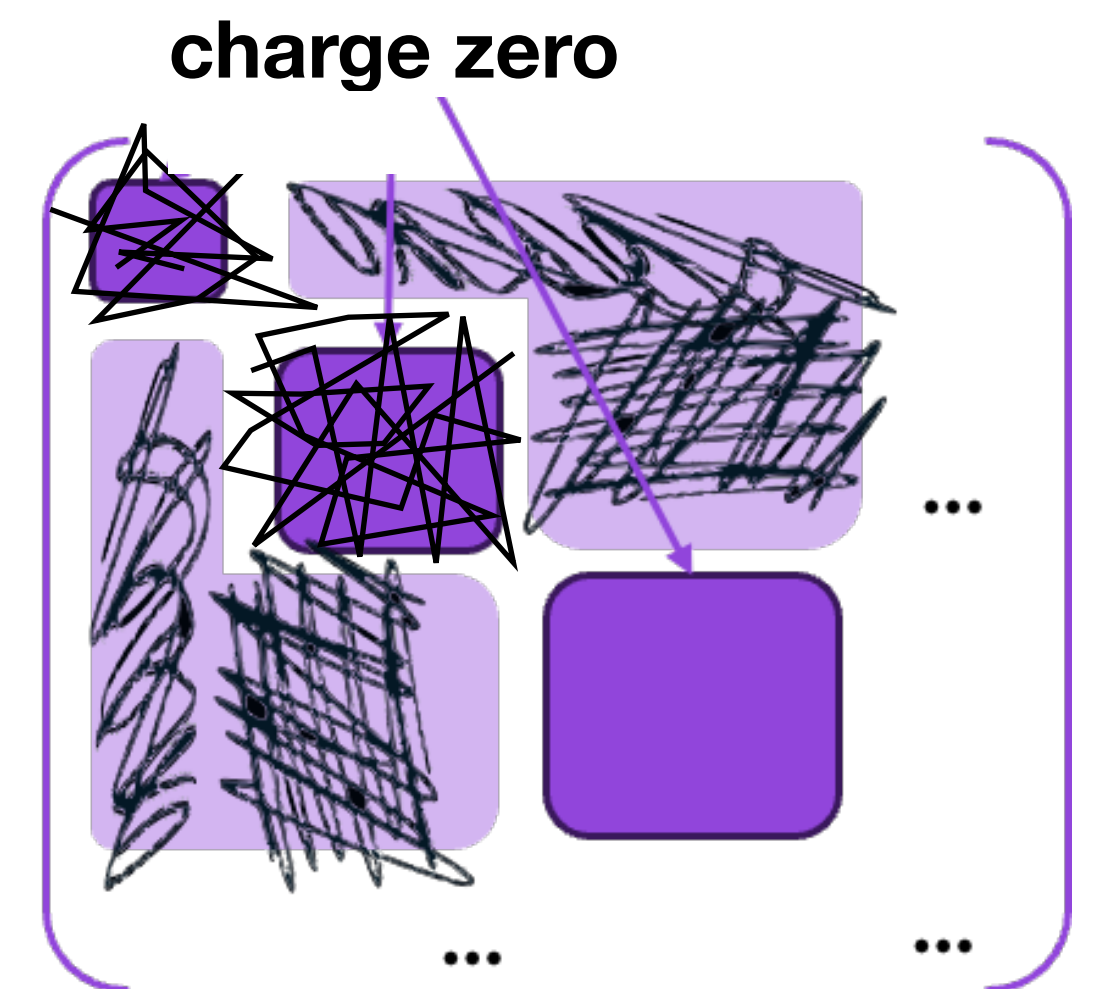
result of **two** equivalence relations:

$$(*) + |\psi\rangle - U_{RS}^g |\psi\rangle \hat{\sim} 0, g \in G$$

(“observer can also not distinguish their difference from zero”)

operational consequence:

CAN'T distinguish superpositions of external frame transformed states from their corresp. weighted mixtures



Two notions of external frame independence

states can be G -invariant in two different ways:

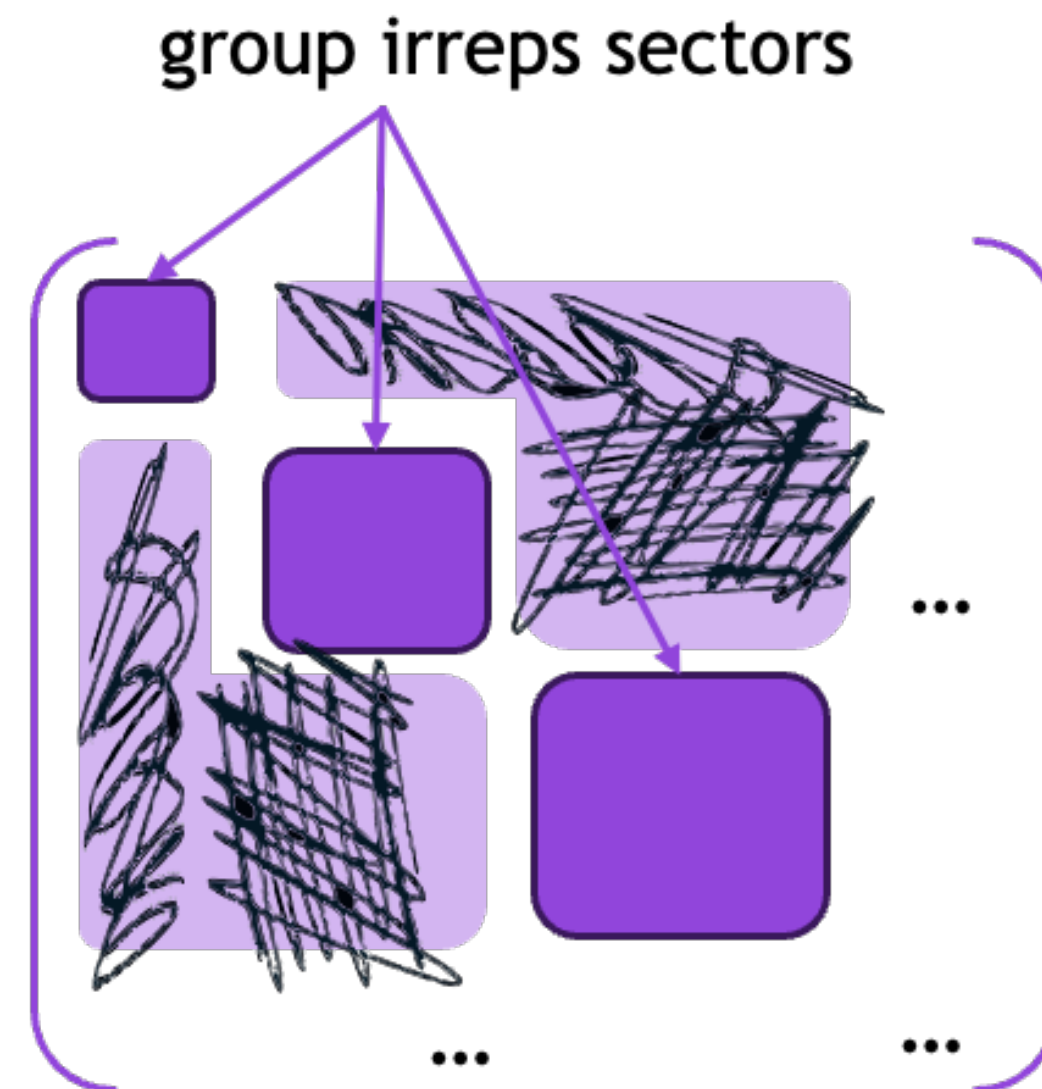
weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

result of **one** equivalence relation:

$$|\psi\rangle \sim |\psi'\rangle \Leftrightarrow |\psi'\rangle = U_{RS}^g |\psi\rangle \text{ for some } g \in G \quad (*)$$

(“observer can’t distinguish between external frame transformed states”)

QRF orientation data removed
(conjugate to U_{RS}^g),
but charge/irrep. data preserved

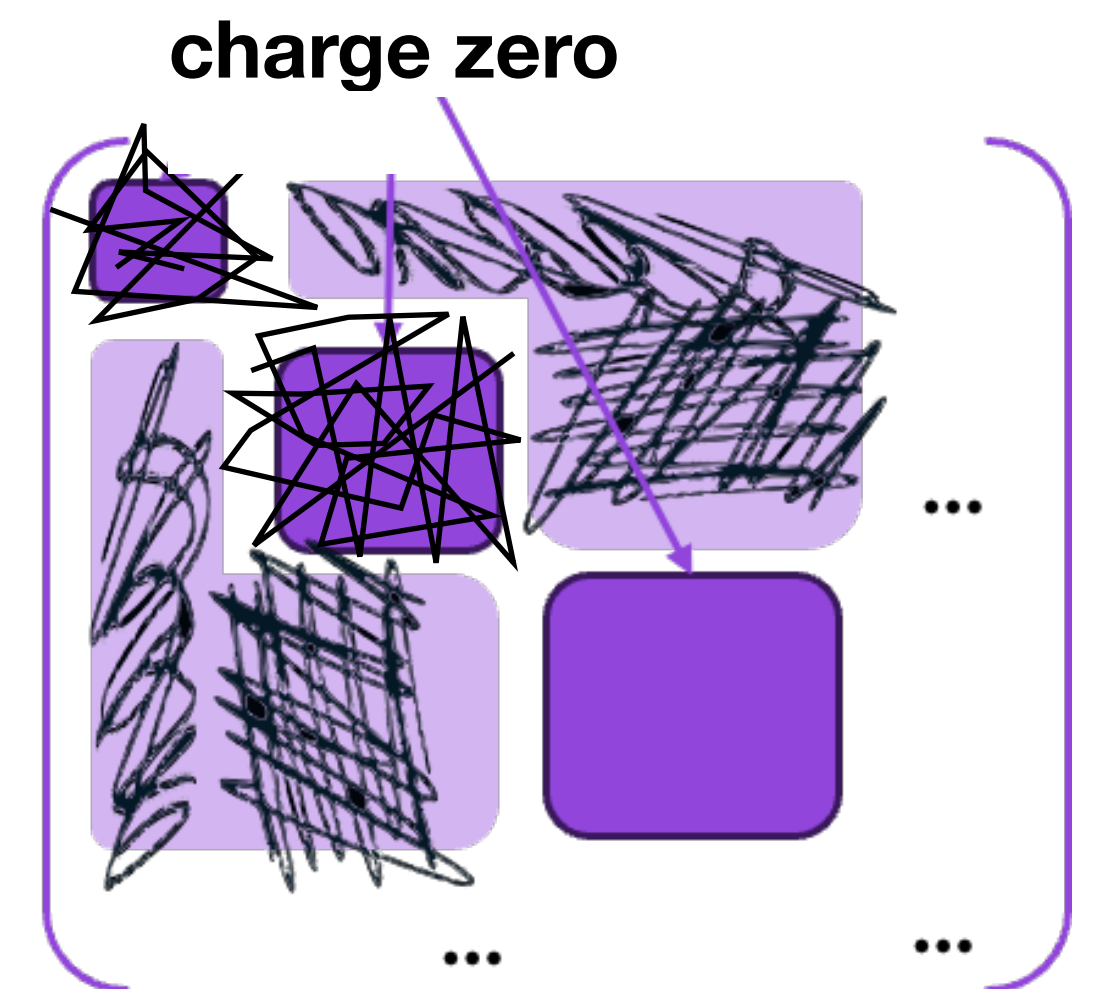


strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

result of **two** equivalence relations:

$$(*) + |\psi\rangle - U_{RS}^g |\psi\rangle \hat{\sim} 0, g \in G$$

(“observer can also not distinguish their difference from zero”)



Two notions of external frame independence

states can be G -invariant in two different ways:

weak: $[\rho_{\text{inv}}, U_{RS}^g] = 0, \forall g \in G$

result of **one** equivalence relation:

$$|\psi\rangle \sim |\psi'\rangle \Leftrightarrow |\psi'\rangle = U_{RS}^g |\psi\rangle \text{ for some } g \in G \quad (*)$$

(“observer can’t distinguish between external frame transformed states”)

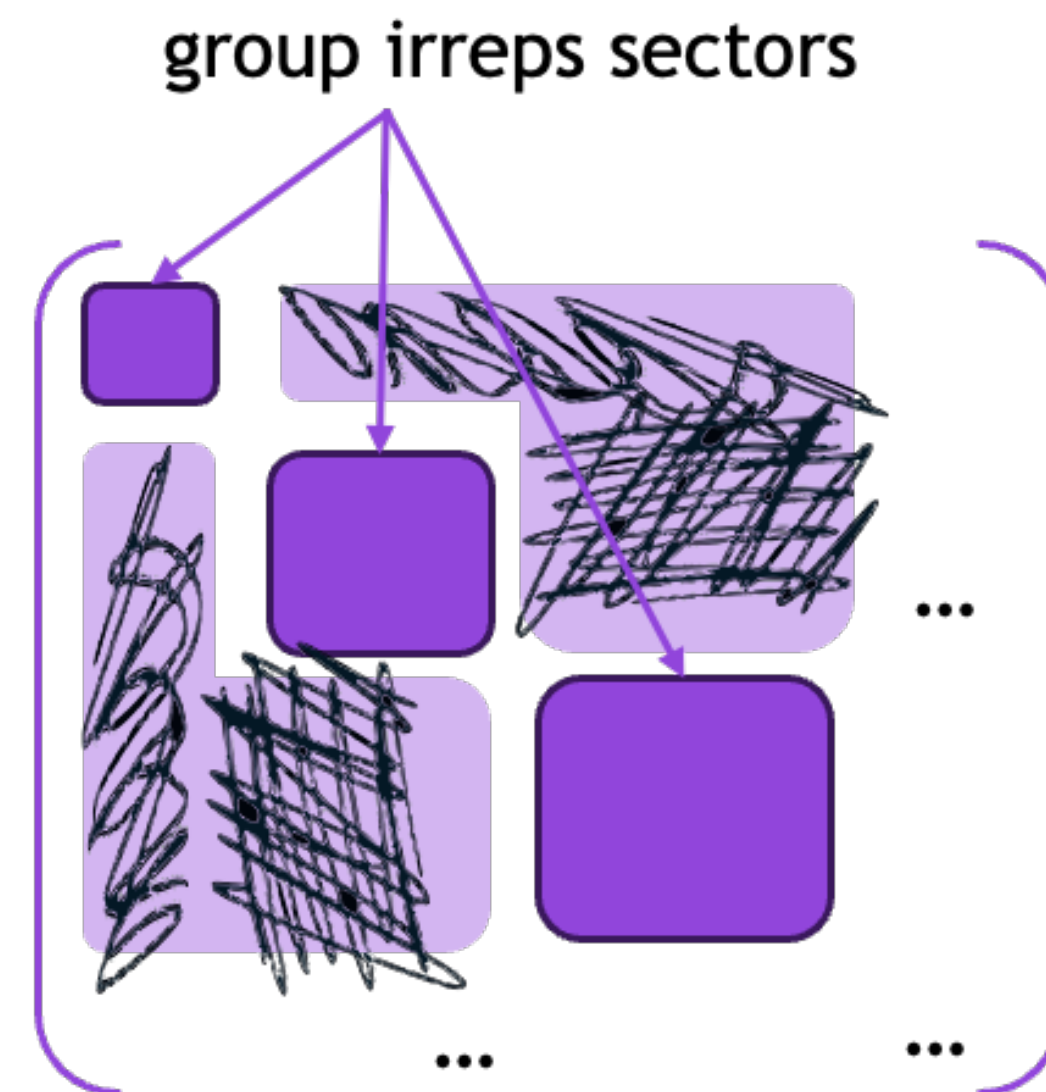
strong: $U_{RS}^g |\psi\rangle = |\psi\rangle, \forall g \in G$

result of **two** equivalence relations:

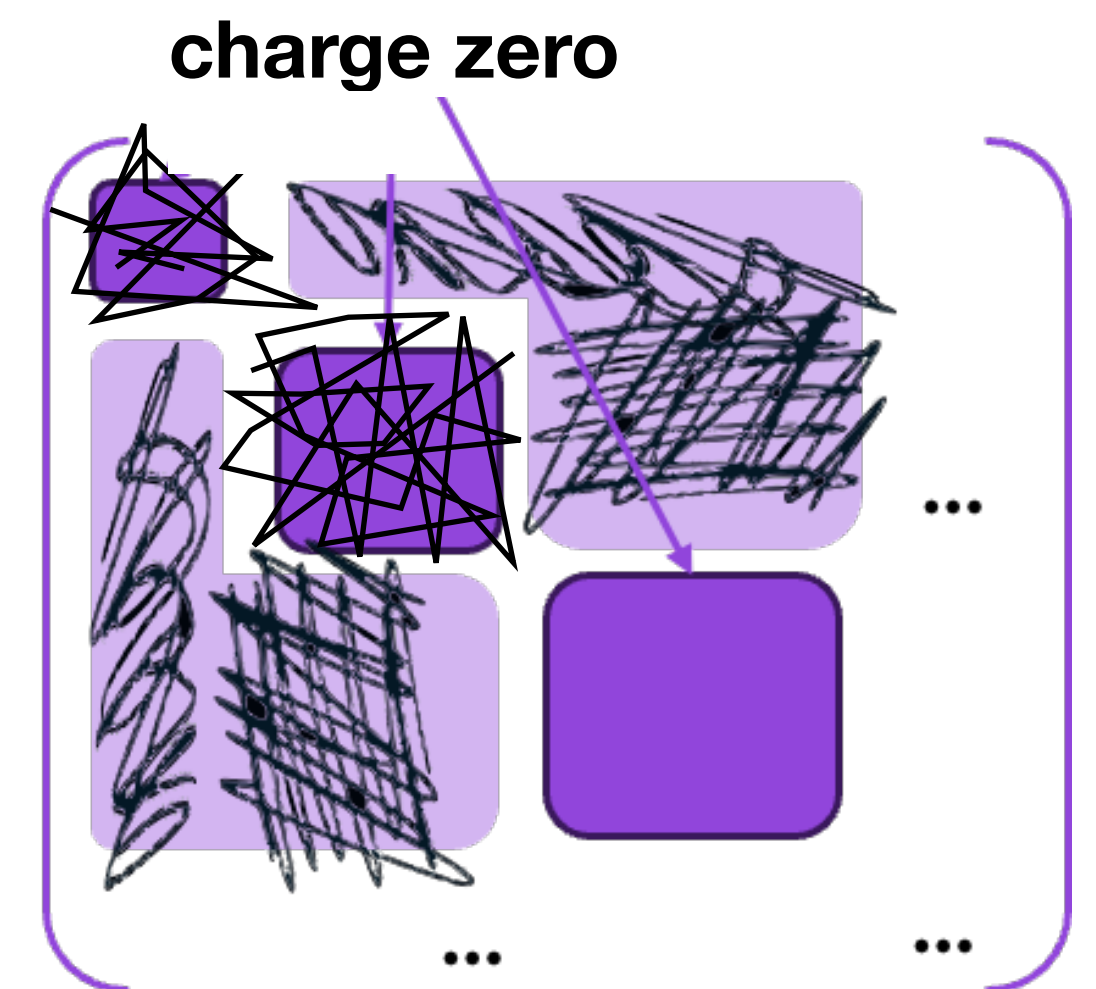
$$(*) + |\psi\rangle - U_{RS}^g |\psi\rangle \hat{\sim} 0, g \in G$$

(“observer can also not distinguish their difference from zero”)

QRF orientation data removed
(conjugate to U_{RS}^g),
but charge/irrep. data preserved



QRF orientation data removed
and
charge/irrep. data fixed



Example: particles on a line

$G = (\mathbb{R}, +)$ translation group



$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_R \otimes L^2(\mathbb{R})_S$$

and

$$U_{RS}^g \equiv U_{RS}(x) = \exp(-ixP), \quad x \in \mathbb{R}, \quad P = p_R + p_S$$

Example: particles on a line

$G = (\mathbb{R}, +)$ translation group



$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_R \otimes L^2(\mathbb{R})_S$$

and

$$U_{RS}^g \equiv U_{RS}(x) = \exp(-ixP), \quad x \in \mathbb{R}, \quad P = p_R + p_S$$

$$\mathcal{A}_{\text{kin}}^G = \langle q_S - q_R, p_S, P \rangle$$

$\Rightarrow$

frame orientation observable q_R eliminated

Example: particles on a line

$G = (\mathbb{R}, +)$ translation group



$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_R \otimes L^2(\mathbb{R})_S$$

and

$$U_{RS}^g \equiv U_{RS}(x) = \exp(-ixP), \quad x \in \mathbb{R}, \quad P = p_R + p_S$$

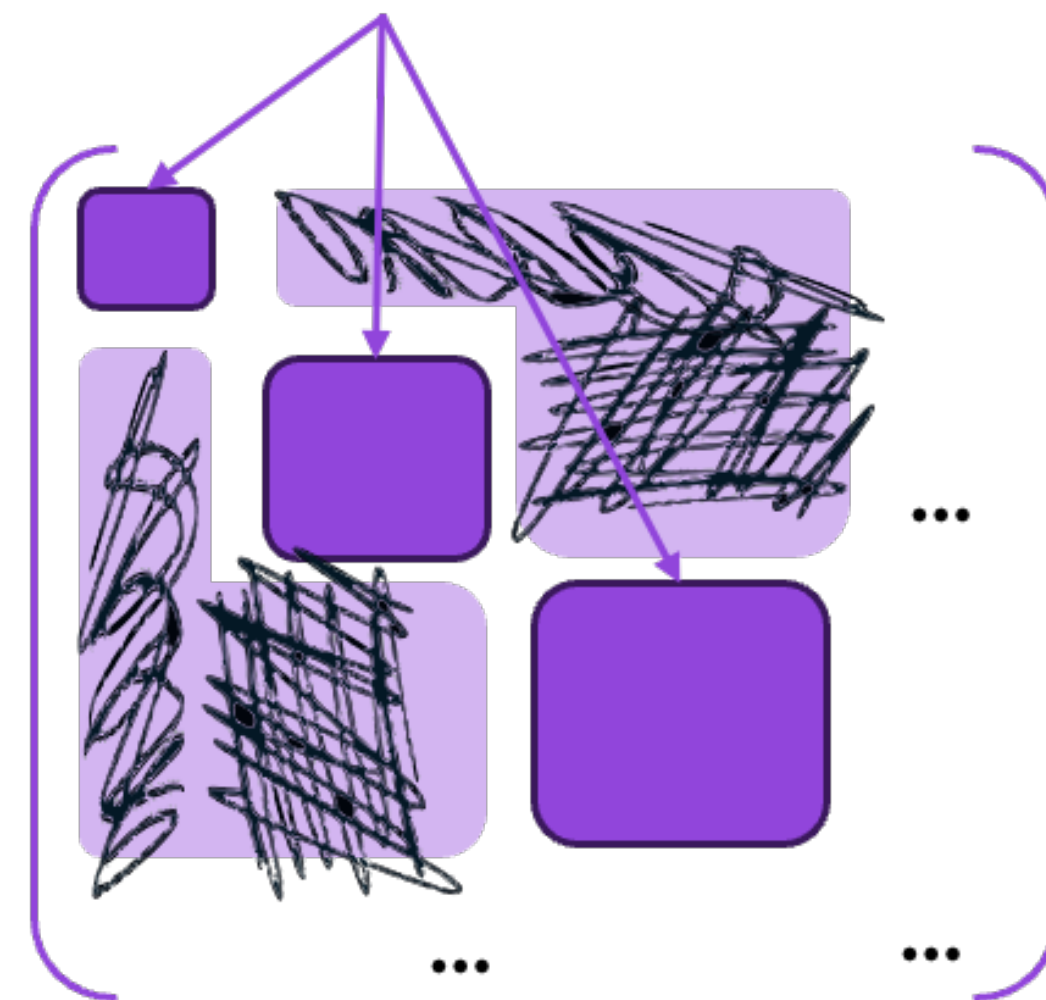
$$\mathcal{A}_{\text{kin}}^G = \langle q_S - q_R, p_S, P \rangle$$

$\Rightarrow$

frame orientation observable q_R eliminated

weak ext. frame indep.

P eigenspaces



$$\int_{\mathbb{R}} dx U_{RS}(x) \rho U_{RS}(x)^\dagger =$$

Example: particles on a line

$G = (\mathbb{R}, +)$ translation group



$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_R \otimes L^2(\mathbb{R})_S$$

and

$$U_{RS}^g \equiv U_{RS}(x) = \exp(-ixP), \quad x \in \mathbb{R}, \quad P = p_R + p_S$$

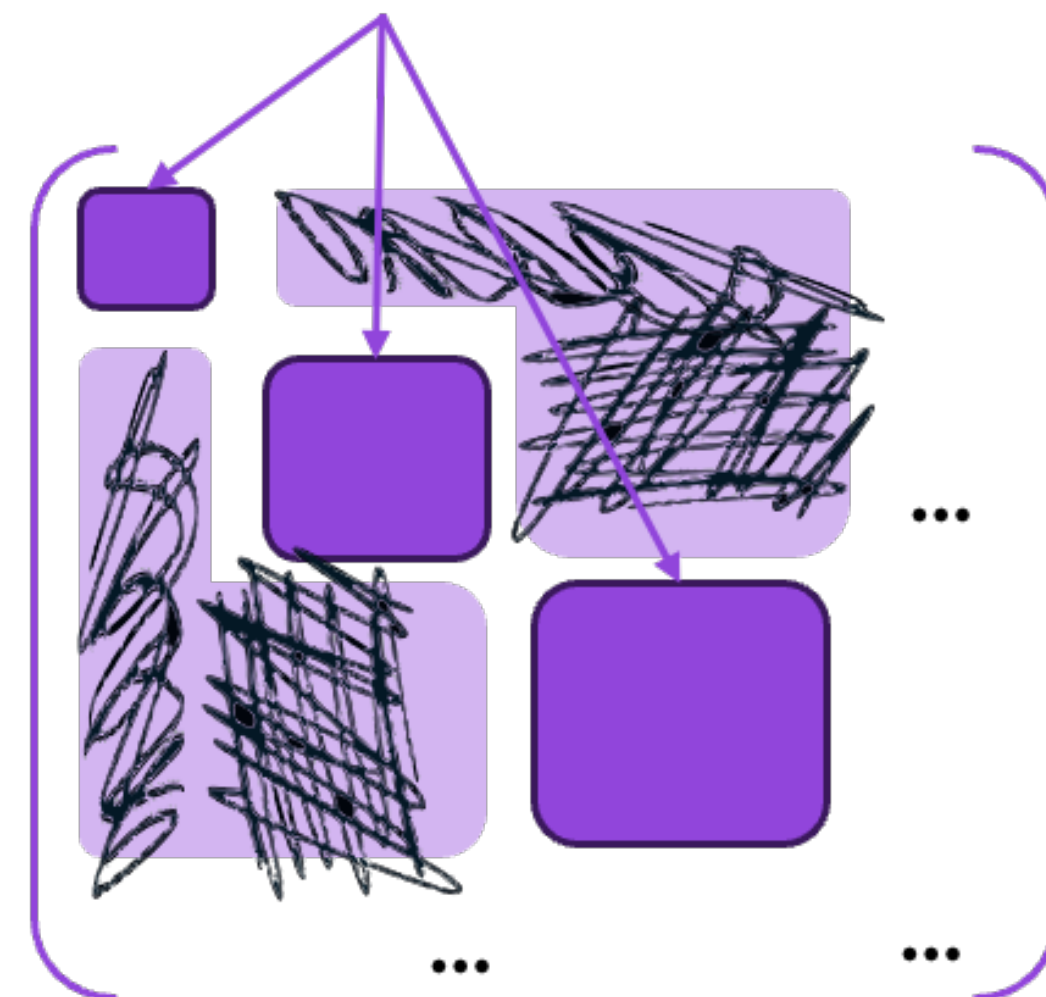
$$\mathcal{A}_{\text{kin}}^G = \langle q_S - q_R, p_S, P \rangle$$

$\Rightarrow$

frame orientation observable q_R eliminated

weak ext. frame indep.

P eigenspaces



$$\int_{\mathbb{R}} dx U_{RS}(x) \rho U_{RS}(x)^\dagger =$$

P preserved and conjugate to q_R
 $\Rightarrow$ “**half** of the QRF subsystem”
 eliminated

Example: particles on a line

$G = (\mathbb{R}, +)$ translation group



$$\mathcal{H}_{\text{kin}} = L^2(\mathbb{R})_R \otimes L^2(\mathbb{R})_S$$

and

$$U_{RS}^g \equiv U_{RS}(x) = \exp(-ixP), \quad x \in \mathbb{R}, \quad P = p_R + p_S$$

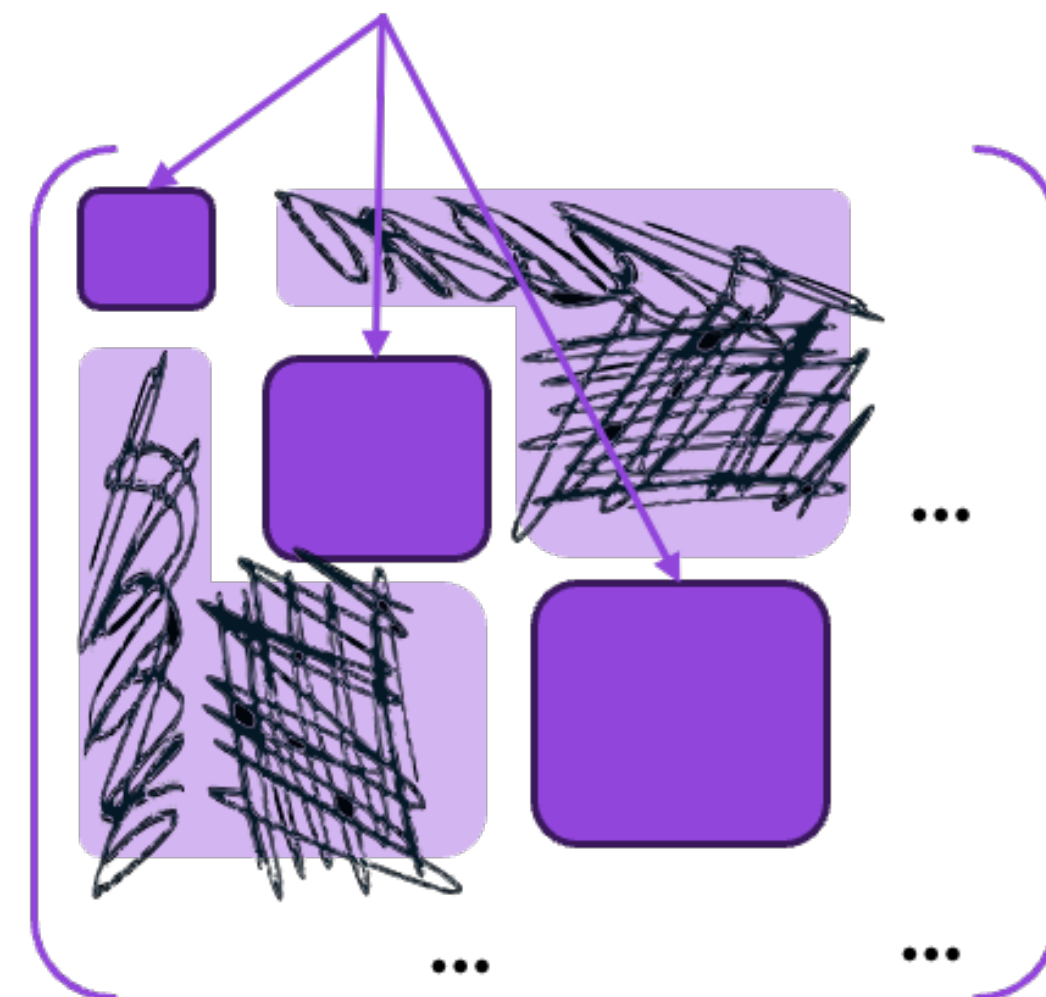
$$\mathcal{A}_{\text{kin}}^G = \langle q_S - q_R, p_S, P \rangle$$

$\Rightarrow$

frame orientation observable q_R eliminated

weak ext. frame indep.

P eigenspaces

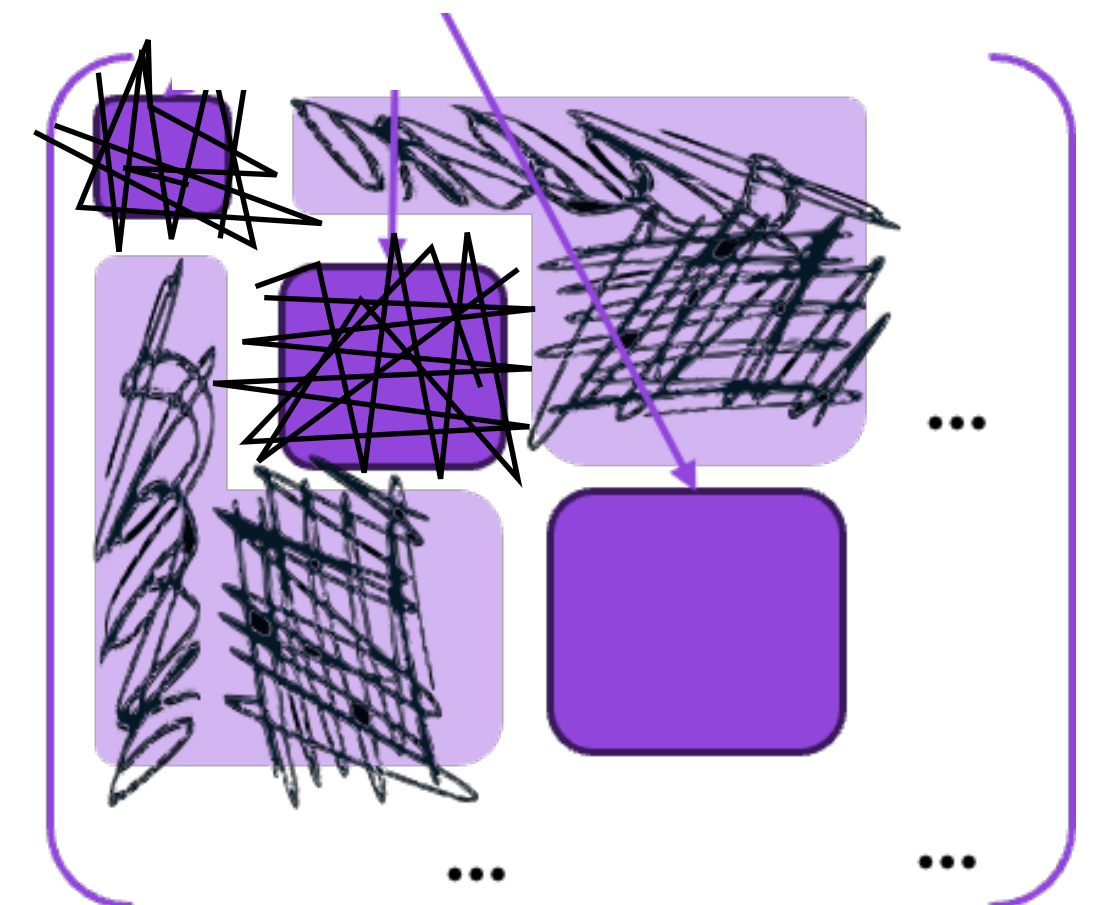


$$\int_{\mathbb{R}} dx U_{RS}(x) \rho U_{RS}(x)^\dagger =$$

P preserved and conjugate to q_R
 $\Rightarrow$ “**half** of the QRF subsystem”
 eliminated

strong ext. frame indep.

$P = 0$



$$\int_{\mathbb{R}} dx U_{RS}(x) \rho \int_{\mathbb{R}} dy U_{RS}(y) =$$

P also eliminated
 $\Rightarrow$ “**entire** QRF subsystem”
 eliminated $\Rightarrow$ fully redundant

Physical symmetry vs. redundancy

Given: symmetry group G (“external frame transformations”)

weak ext. frame indep.:

G physical

- lab frame transformations
- large gauge transformations in gauge theory and gravity
- corner symmetries in gauge theory and gravity
- ...

strong ext. frame indep.:

G param. redundancy

- small gauge transformations in gauge theory and gravity
- stabilizer symmetry in QEC
- quantum simulations
- operational restrictions
- ...

are the orientations of the QRF (in principle) measurable?

Yes

No

(but requires access to an ext. frame)

Potpourri of approaches to QRF covariance

perspectival

[Giacomini, Castro-Ruiz, Brukner '17; de la Hamette, Galley '20; de la Hamette, Kabel, Castro-Ruiz, Brukner '22,...]

extra-particle

[Castro-Ruiz, Oreshkov '21; Garnier, Hausmann, Castro-Ruiz '25]

quantum information-theoretic

[Aharonov, Susskind '67; Bartlett, Rudolph, Spekkens '06; Krumm, PH, Müller '20]

effective (semiclassical)

[Bojowald, PH, Tsobanjan '10; PH, Kubalova, Tsobanjan '12]

algebraic I

[Bojowald, Tsobanjan '19; '22; De Vuyst, PH, Tsobanjan *to appear*]

operational

[Fewster et al. '24; Carette, Glowacki, Loveridge '23; Glowacki, Loveridge, Waldron '23]

algebraic II

Ahmad, Klinger, Leigh '24

perspective-neutral

[Vanrietvelde, PH, Giacomini, Castro-Ruiz '18; de la Hamette, Galley, PH, Loveridge, Müller '21; PH, Smith, Lock '19; PH, Kotecha, Mele '23; ...]



Potpourri of approaches to QRF covariance

weak ext. frame indep.

perspectival

[Giacomini, Castro-Ruiz, Brukner '17; de la Hamette, Galley '20;
de la Hamette, Kabel, Castro-Ruiz, Brukner '22,...]

**equiv. to strong ext. frame
indep. for ideal QRFs**

extra-particle

[Castro-Ruiz, Oreshkov '21; Garnier, Hausmann, Castro-Ruiz '25]

effective (semiclassical)

[Bojowald, PH, Tsobanjan '10; PH, Kubalova, Tsobanjan '12]

quantum information-theoretic

[Aharonov, Susskind '67; Bartlett, Rudolph, Spekkens '06;
Krumm, PH, Müller '20]

algebraic I

[Bojowald, Tsobanjan '19; '22; De Vuyst, PH, Tsobanjan '25]

operational

[Fewster et al. '24; Carette, Glowacki, Loveridge '23;
Glowacki, Loveridge, Waldron '23]

strong ext. frame indep.

algebraic II

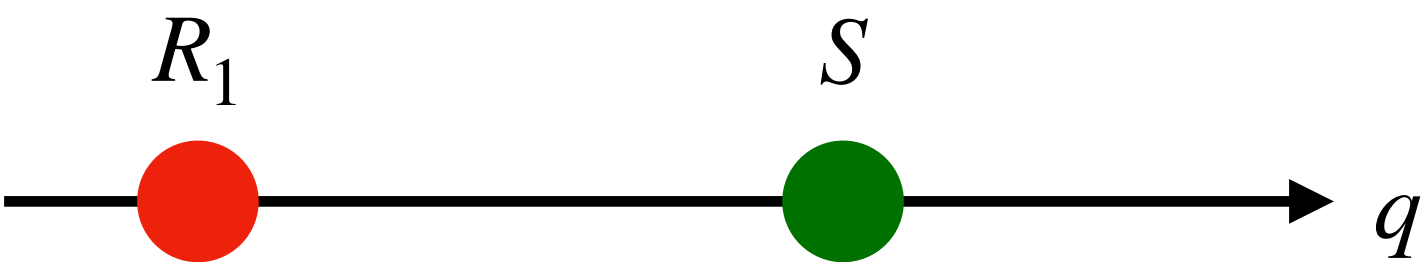
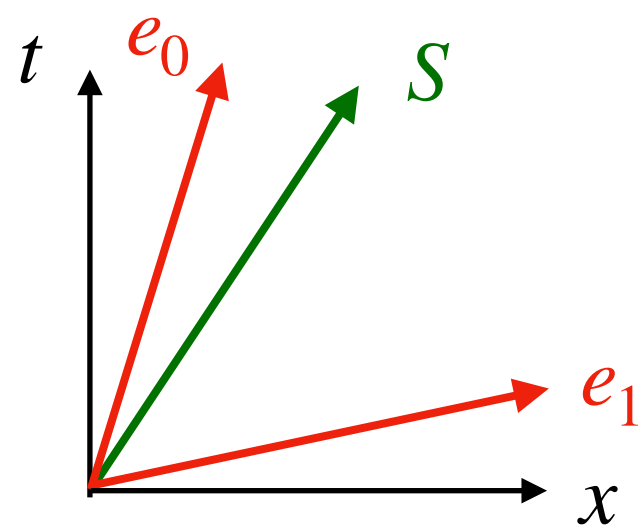
Ahmad, Klinger, Leigh '24

perspective-neutral

[Vanrietvelde, PH, Giacomini, Castro-Ruiz '18; de la Hamette, Galley,
PH, Loveridge, Müller '21; PH, Smith, Lock '19; PH, Kotecha, Mele '23; ...]



internal frames in SR vs. QT



**internally
distinguishable
configurations**
(extern frame indep.)

$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$
relational/dressed observables

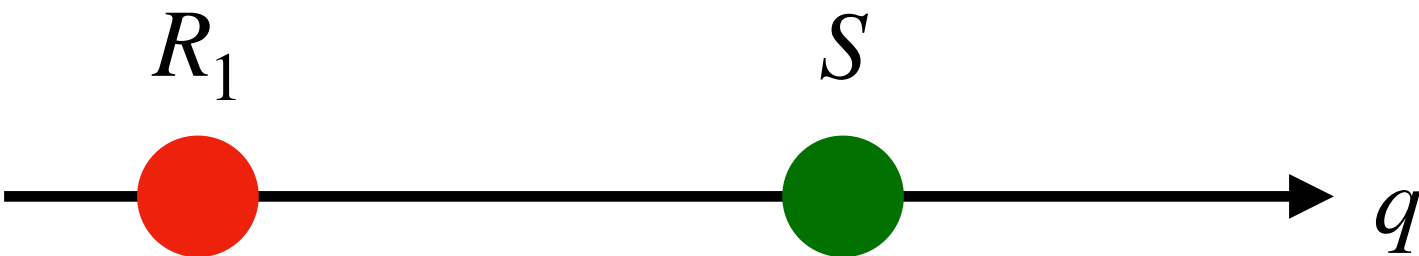
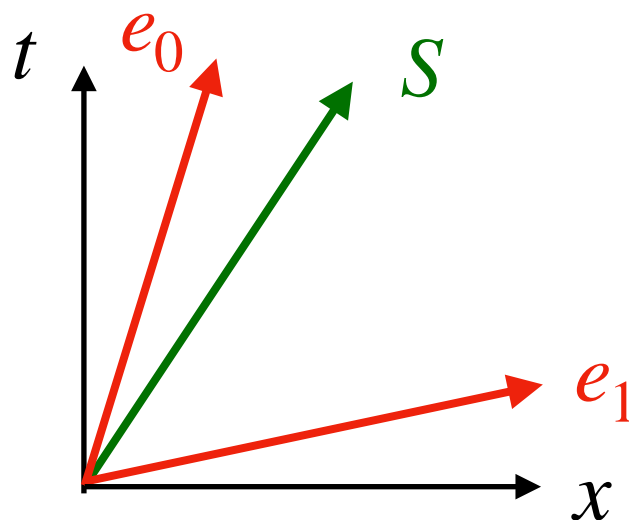
weak or strong indep. (dep. on approach)

$$[\rho_{\text{inv}}, U_{RS}^g] = 0 \quad \text{or} \quad U_{RS}^g |\psi\rangle = |\psi\rangle$$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g')

internal frames in SR vs. QT



weak or strong indep. (dep. on approach)

$$[\rho_{\text{inv}}, U_{RS}^g] = 0 \quad \text{or} \quad U_{RS}^g |\psi\rangle = |\psi\rangle$$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g')

internally
distinguishable
configurations
(extern frame indep.)

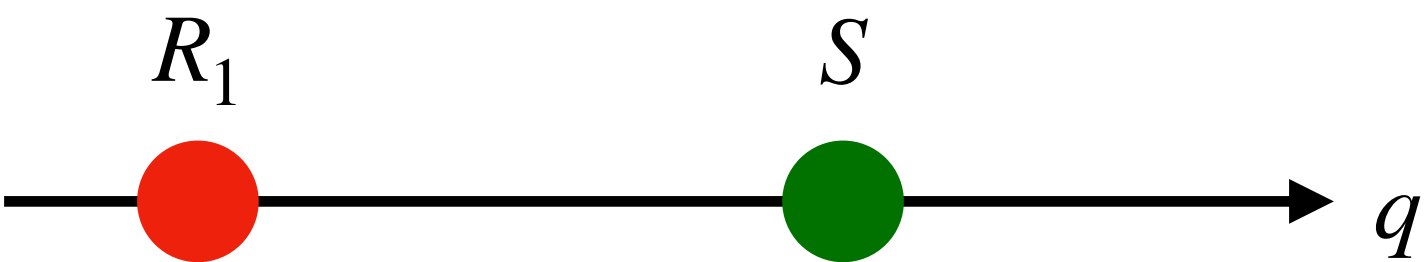
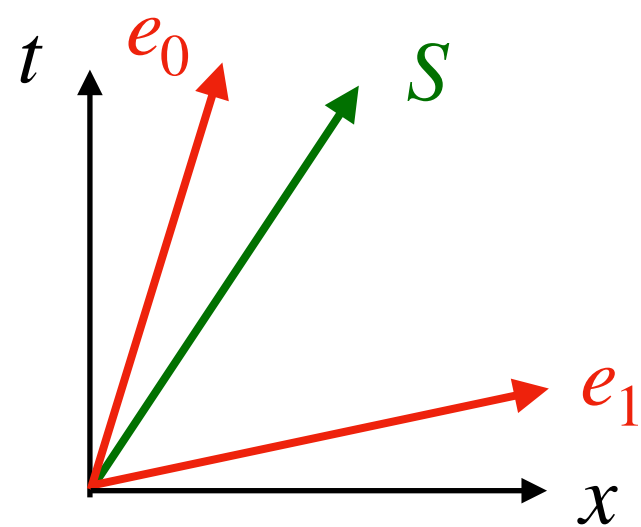
$$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$$

relational/dressed observables

a) relational observables

“jumping into internal
frame”

internal frames in SR vs. QT



weak or strong indep. (dep. on approach)

$$[\rho_{\text{inv}}, U_{RS}^g] = 0 \quad \text{or} \quad U_{RS}^g |\psi\rangle = |\psi\rangle$$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g')

internally distinguishable configurations
(extern frame indep.)

$$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$$

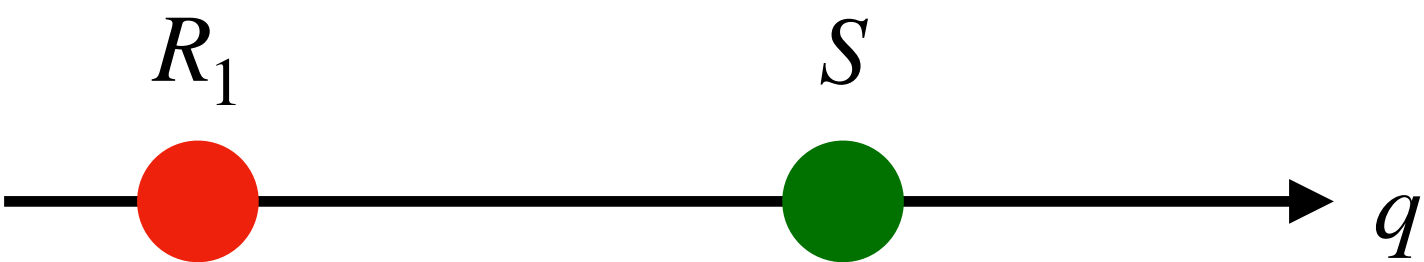
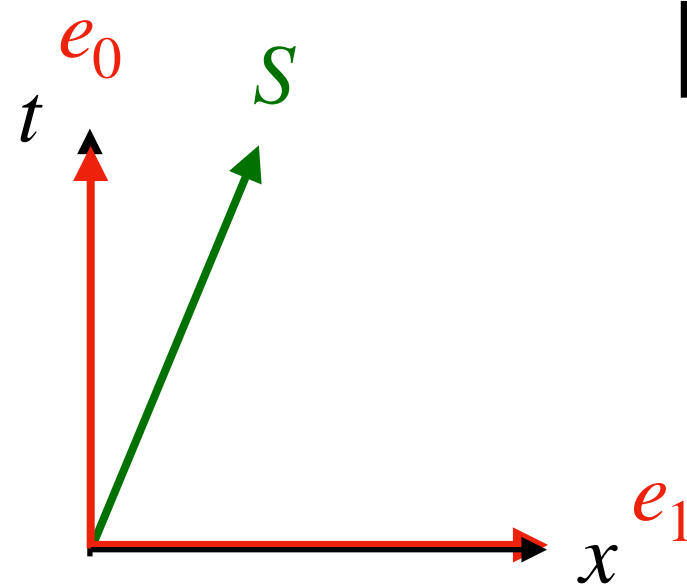
relational/dressed observables

“jumping into internal frame”

- a) relational observables**
- b) align int. & ext. frames (gauge fix)**

e.g. $e_a^\mu = \delta_a^\mu$

internal frames in SR vs. QT



**internally
distinguishable
configurations**
(extern frame indep.)

$$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$$

relational/dressed observables

weak or strong indep. (dep. on approach)

$$[\rho_{\text{inv}}, U_{RS}^g] = 0 \quad \text{or} \quad U_{RS}^g |\psi\rangle = |\psi\rangle$$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g')

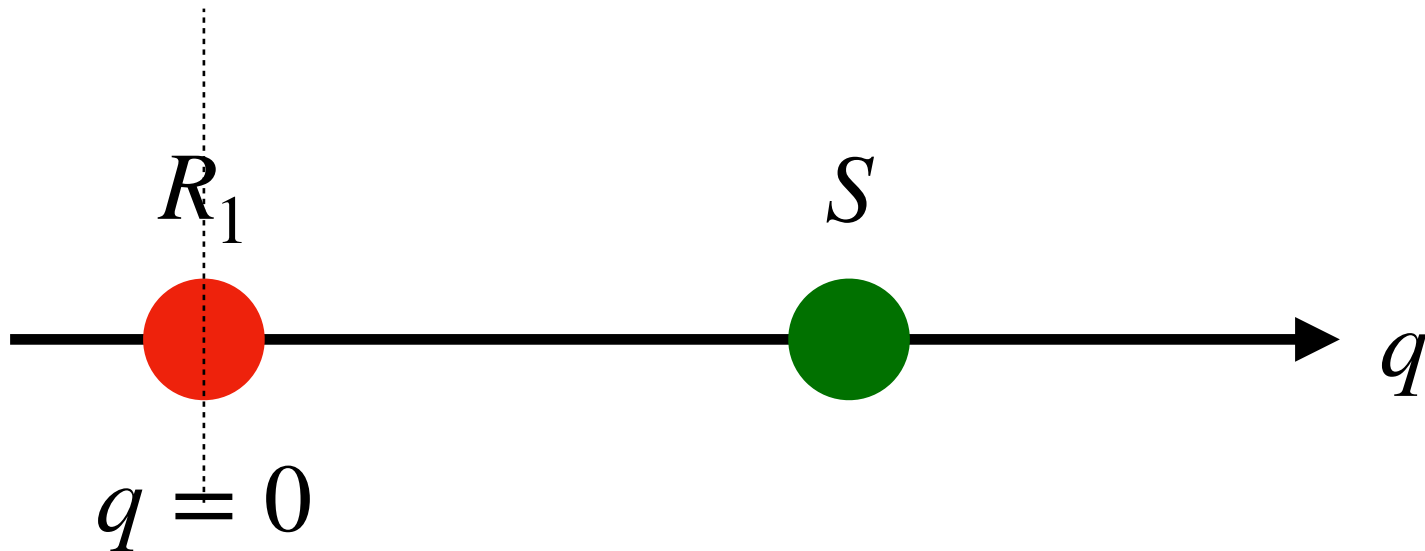
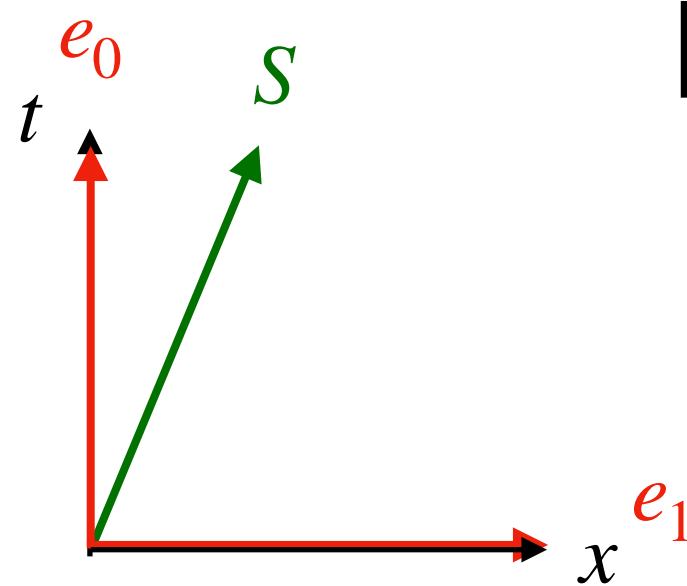
**“jumping into internal
frame”**

a) relational observables

b) align int. & ext. frames (gauge fix)

e.g. $e_a^\mu = \delta_a^\mu$

internal frames in SR vs. QT



**internally
distinguishable
configurations**
(extern frame indep.)

$$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$$

relational/dressed observables

weak or strong indep. (dep. on approach)

$$[\rho_{\text{inv}}, U_{RS}^g] = 0 \quad \text{or} \quad U_{RS}^g |\psi\rangle = |\psi\rangle$$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g)

**“jumping into internal
frame”**

a) relational observables

b) align int. & ext. frames (gauge fix)

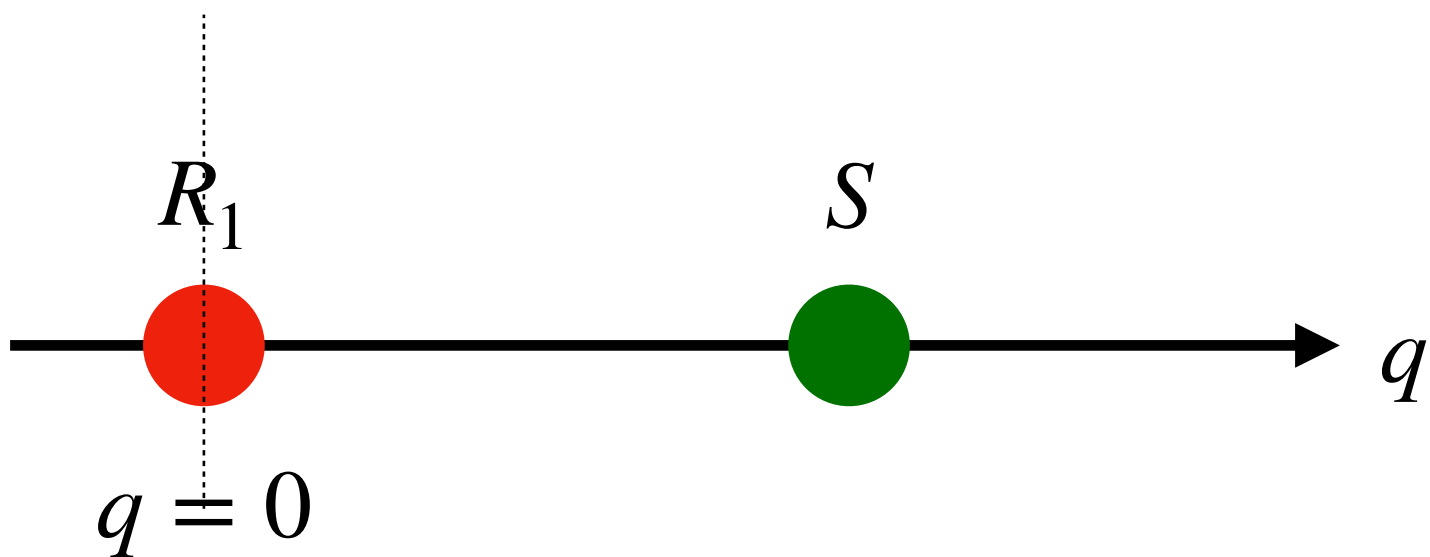
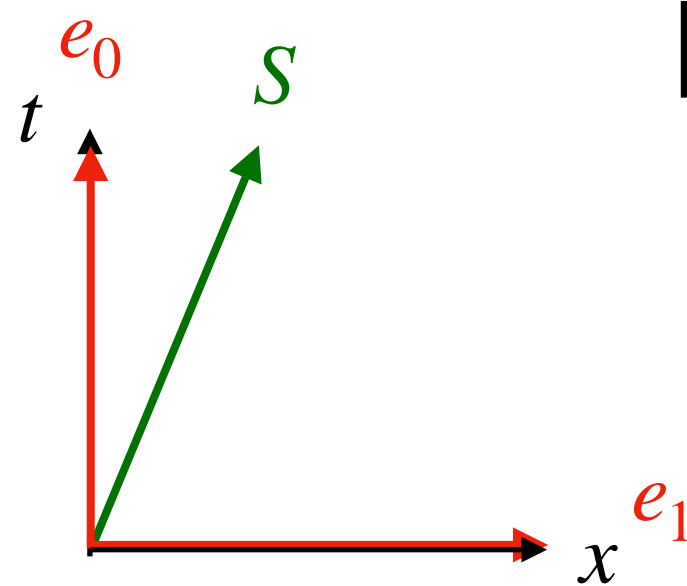
e.g. $e_a^\mu = \delta_a^\mu$

a) relational observables

b) align int. & ext. frames (gauge fix)

e.g. Page-Wootters reduction $\mathcal{R}^g = \langle g = e |_R \otimes I_S$

internal frames in SR vs. QT



weak or strong indep. (dep. on approach)

$$[\rho_{\text{inv}}, U_{RS}^g] = 0 \quad \text{or} \quad U_{RS}^g |\psi\rangle = |\psi\rangle$$

relational observables $\int_G dg U_{RS}^g (|g'\rangle\langle g'| \otimes f_S) (U_{RS}^g)^\dagger$

(f_S when QRF in orientation g)

internally distinguishable configurations
(extern frame indep.)

$$\{S_\mu S^\mu, S_\mu e_a^\mu, e_a^\mu e_\mu^b\}$$

relational/dressed observables

b) align int. & ext. frames (gauge fix)

e.g. Page-Wootters reduction $\mathcal{R}^g = \langle g = e |_R \otimes I_S$

can only be done unitarily charge sector wise!

frame perspective map of the perspective-neutral approach (strong inv.)

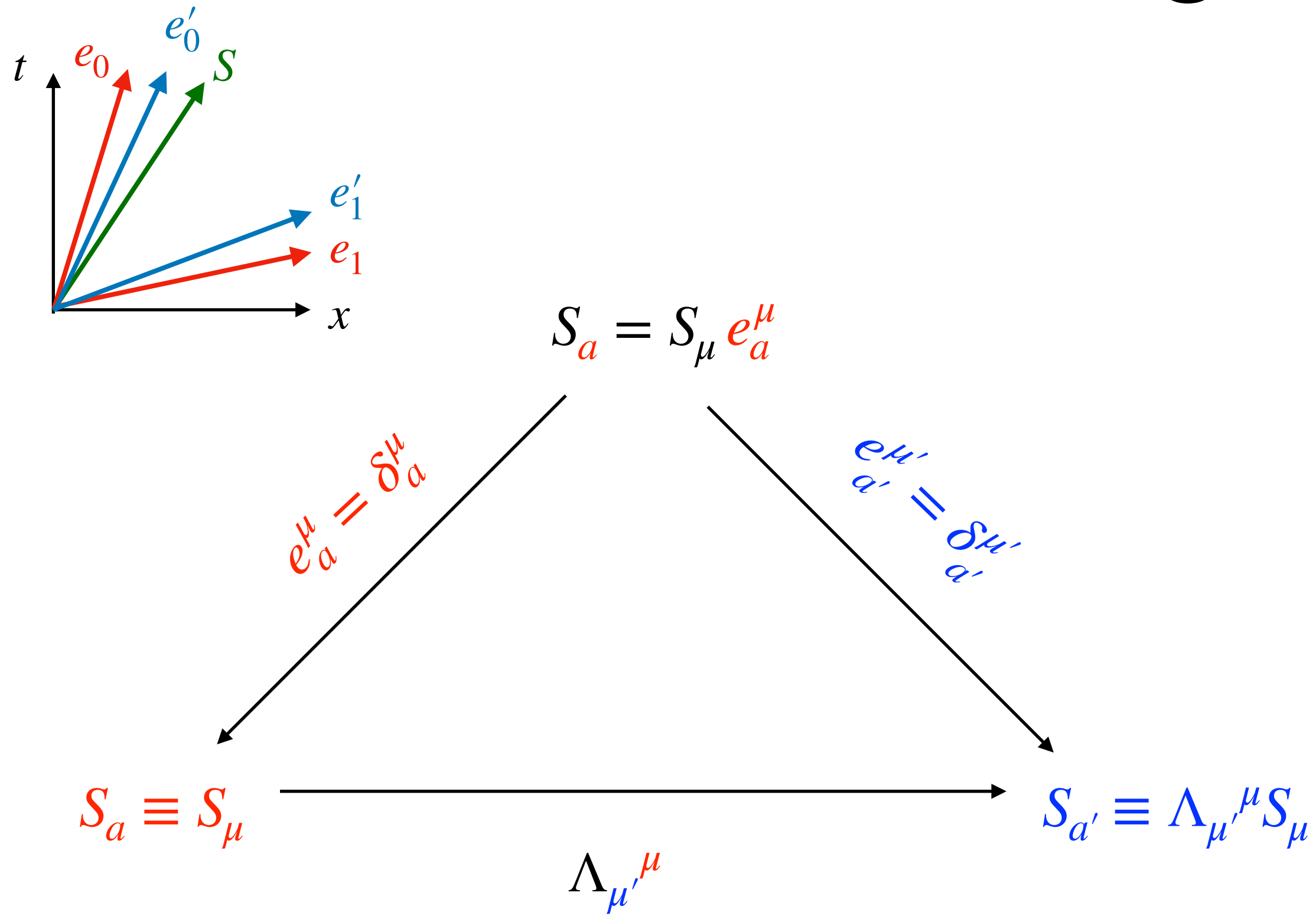
“jumping into internal frame”

a) relational observables

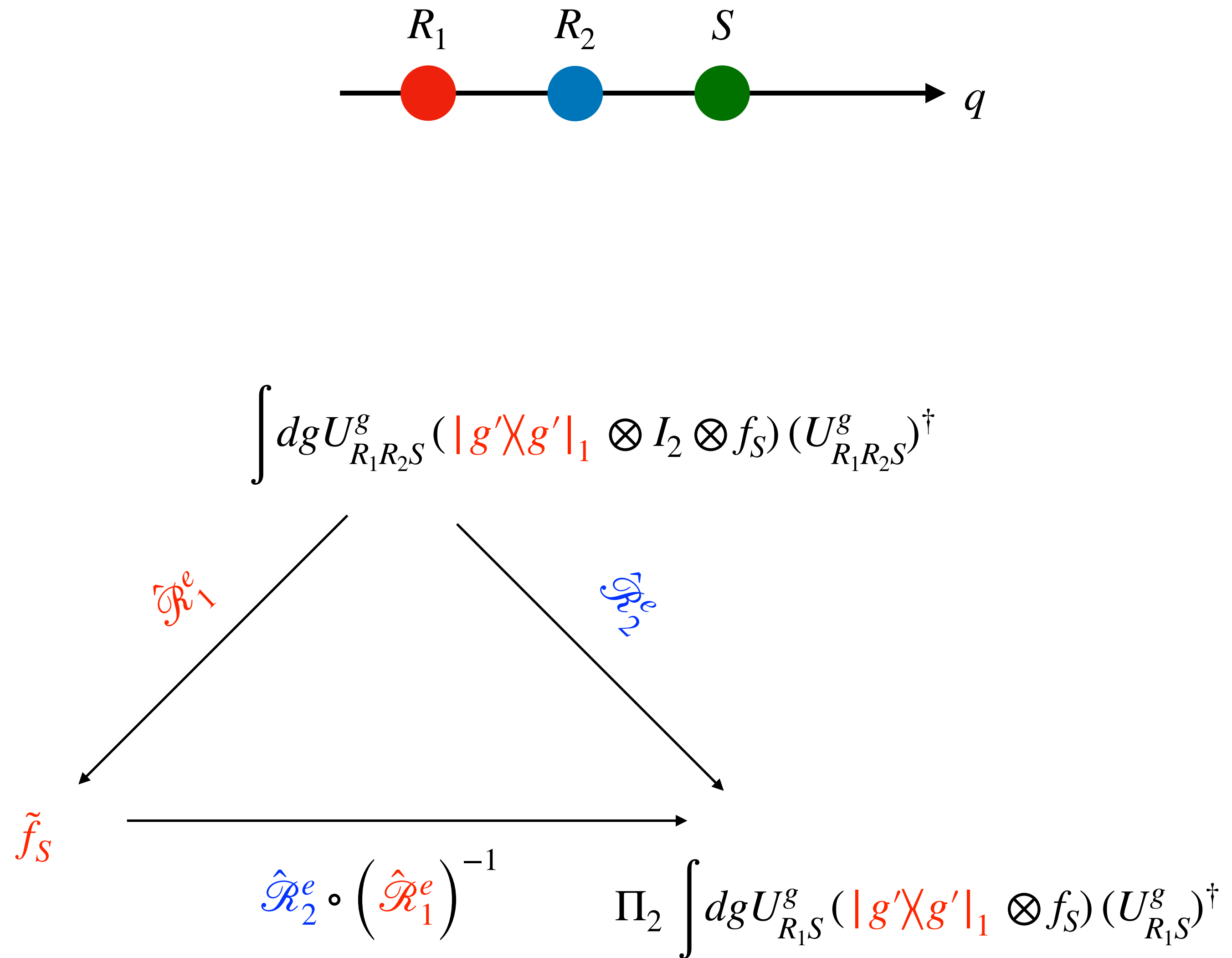
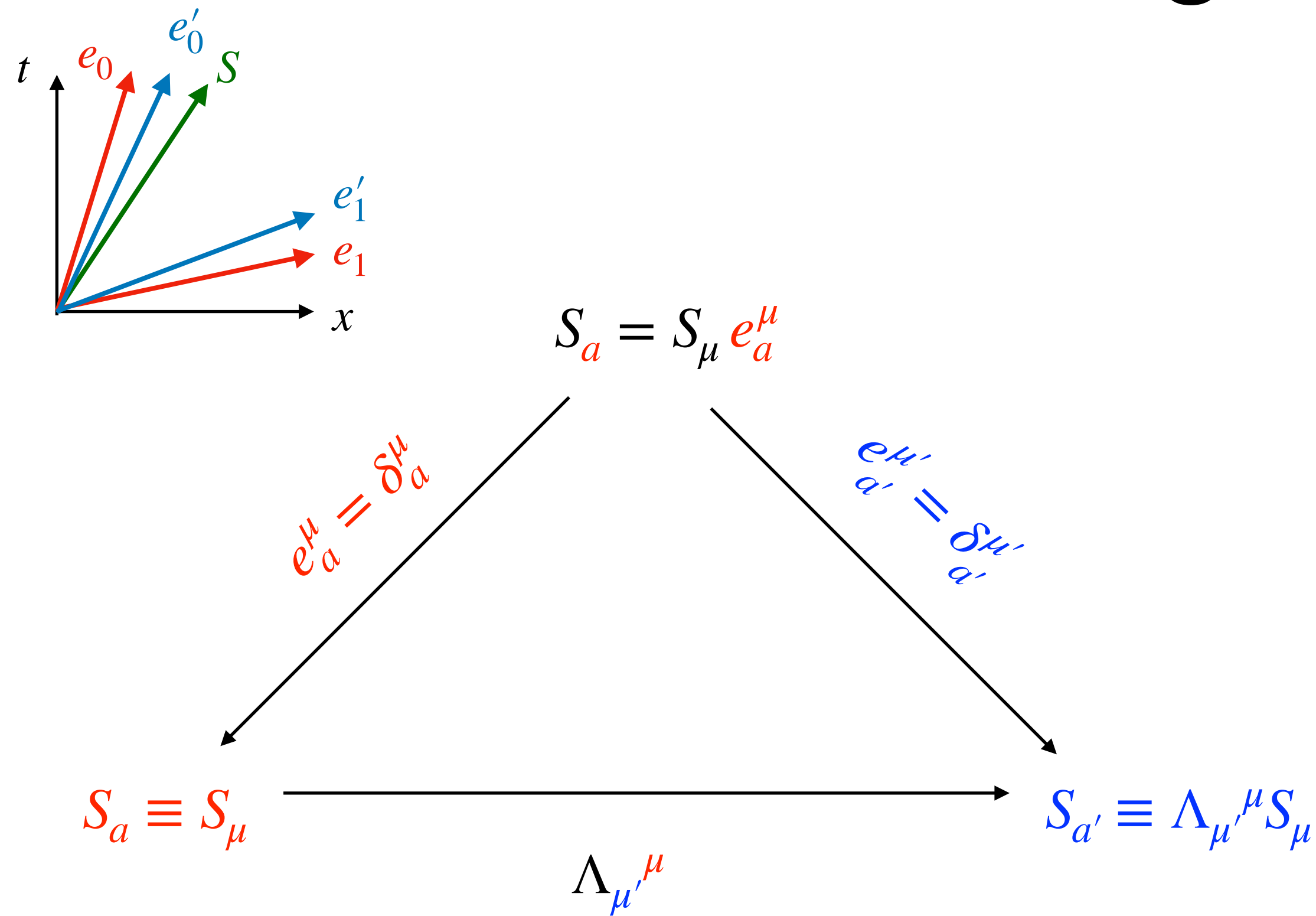
b) align int. & ext. frames (gauge fix)

e.g. $e_a^\mu = \delta_a^\mu$

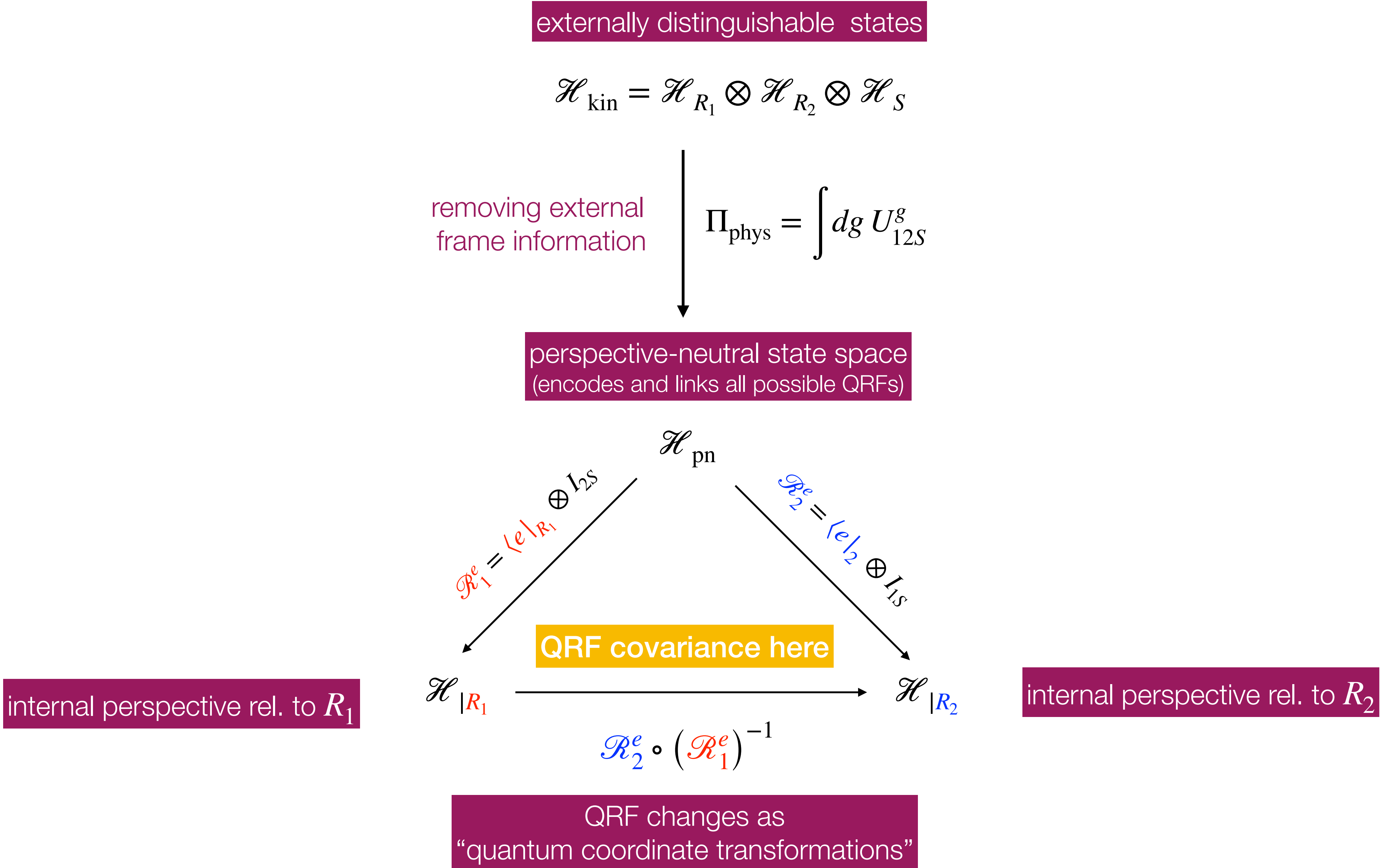
Internal frame changes in SR and the PN approach



Internal frame changes in SR and the PN approach



Summary: perspective-neutral QRF framework



Subsystem relativity

[Ahmad, Galley, PH, Lock, Smith '21; Araújo-Regado, PH, Sartini '25;
Castro-Ruiz, Oreshkov '21; PH, Kotecha, Mele '23; De Vuyst, Eccles, PH, Kirklin '24;]

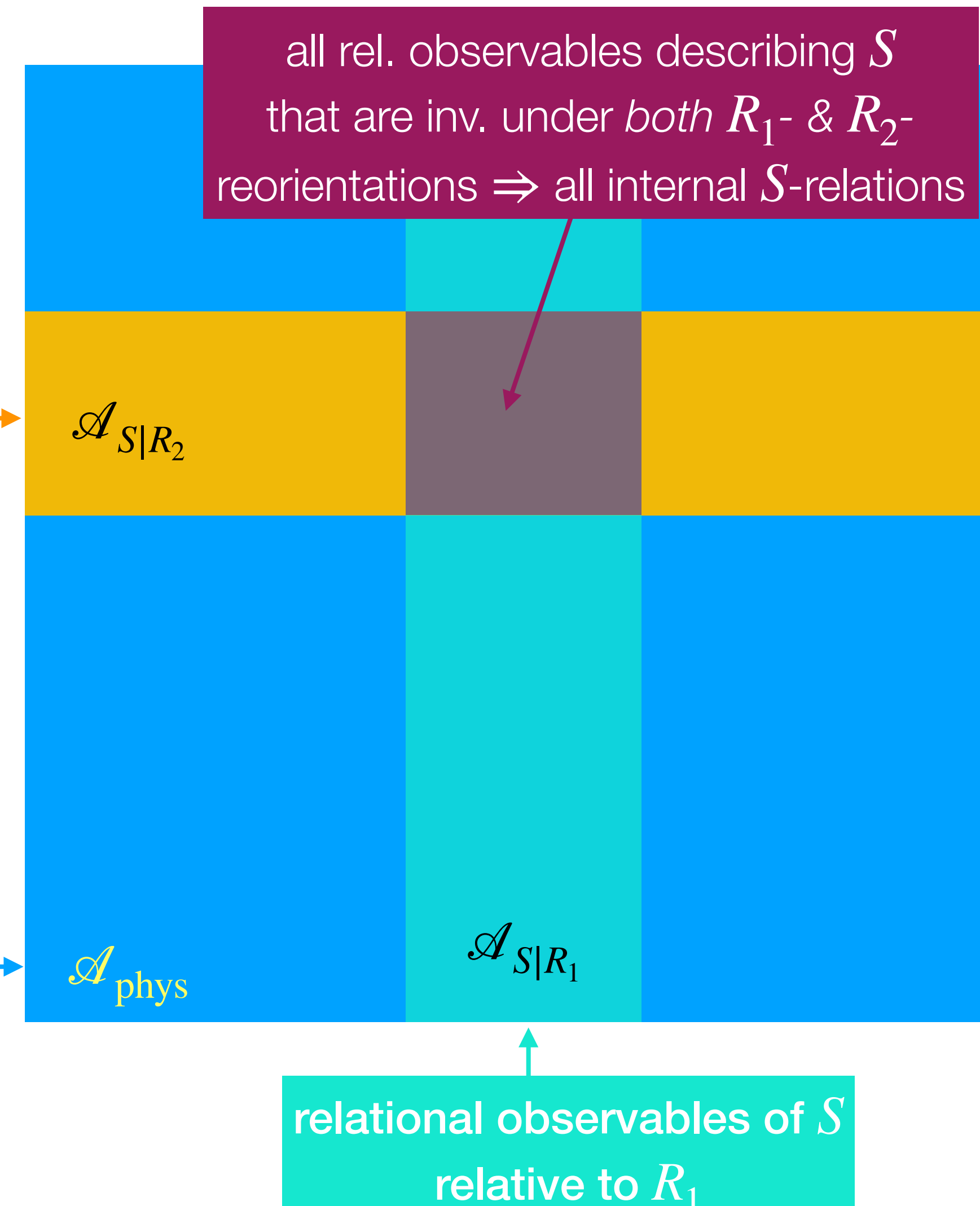
Relativity of subsystems

$$\mathcal{A}_{S|R_1} \cap \mathcal{A}_{S|R_2} = \mathcal{A}_S^G = \{\text{internal rel. obs. of } S\}$$

“frames R_1 and R_2 mean different gauge inv. DoFs when they refer to subsystem S ”

relational observables of
 S relative to R_2

full algebra of
relational observables



Relativity of subsystems

$$\mathcal{A}_{S|R_1} \cap \mathcal{A}_{S|R_2} = \mathcal{A}_S^G = \{\text{internal rel. obs. of } S\}$$

relational observables of
 S relative to R_2

$\mathcal{A}_{S|R_2}$

all rel. observables describing S
that are inv. under *both* R_1 - & R_2 -
reorientations $\Rightarrow$ all internal S -relations

“frames R_1 and R_2 mean different gauge inv. DoFs when they refer to subsystem S ”

full algebra of
relational observables

$\mathcal{A}_{\text{phys}}$

$\mathcal{A}_{S|R_1}$

relational observables of S
relative to R_1

$\Rightarrow$ correlations, entropies, thermodynamic properties, ...
become QRF-dependent

[systematic investigation in
PH, Kotecha, Mele '23]

QRF perspectives are TPSs on the PN-space

Kotecha, Mele, PH '23;
De Vuyst, Eccles, PH, Kirklin '24;
Araújo-Regado, PH, Sartini '25

TPS relative to external frame

$$\mathcal{H}_{\text{kin}} = \mathcal{H}_{R_1} \otimes \mathcal{H}_{R_2} \otimes \mathcal{H}_S$$

removing external
frame information

Π_{phys}

perspective/TPS-neutral state space
(does not inherit kin. TPS)

ideal (sharp) QRFs ($\langle g | g' \rangle = \delta(g, g')$)

**QRF perspective is a
gauge-inv. subsystem partition**

$\mathcal{H}_{\text{pn}}$

$\mathcal{R}_1^e$

$\mathcal{R}_2^e$

internal TPS rel. to R_1

$$\mathcal{H}_{|R_1} = \mathcal{H}_{R_2} \otimes \mathcal{H}_S$$

$$\mathcal{H}_{|R_2} = \mathcal{H}_{R_1} \otimes \mathcal{H}_S$$

internal TPS rel. to R_2

$$\mathcal{R}_2^e \circ (\mathcal{R}_1^e)^{-1}$$

QRF transformations are
changes of TPS

QRF perspectives are TPSs on the PN-space

Kotecha, Mele, PH '23;
De Vuyst, Eccles, PH, Kirklin '24;
Araújo-Regado, PH, Sartini '25

TPS relative to external frame

$$\mathcal{H}_{\text{kin}} = \mathcal{H}_{R_1} \otimes \mathcal{H}_{R_2} \otimes \mathcal{H}_S$$

removing external
frame information

Π_{phys}

perspective/TPS-neutral state space
(does not inherit kin. TPS)

ideal (sharp) QRFs ($\langle g | g' \rangle = \delta(g, g')$)

**QRF perspective is a
gauge-inv. subsystem partition**

generated by relational observable algebras

$$\mathcal{A}_{\text{pn}} \simeq \mathcal{A}_{S|R_1} \otimes \mathcal{A}_{R_2|R_1} \simeq \mathcal{A}_{S|R_2} \otimes \mathcal{A}_{R_1|R_2}$$

internal TPS rel. to R_1

$$\mathcal{H}_{|R_1} = \mathcal{H}_{R_2} \otimes \mathcal{H}_S$$

$$\mathcal{H}_{|R_2} = \mathcal{H}_{R_1} \otimes \mathcal{H}_S$$

internal TPS rel. to R_2

$$\mathcal{R}_2^e \circ (\mathcal{R}_1^e)^{-1}$$

QRF transformations are
changes of TPS

Example: particles on a line

Vanrietvelde, PH, Giacomini, Castro-Ruiz '19



perspective-neutral space

$\mathcal{H}_{\text{pn}}$

$\mathcal{R}_1^e = \langle q_1 = 0 \rangle \oplus I_{2s}$

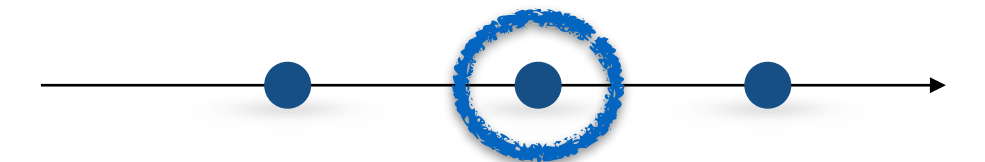
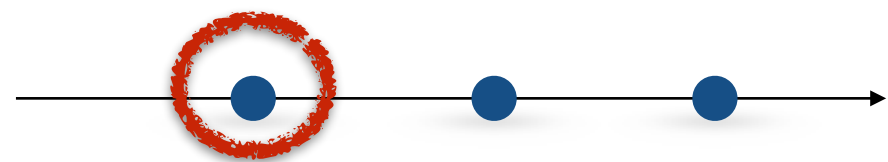
$\mathcal{R}_2^e = \langle q_2 = 0 \rangle \oplus I_{1s}$

internal TPS rel. to R_1

$$\mathcal{H}_{|R_1} = L^2(\mathbb{R})_{R_2} \otimes L^2(\mathbb{R})_S$$

$$\mathcal{H}_{|R_2} = L^2(\mathbb{R})_{R_1} \otimes L^2(\mathbb{R})_S$$

internal TPS rel. to R_2



Example: particles on a line

Vanrietvelde, PH, Giacomini, Castro-Ruiz '19



perspective-neutral space

$\mathcal{H}_{\text{pn}}$

$\mathcal{R}_1^e = \langle q_1 = 0 | \oplus I_{2S}$

$\mathcal{R}_2^e = \langle q_2 = 0 | \oplus I_{1S}$

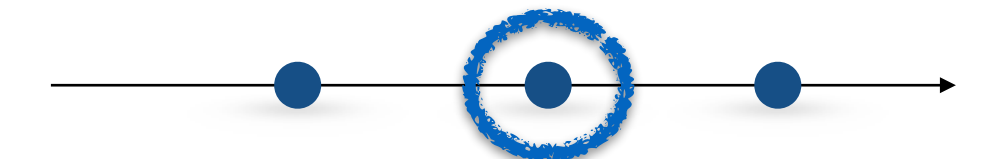
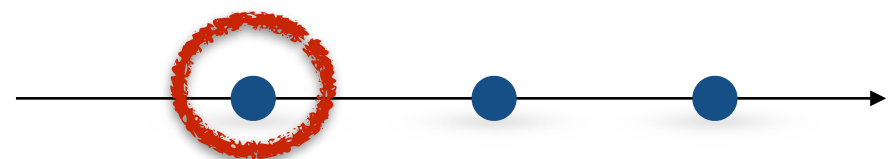
internal TPS rel. to R_1

$$\mathcal{H}_{|R_1} = L^2(\mathbb{R})_{R_2} \otimes L^2(\mathbb{R})_S$$

$$\mathcal{H}_{|R_2} = L^2(\mathbb{R})_{R_1} \otimes L^2(\mathbb{R})_S$$

internal TPS rel. to R_2

$$\mathcal{R}_2^e \circ (\mathcal{R}_1^e)^{-1} = \mathbb{F}_{12} \int dq | -q \rangle \langle q |_{R_2} \otimes U_S(q)$$



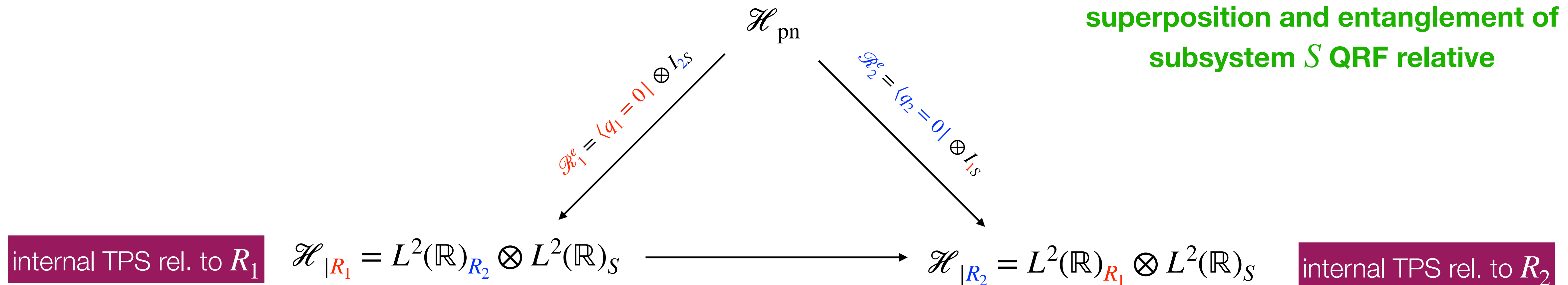
Example: particles on a line

Vanrietvelde, PH, Giacomini, Castro-Ruiz '19

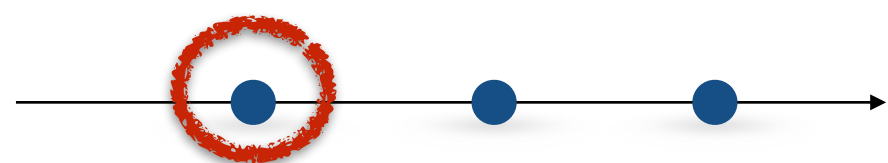


perspective-neutral space

superposition and entanglement of
subsystem S QRF relative

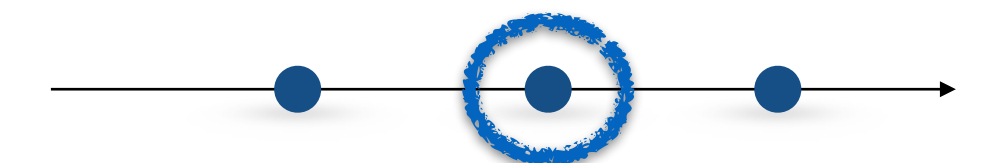


$$\mathcal{R}_2^e \circ (\mathcal{R}_1^e)^{-1} = \mathbb{F}_{12} \int dq | -q \rangle \langle q |_{R_2} \otimes U_S(q)$$



$$(|q_1\rangle_{R_2} + |q_2\rangle_{R_2}) \otimes |x\rangle_S$$

superposition + product



$$| -q_1 \rangle_{R_1} \otimes |x - q_1\rangle_S + | -q_2 \rangle_{R_1} \otimes |x - q_2\rangle_S$$

entanglement

QRF perspectives are TPSs on the PN-space

Kotecha, Mele, PH '23;
De Vuyst, Eccles, PH, Kirklin '24;
Araújo-Regado, PH, Sartini '25

TPS relative to external frame

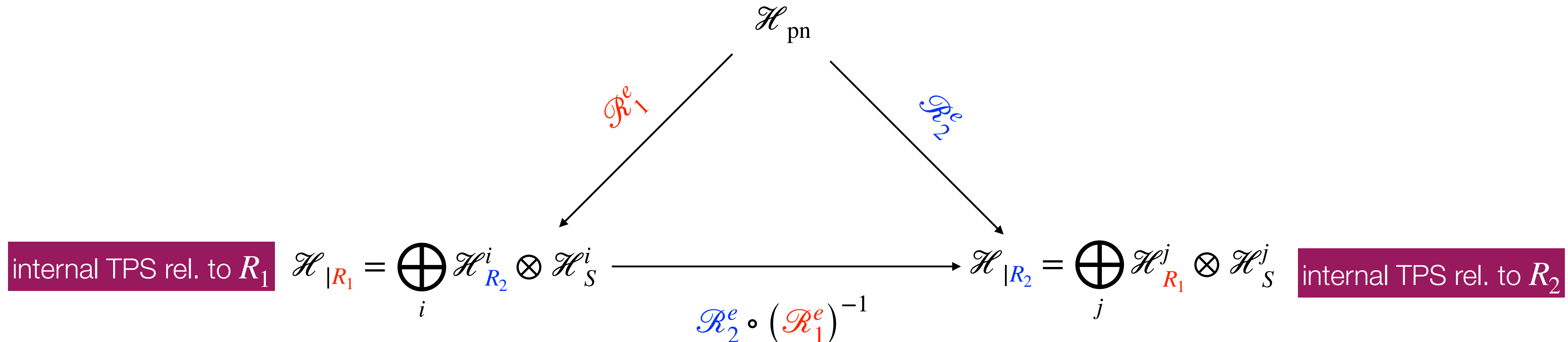
$$\mathcal{H}_{\text{kin}} = \mathcal{H}_{R_1} \otimes \mathcal{H}_{R_2} \otimes \mathcal{H}_S$$

removing external
frame information

Π_{phys}

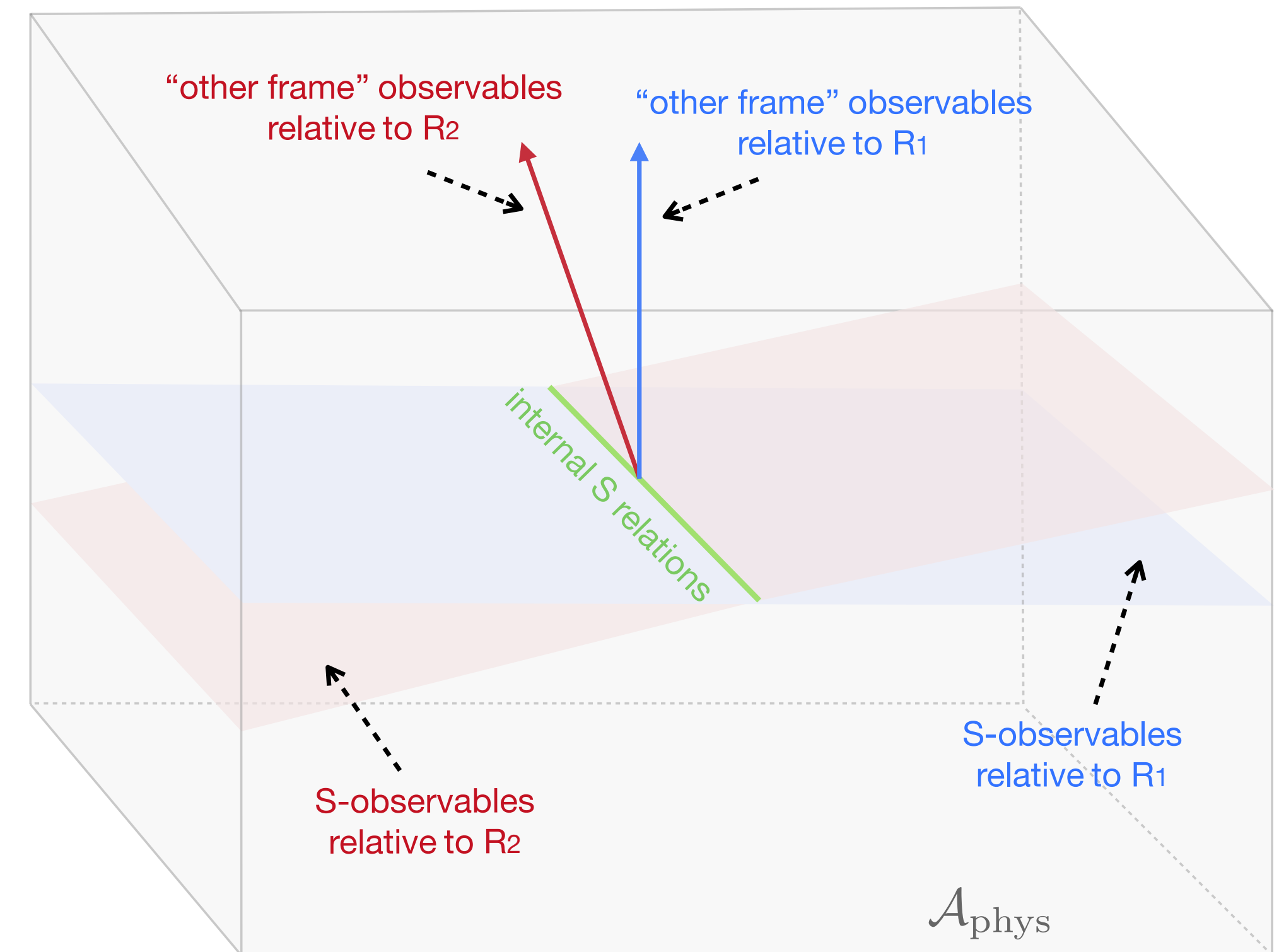
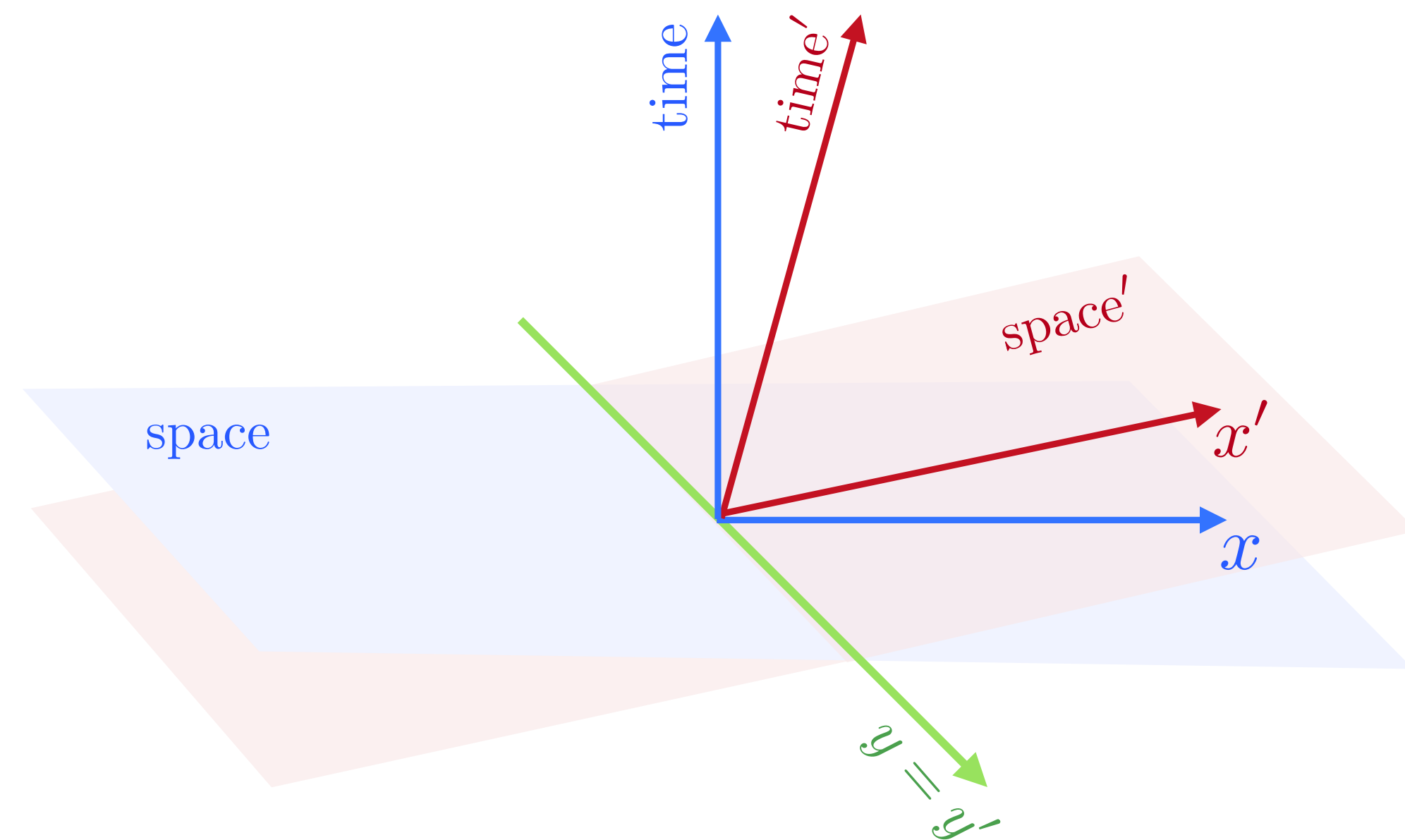
perspective/TPS-neutral state space
(does not inherit kin. TPS)

fuzzy QRFs



QRF transformations are
changes of TPS

Relativity of simultaneity vs. subsystem relativity



⇒ physical consequences should be seen in similar light:
Just like the relativity of simultaneity is the root of all characteristic special relativistic phenomena,
so is subsystem relativity root of all novel QRF relative effects

Gauge and physical QRFs in lattice gauge theory

[Araújo-Regado, PH, Sartini 2506.23459]

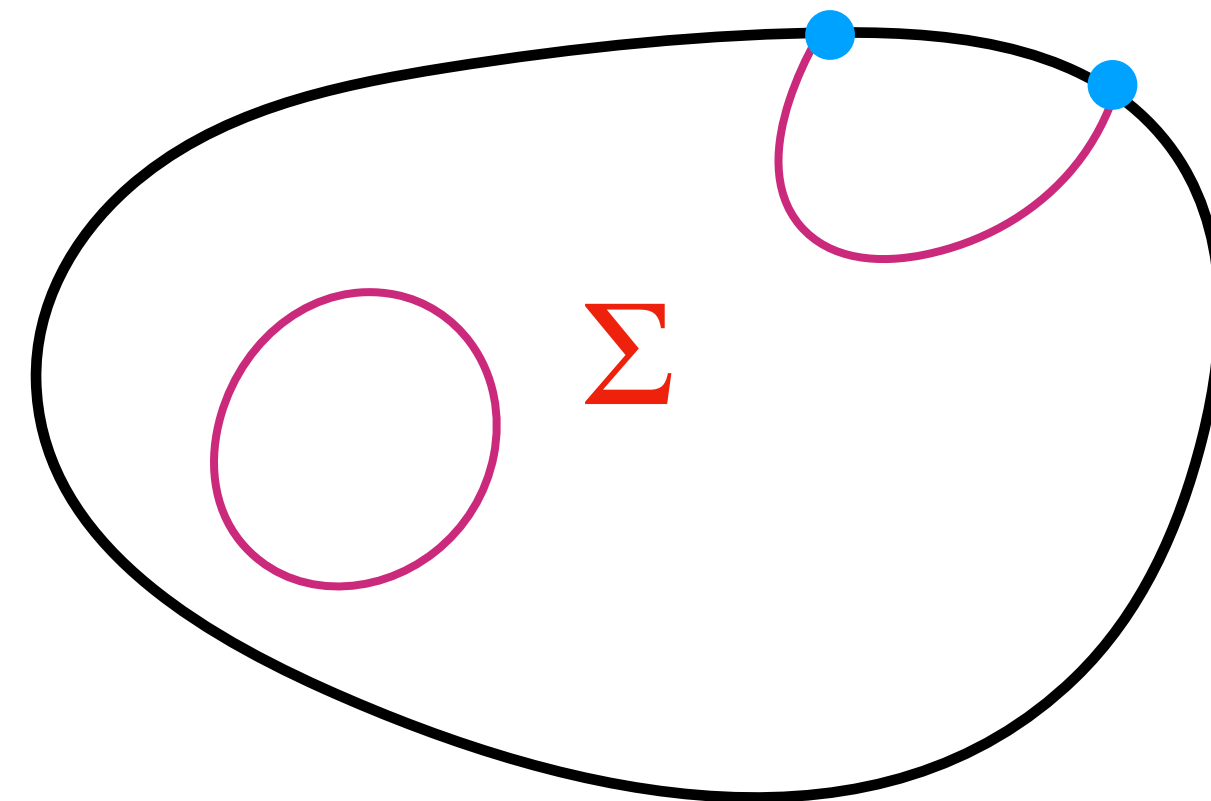
Edge modes and corner symmetries

small gauge transf. generated by Gauß law (redundancy)

$$C[\alpha] = \int_{\Sigma} \text{Tr}(\alpha d_A \star F)$$

electric corner symmetries (physical)

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_{\Phi})$$



edge mode Φ

[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

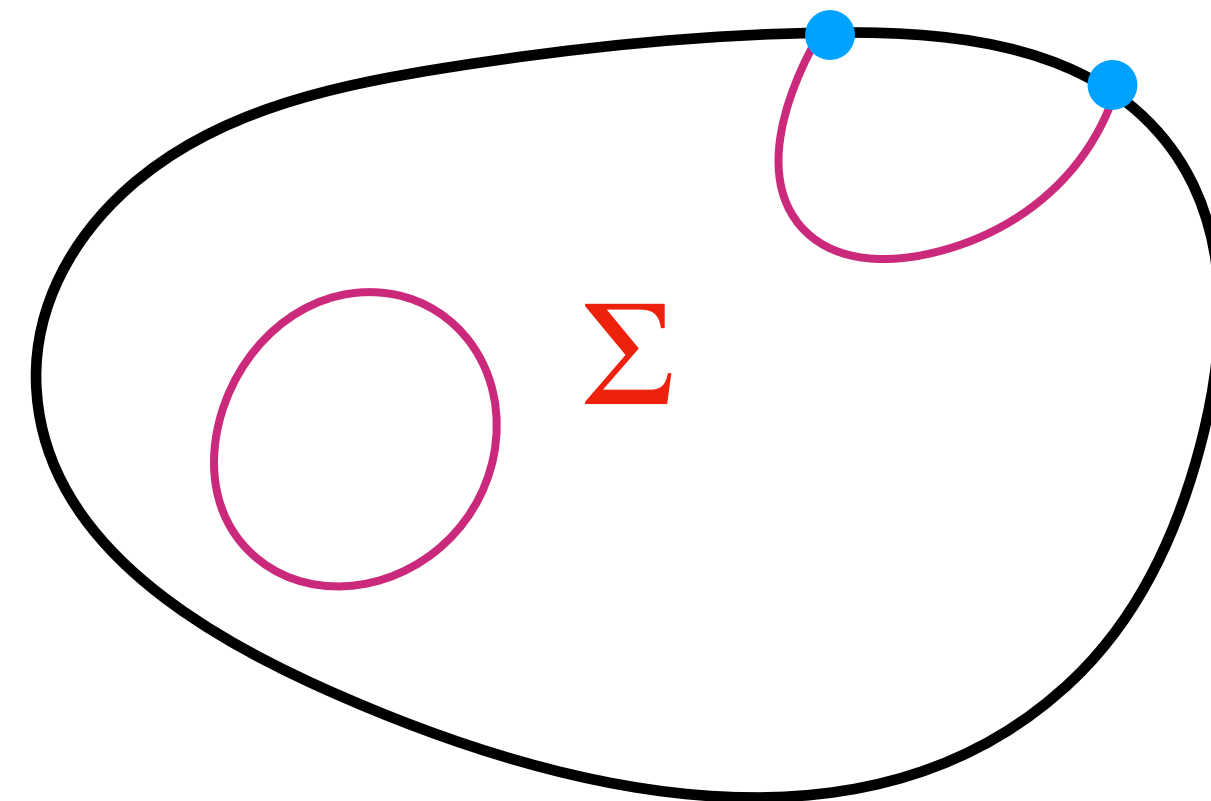
Edge modes and corner symmetries

small gauge transf. generated by Gauß law (redundancy)

$$C[\alpha] = \int_{\Sigma} \text{Tr}(\alpha d_A \star F)$$

electric corner symmetries (physical)

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_{\Phi})$$



edge mode Φ

[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

$\Rightarrow$ how do you realize them?

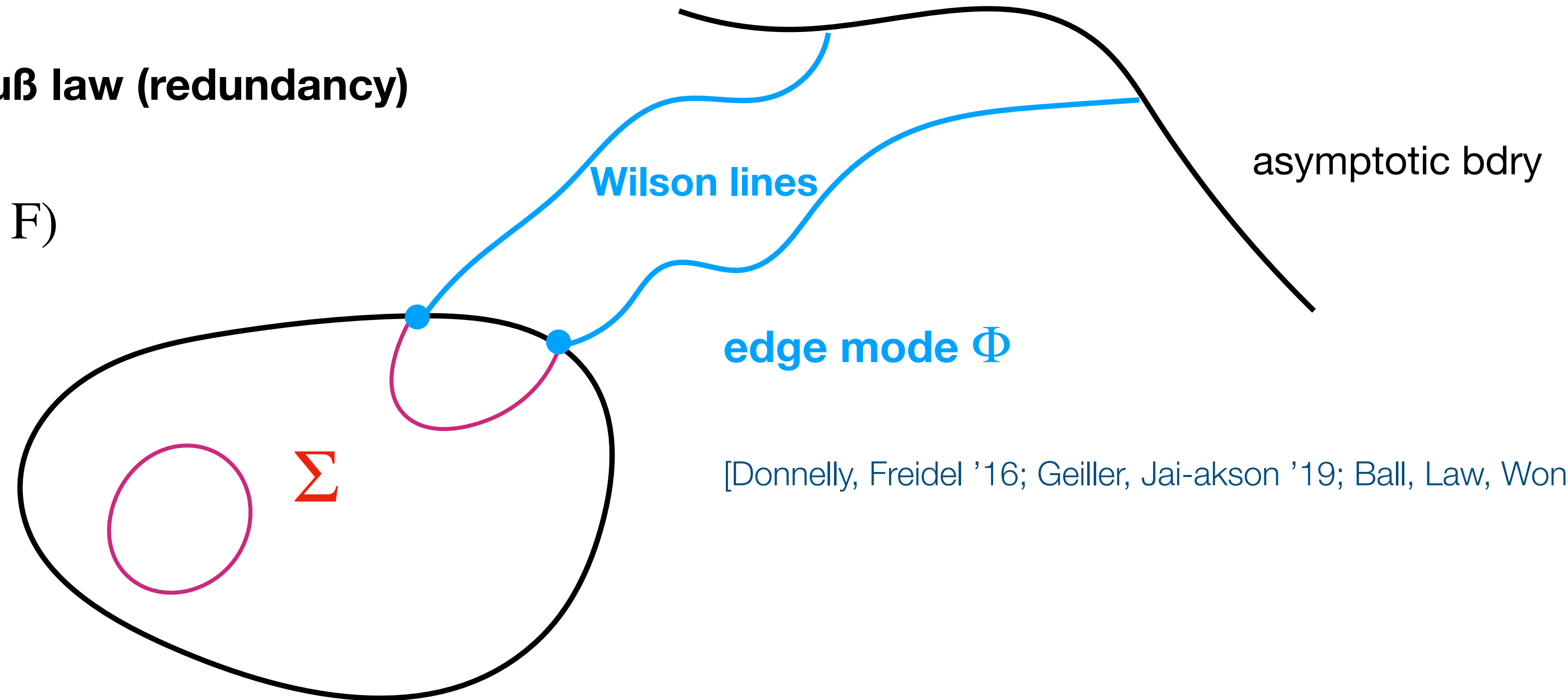
Edge modes as QRFs

small gauge transf. generated by Gauß law (redundancy)

$$C[\alpha] = \int_{\Sigma} \text{Tr}(\alpha d_A \star F)$$

electric corner symmetries (physical)

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_{\Phi})$$



[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

⇒ how do you realize them? as QRFs

extrinsic edge mode QRFs:

e.g. via extrinsic Wilson lines

[Carrozza, PH '21; Araújo-Regado, PH, Sartini, Tomova '24]

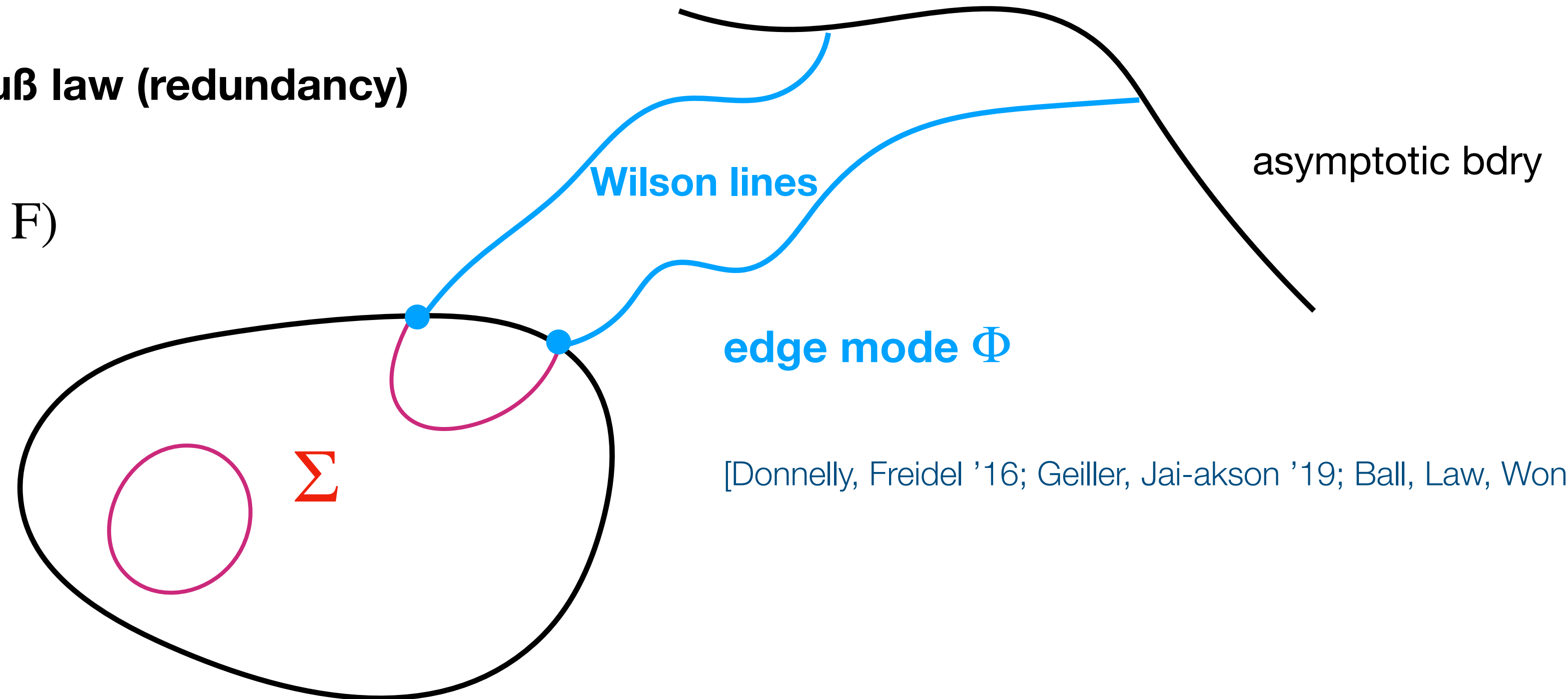
Edge modes as QRFs

small gauge transf. generated by Gauß law (redundancy)

$$C[\alpha] = \int_{\Sigma} \text{Tr}(\alpha d_A \star F)$$

electric corner symmetries (physical)

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_{\Phi})$$



[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

⇒ how do you realize them? as QRFs

extrinsic edge mode QRFs:

e.g. via extrinsic Wilson lines

[Carrozza, PH '21; Araújo-Regado, PH, Sartini, Tomova '24]

intrinsic edge mode QRFs:

e.g. via Wilson lines anchored on the corner $\partial\Sigma$ or a Hodge decomp.

[Araújo-Regado, PH, Sartini, Tomova '24]

Edge modes as QRFs

small gauge transf. generated by Gauß law (redundancy)

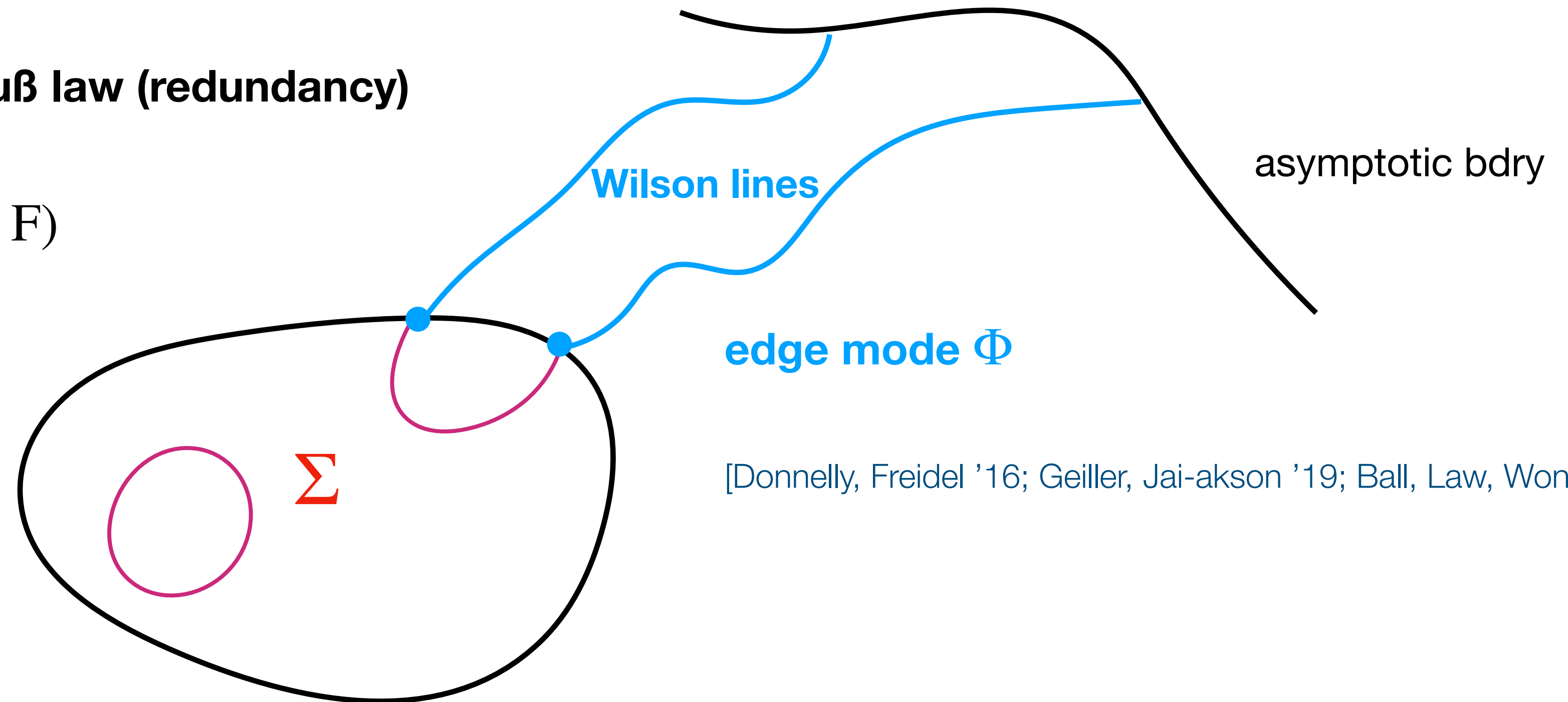
$$C[\alpha] = \int_{\Sigma} \text{Tr}(\alpha d_A \star F)$$

electric corner symmetries (physical)

$$Q[\rho] = \int_{\partial\Sigma} \text{Tr}(\rho \star F_{\Phi})$$

⇒ are reorientations of
extrinsic gauge QRFs!

⇒ how do you realize them? as gauge QRFs



[Donnelly, Freidel '16; Geiller, Jai-akson '19; Ball, Law, Wong '24;...]

extrinsic edge mode QRFs:

e.g. via extrinsic Wilson lines

[Carrozza, PH '21; Araújo-Regado, PH, Sartini, Tomova '24]

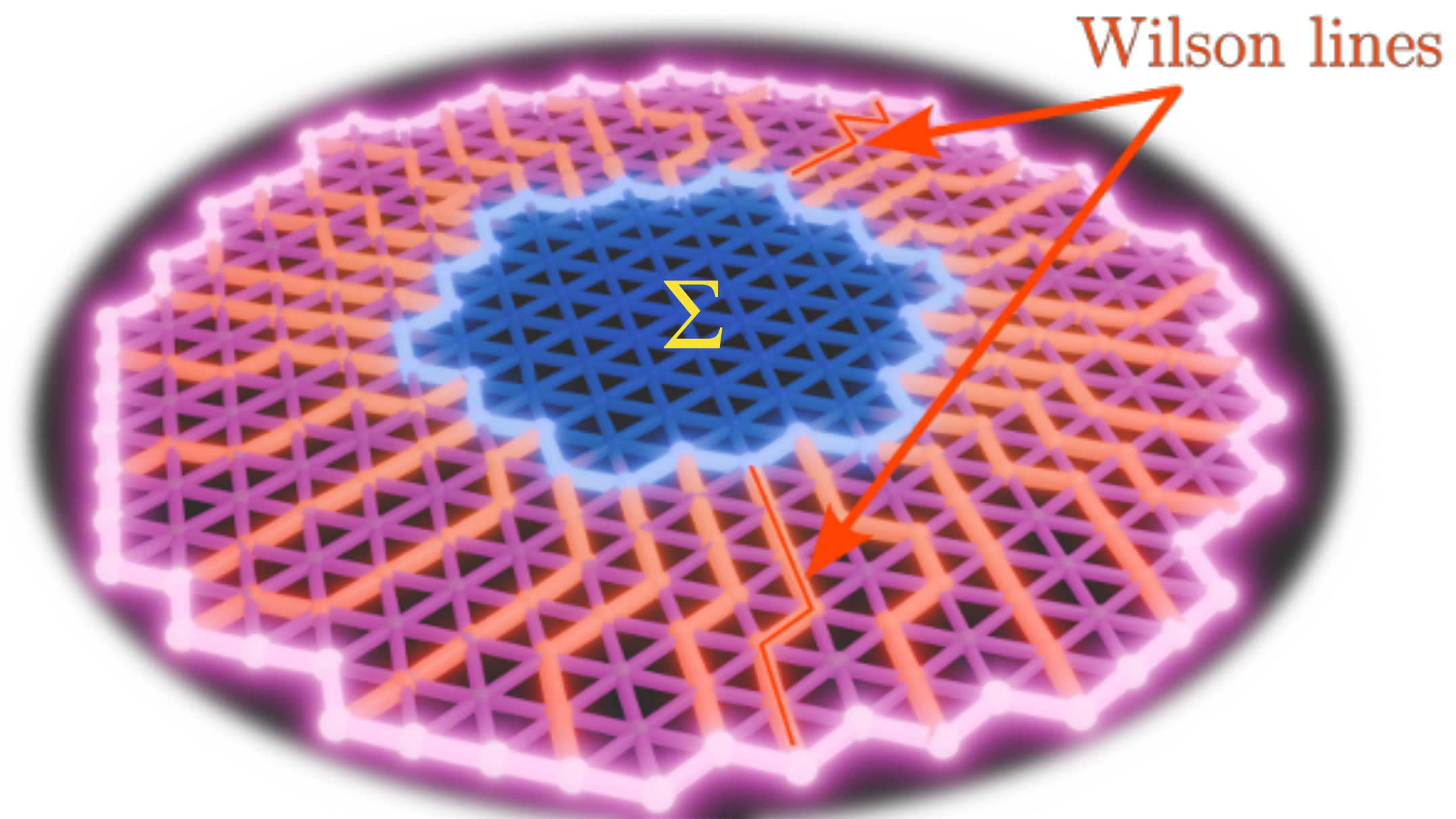
intrinsic edge mode QRFs:

e.g. via Wilson lines anchored on the corner $\partial\Sigma$ or a Hodge decomp.

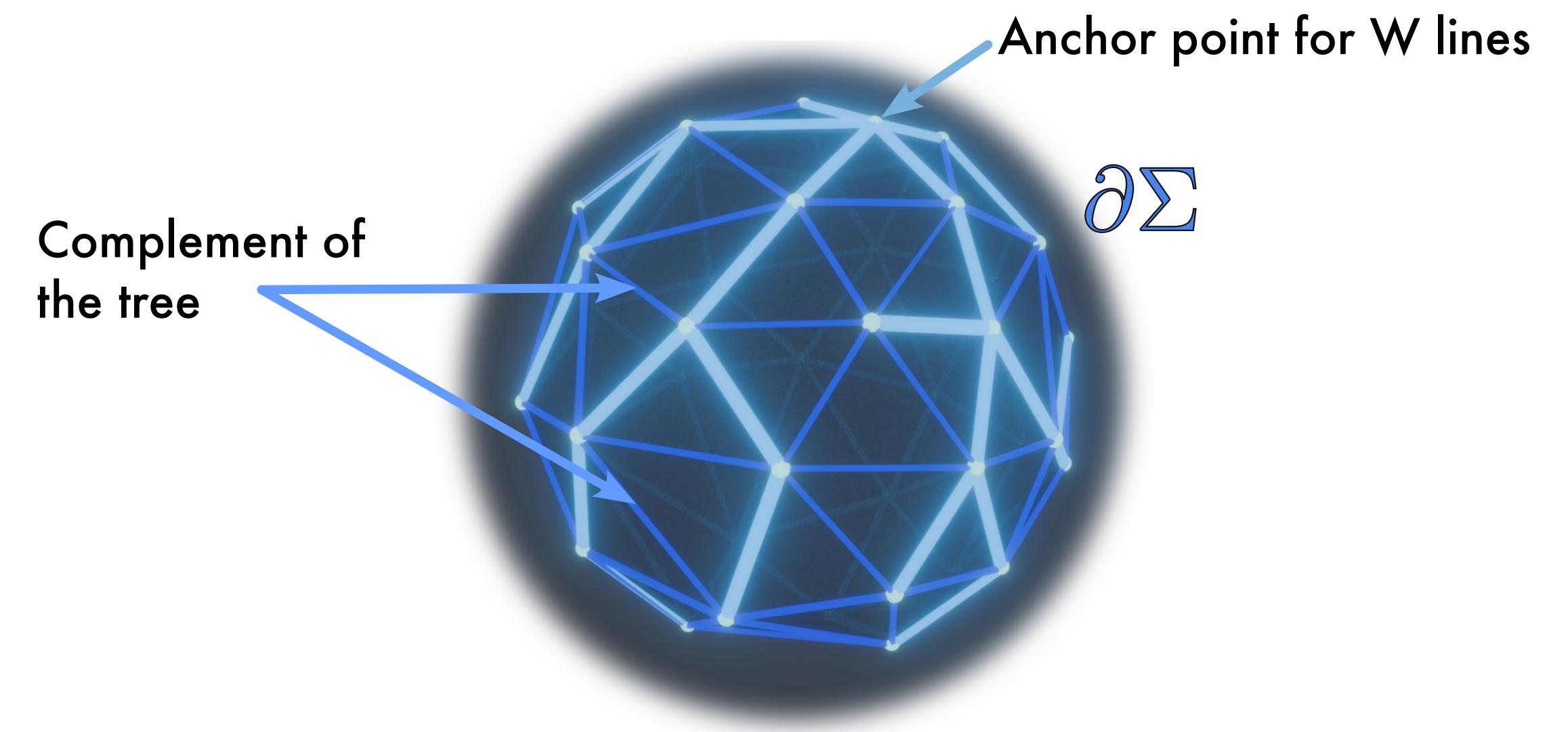
[Araújo-Regado, PH, Sartini, Tomova '24]

Extrinsic and intrinsic gauge QRFs on the lattice

Extrinsic: congruence of Wilson lines in the complement

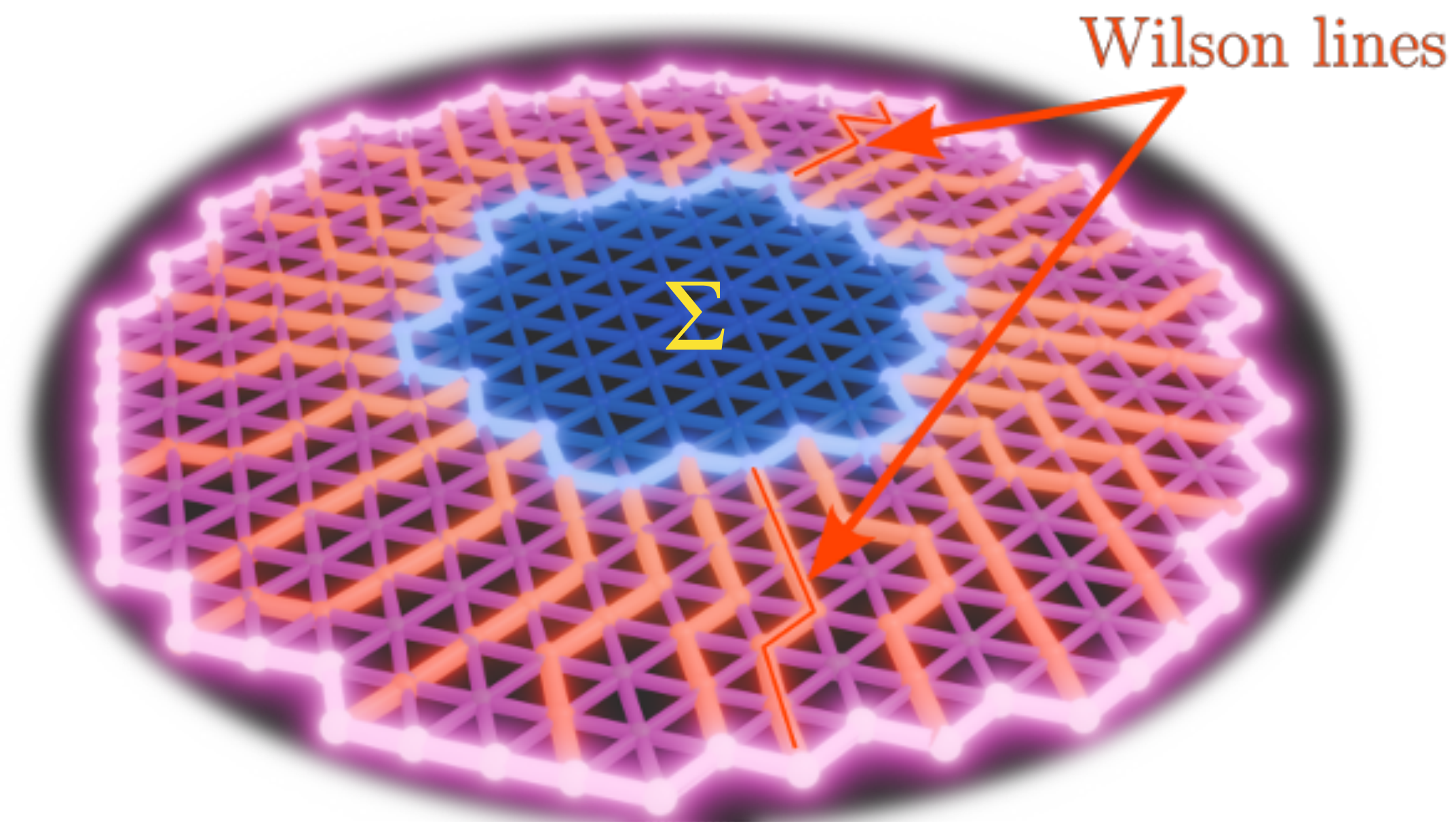


Spanning tree anchored on node N of the corner



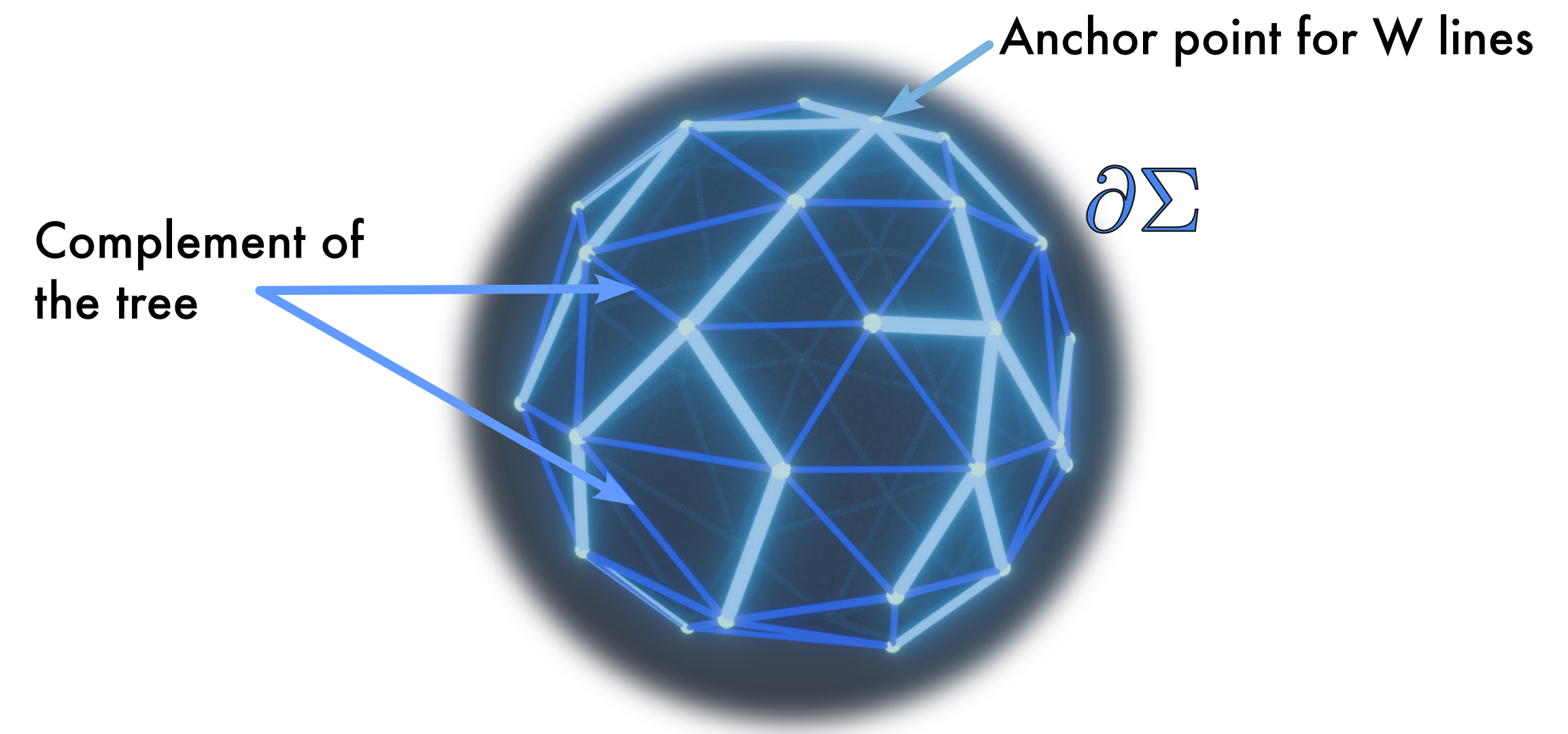
Extrinsic and intrinsic gauge QRFs on the lattice

Extrinsic: congruence of Wilson lines in the complement



complete ideal QRF: $\mathcal{H}_{\Phi} = L^2(G^{V_{\partial\Sigma}})$

Spanning tree anchored on node N of the corner

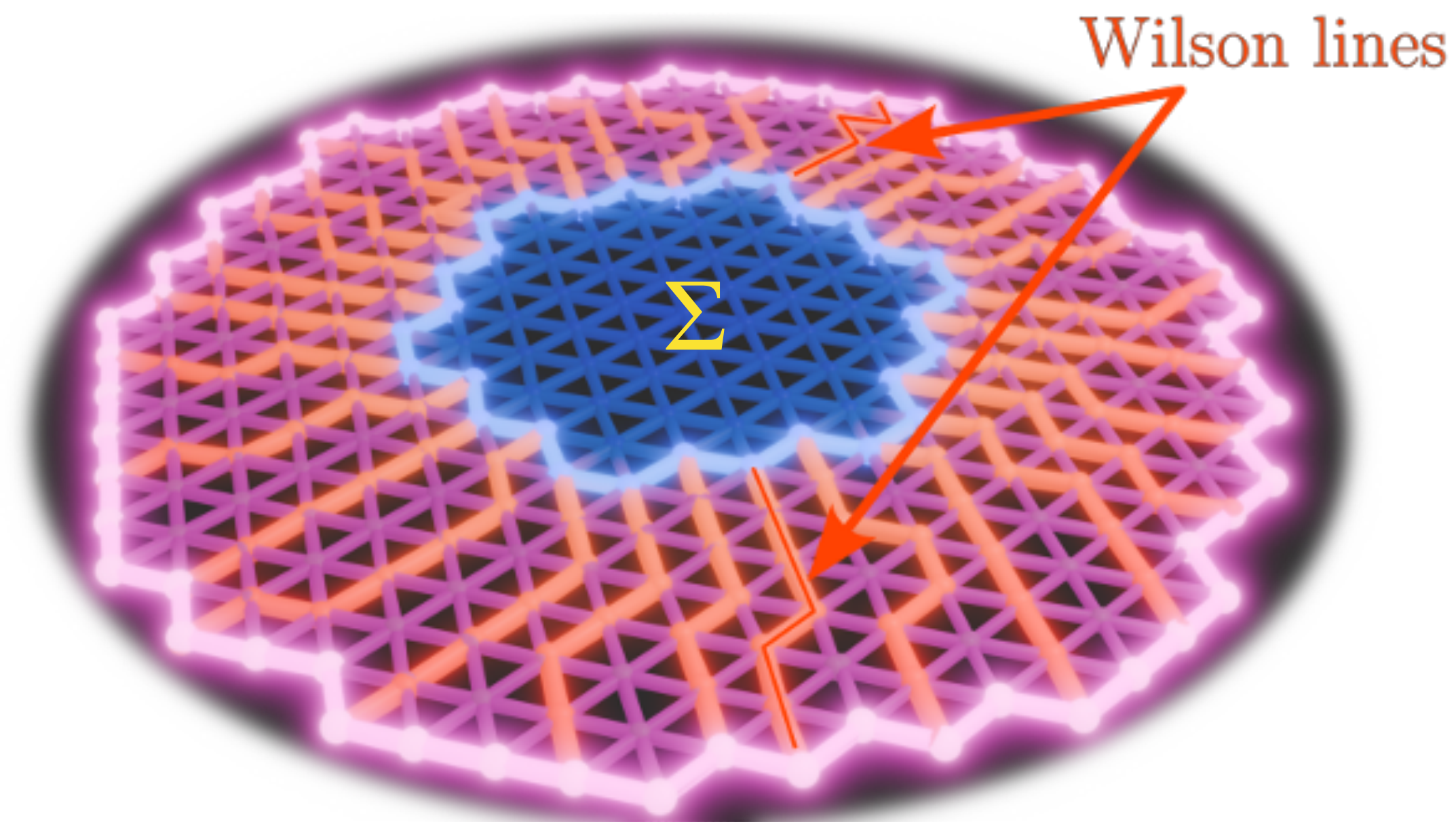


incomplete ideal QRF: $\mathcal{H}_{\tilde{\Phi}} = L^2(G^{V_{\partial\Sigma}^{-1}})$

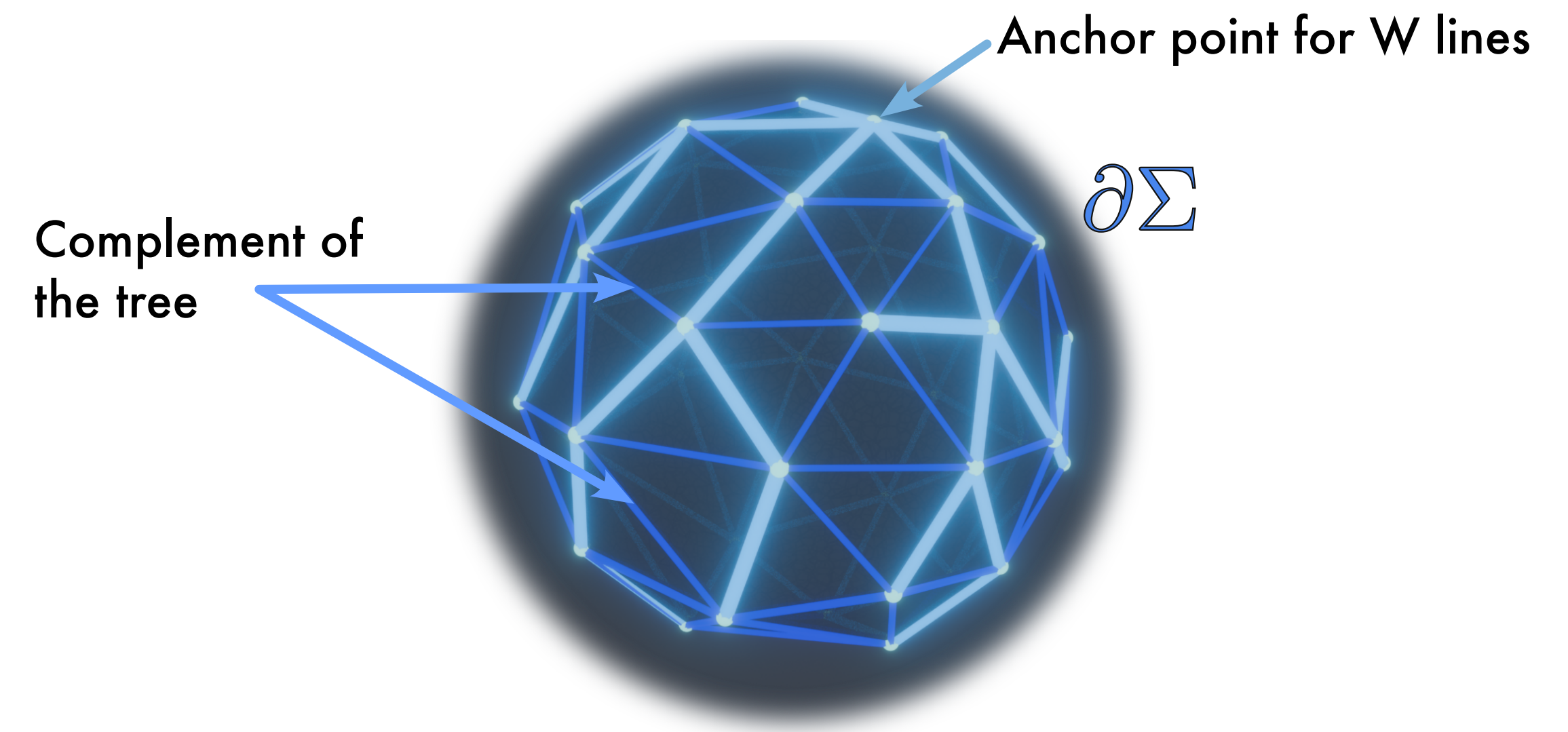
⇒ can't deparametrize/dress at anchor point N

Extrinsic and intrinsic gauge QRFs on the lattice

Extrinsic: congruence of Wilson lines in the complement



Spanning tree anchored on node N of the corner



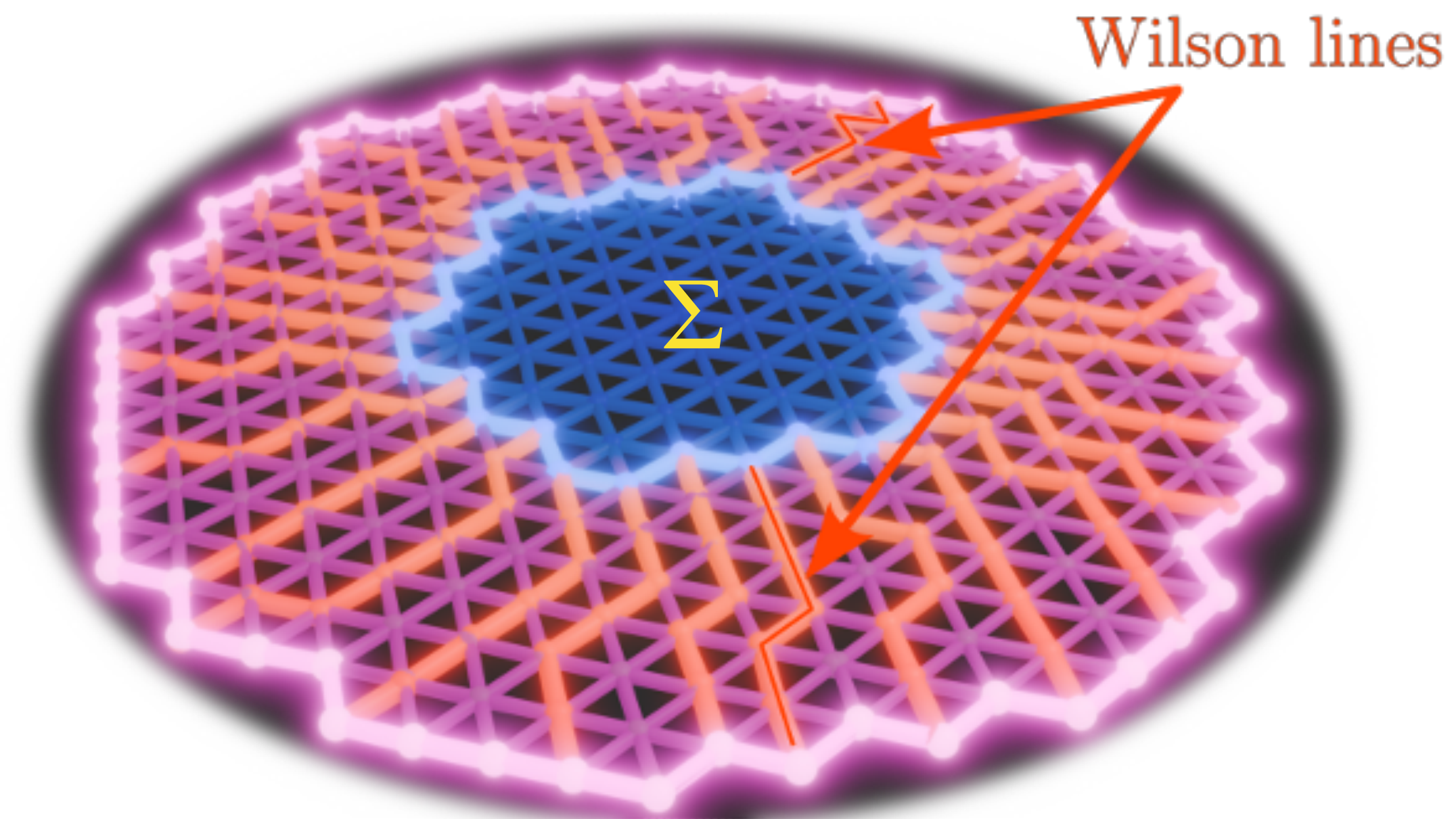
complete ideal QRF: $\mathcal{H}_{\Phi} = L^2(G^{V_{\partial\Sigma}})$

incomplete ideal QRF: $\mathcal{H}_{\tilde{\Phi}} = L^2(G^{V_{\partial\Sigma}^{-1}})$

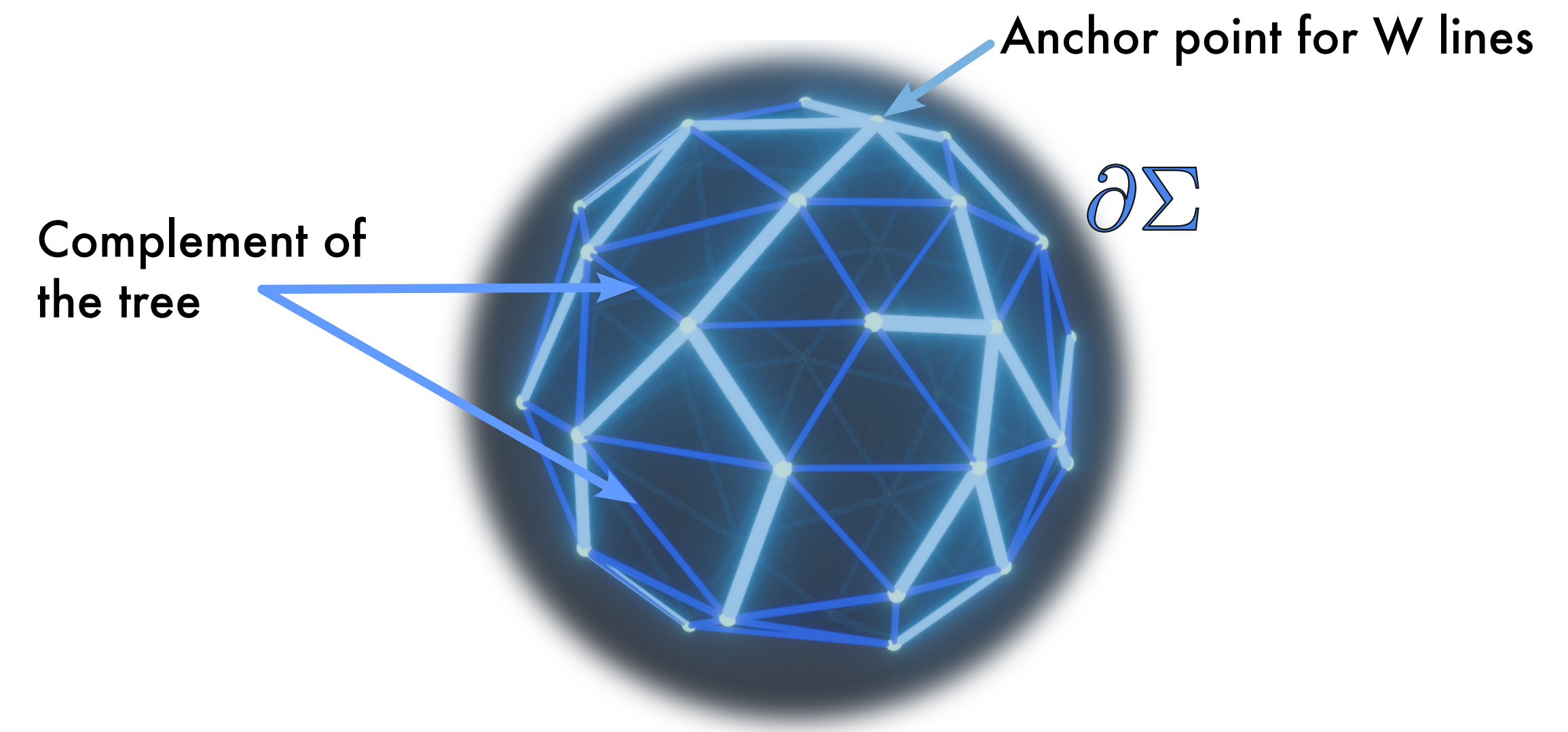
$\Rightarrow$ both can be used to build gauge-inv. description (algebras) for subregion, overall impose strong G -invariance

How do you build physical QRFs for the corner group?

Extrinsic: congruence of Wilson lines in the complement



Spanning tree anchored on node N of the corner

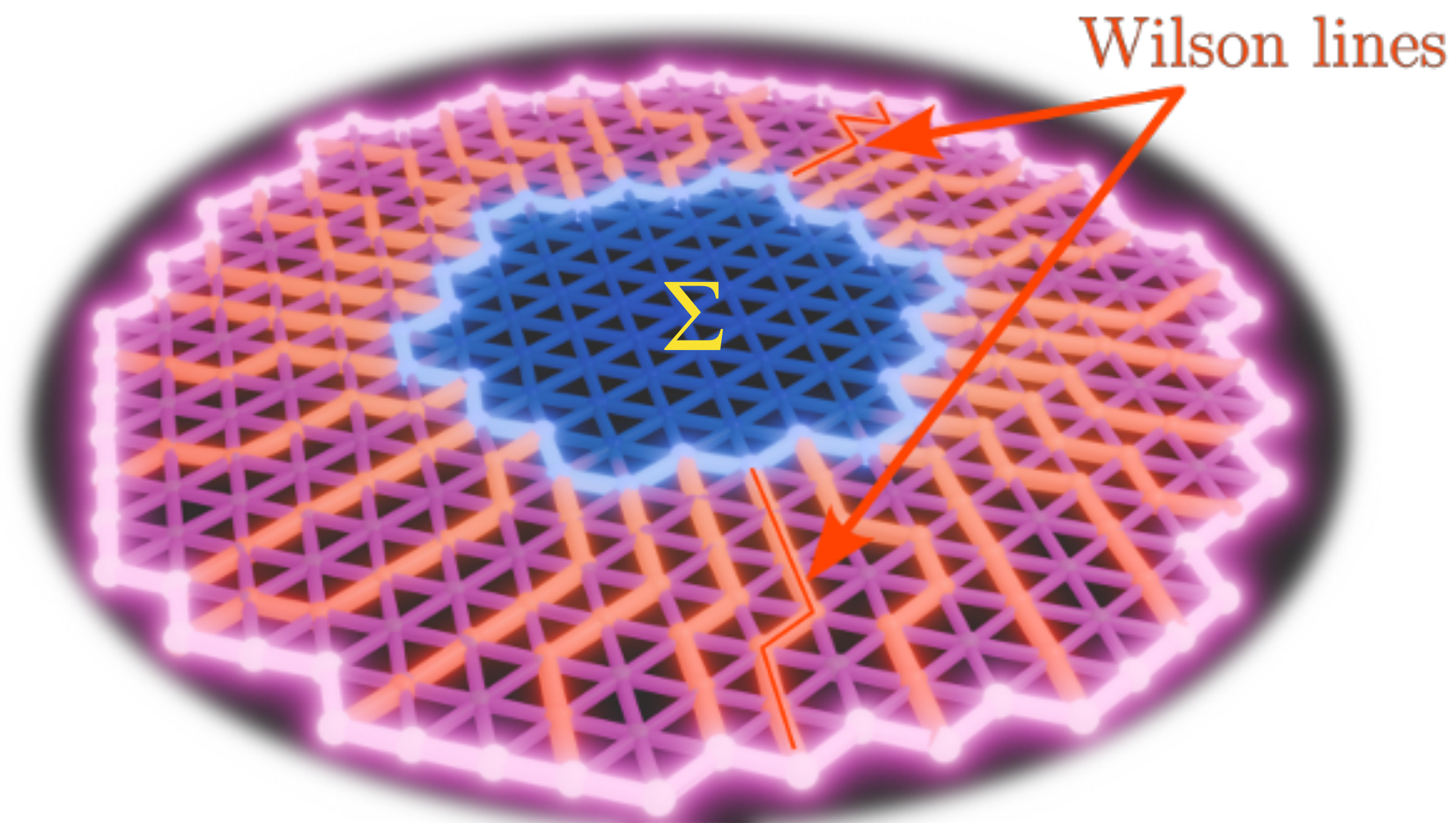


How do you build physical QRFs for the corner group?

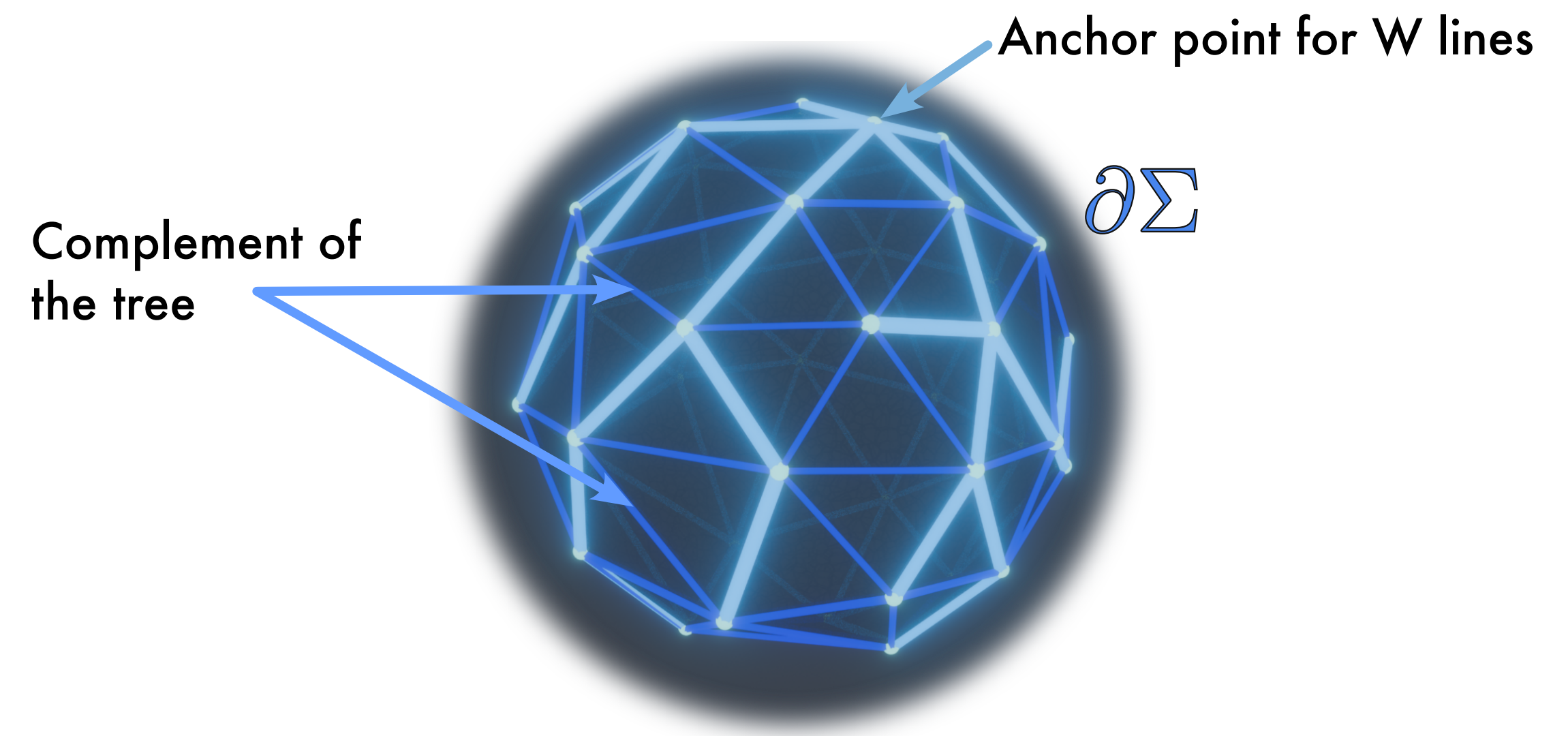
reorientations of extr. QRF Φ

Extrinsic: congruence of Wilson lines in the complement

Spanning tree anchored on node N of the corner



complete ideal QRF: $\mathcal{H}_{\Phi} = L^2(G^{V_{\partial\Sigma}})$



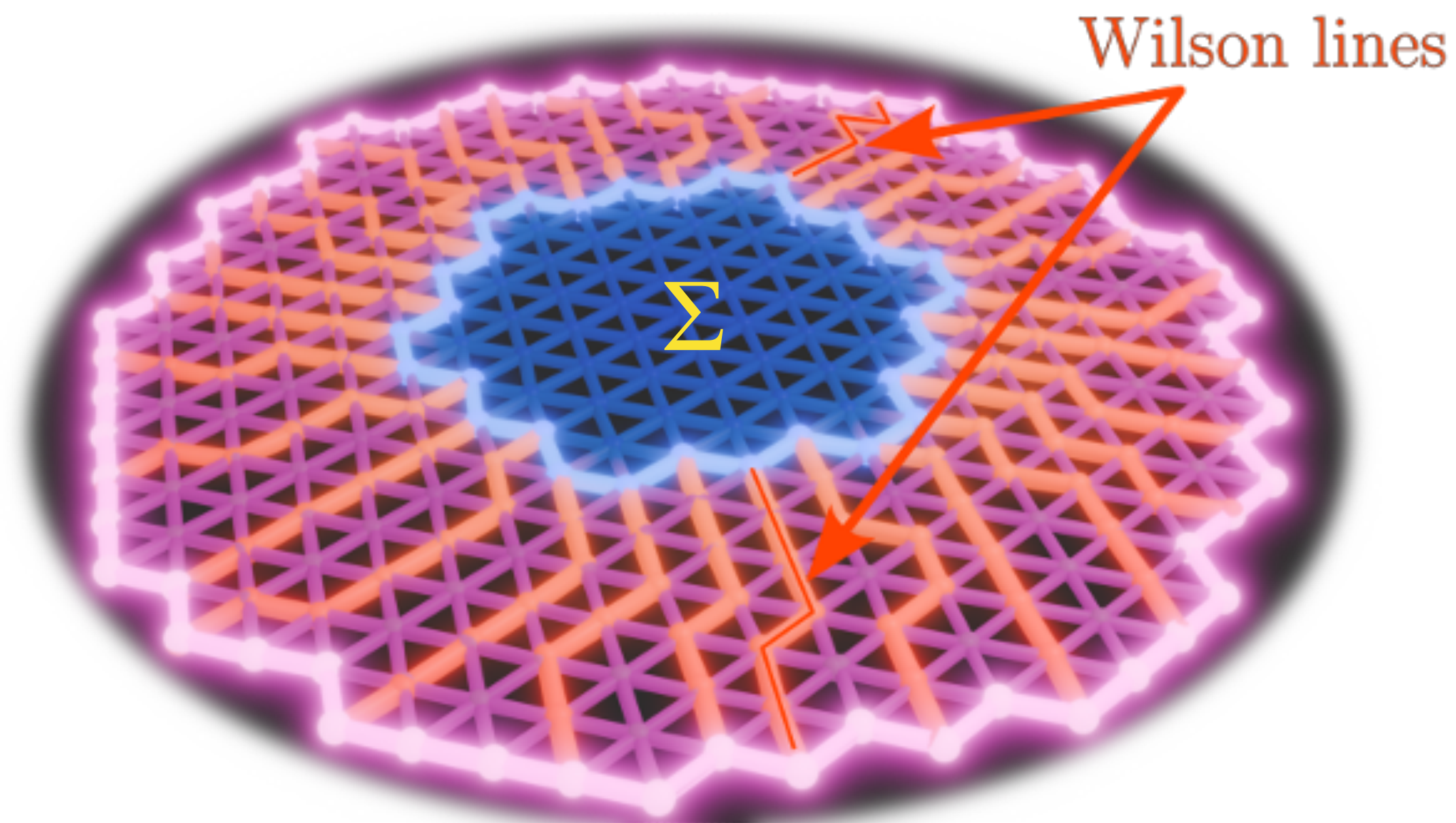
incomplete ideal QRF: $\mathcal{H}_{\tilde{\Phi}} = L^2(G^{V_{\partial\Sigma}^{-1}})$

How do you build physical QRFs for the **corner group**?

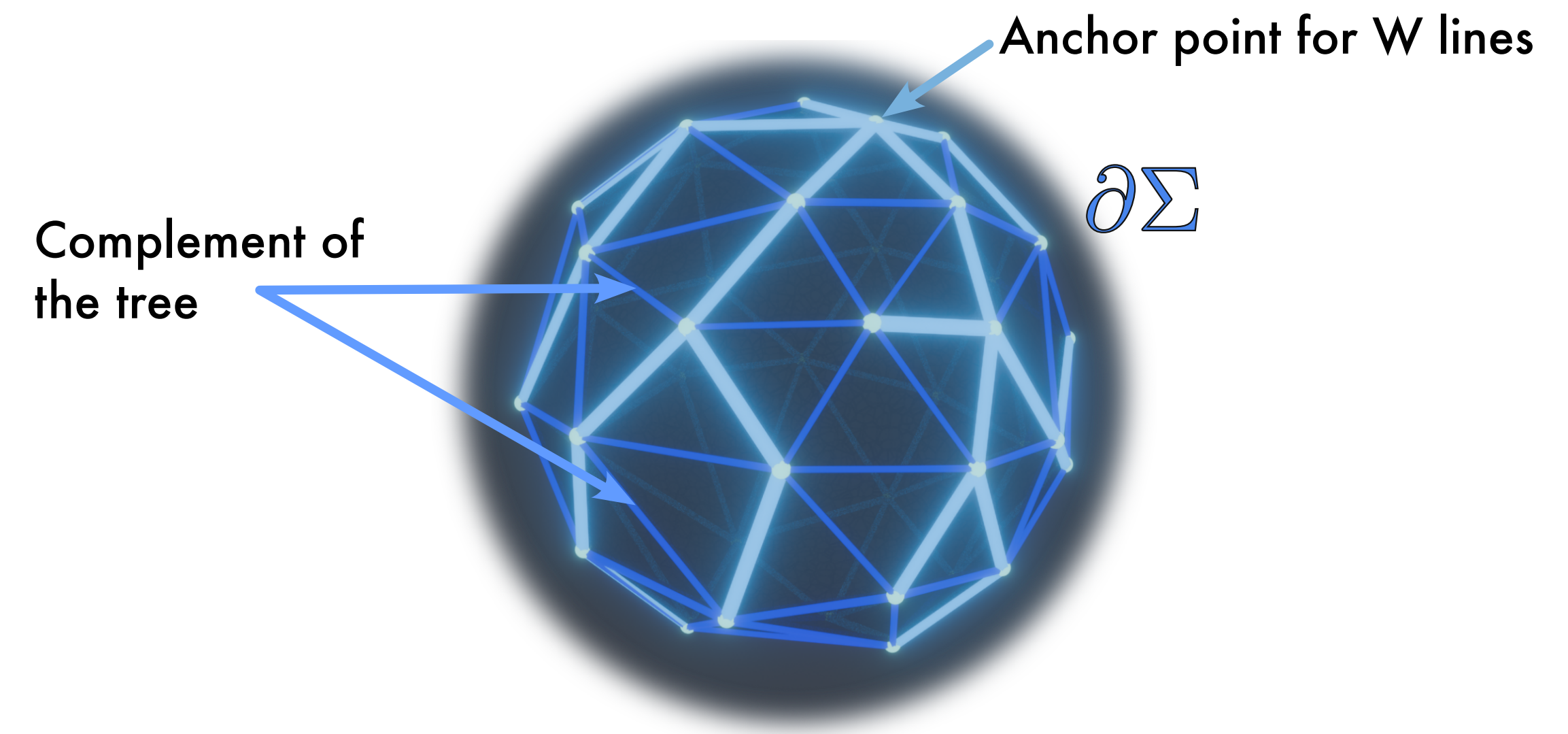
reorientations of extr. QRF Φ

Extrinsic: congruence of Wilson lines in the complement

Spanning tree anchored on node N of the corner



complete ideal QRF: $\mathcal{H}_{\Phi} = L^2(G^{V_{\partial\Sigma}})$



incomplete ideal QRF: $\mathcal{H}_{\tilde{\Phi}} = L^2(G^{V_{\partial\Sigma}^{-1}})$

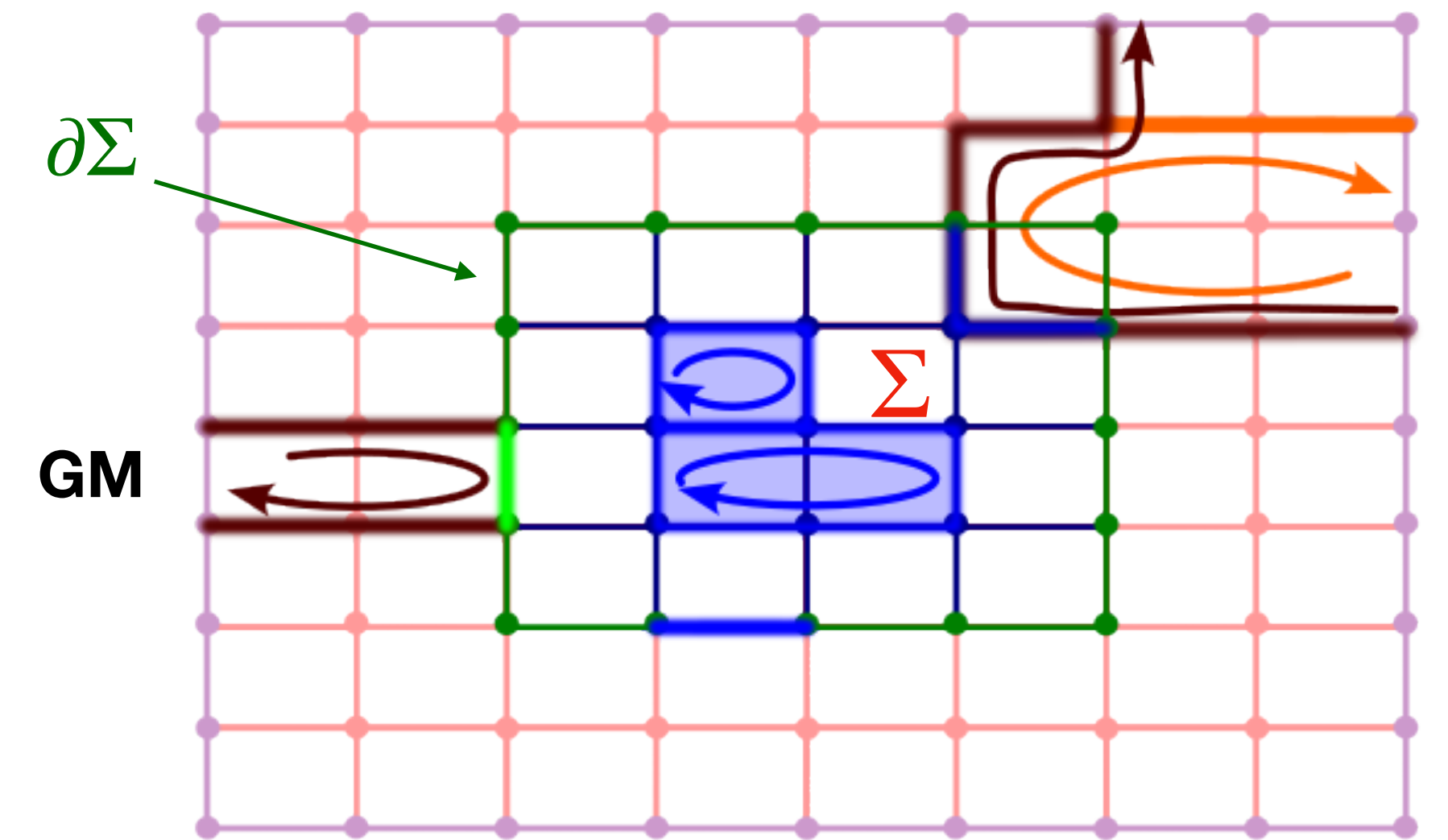
relational observables $O_{|\Phi}(\tilde{\Phi})$ describing intr. relative to extr. gauge QRF
constitute a gauge-inv./phys. QRF for the corner group (Goldstone modes)

Averaging over the Goldstone mode

$\mathcal{A}_{\text{ext}} = (\mathcal{A}_{\Sigma} \otimes \mathcal{A}_{\Phi})^G$ contains the Goldstone mode:

rel. obs. describing int. rel. to ext. QRF $O_{|\Phi}^g(\tilde{\Phi})$

$\Rightarrow$ nonlocal **gauge-inv.** QRF for the electric corner group



Averaging over the Goldstone mode

$\mathcal{A}_{\text{ext}} = (\mathcal{A}_{\Sigma} \otimes \mathcal{A}_{\Phi})^G$ contains the Goldstone mode:

rel. obs. describing int. rel. to ext. QRF $O_{|\Phi}^g(\tilde{\Phi})$

$\Rightarrow$ nonlocal **gauge-inv.** QRF for the electric corner group

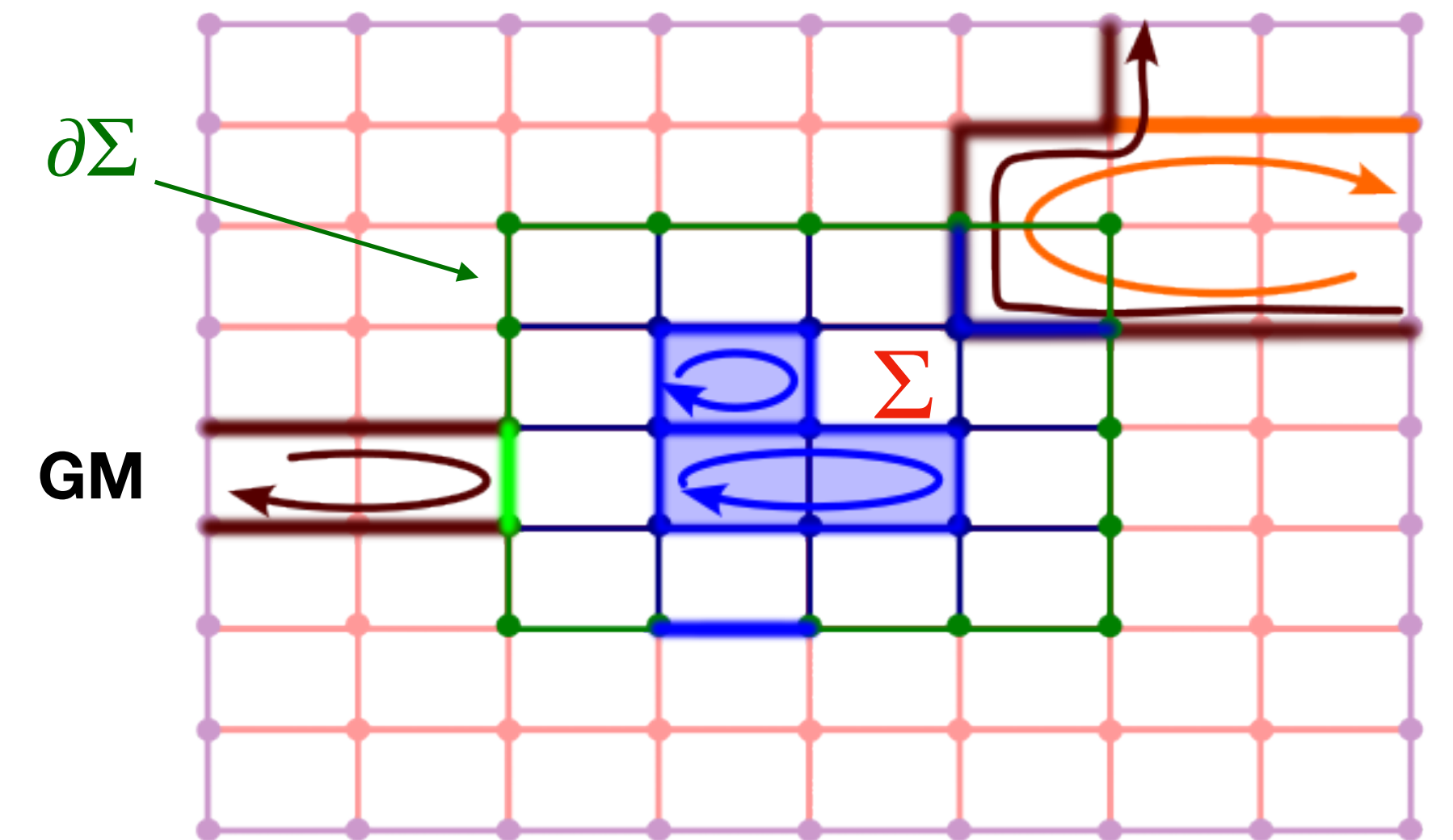
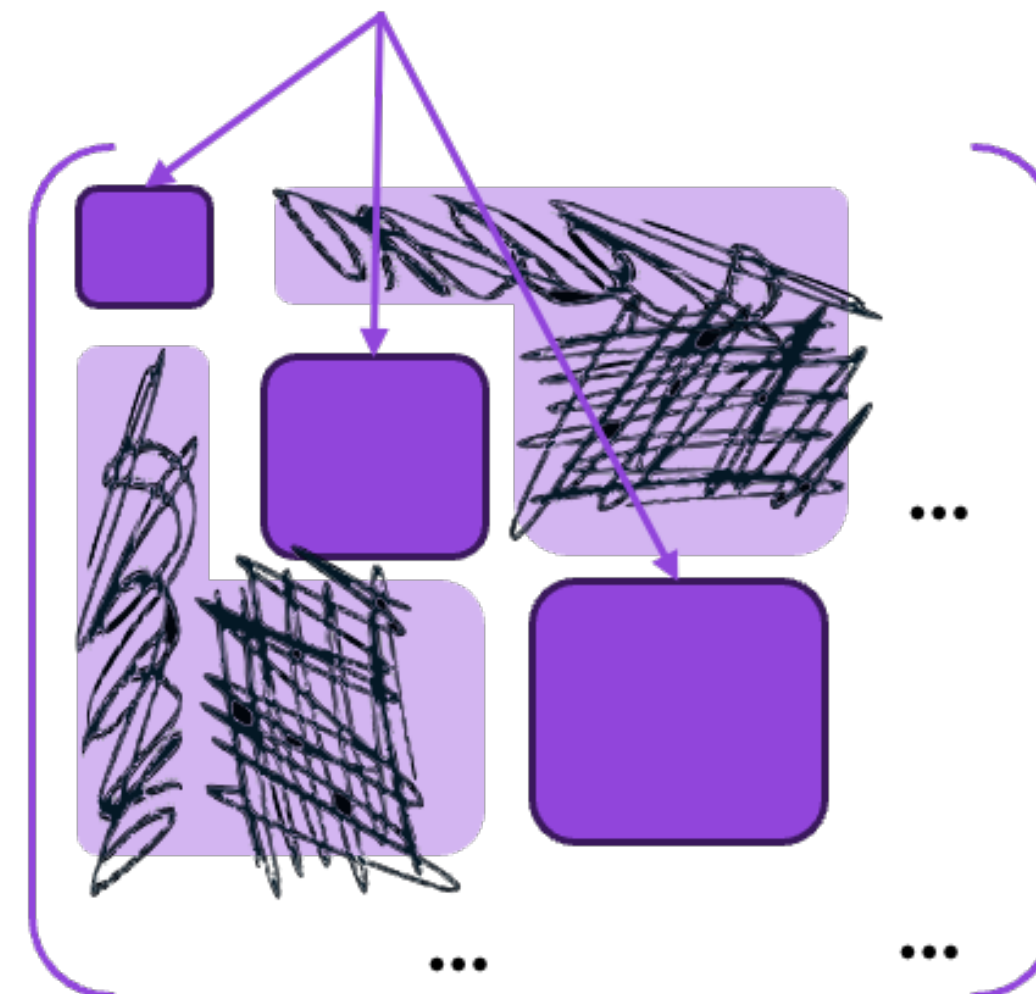
if no access to it $\Rightarrow$ average over its orientations/corner group

$$\mathcal{A}_E = \mathbb{G}_{\partial\Sigma}(\mathcal{A}_{\text{ext}}) = \int_{G^{V_{\partial\Sigma}}} dg V_{\Phi}(g) \mathcal{A}_{\text{ext}} V_{\Phi}^{\dagger}(g)$$

Corner group irreps sectors

electric corner charge superselection:

$\mathcal{A}_E =$



Averaging over the Goldstone mode

$\mathcal{A}_{\text{ext}} = (\mathcal{A}_{\Sigma} \otimes \mathcal{A}_{\Phi})^G$ contains the Goldstone mode:

rel. obs. describing int. rel. to ext. QRF $O_{|\Phi}^g(\tilde{\Phi})$

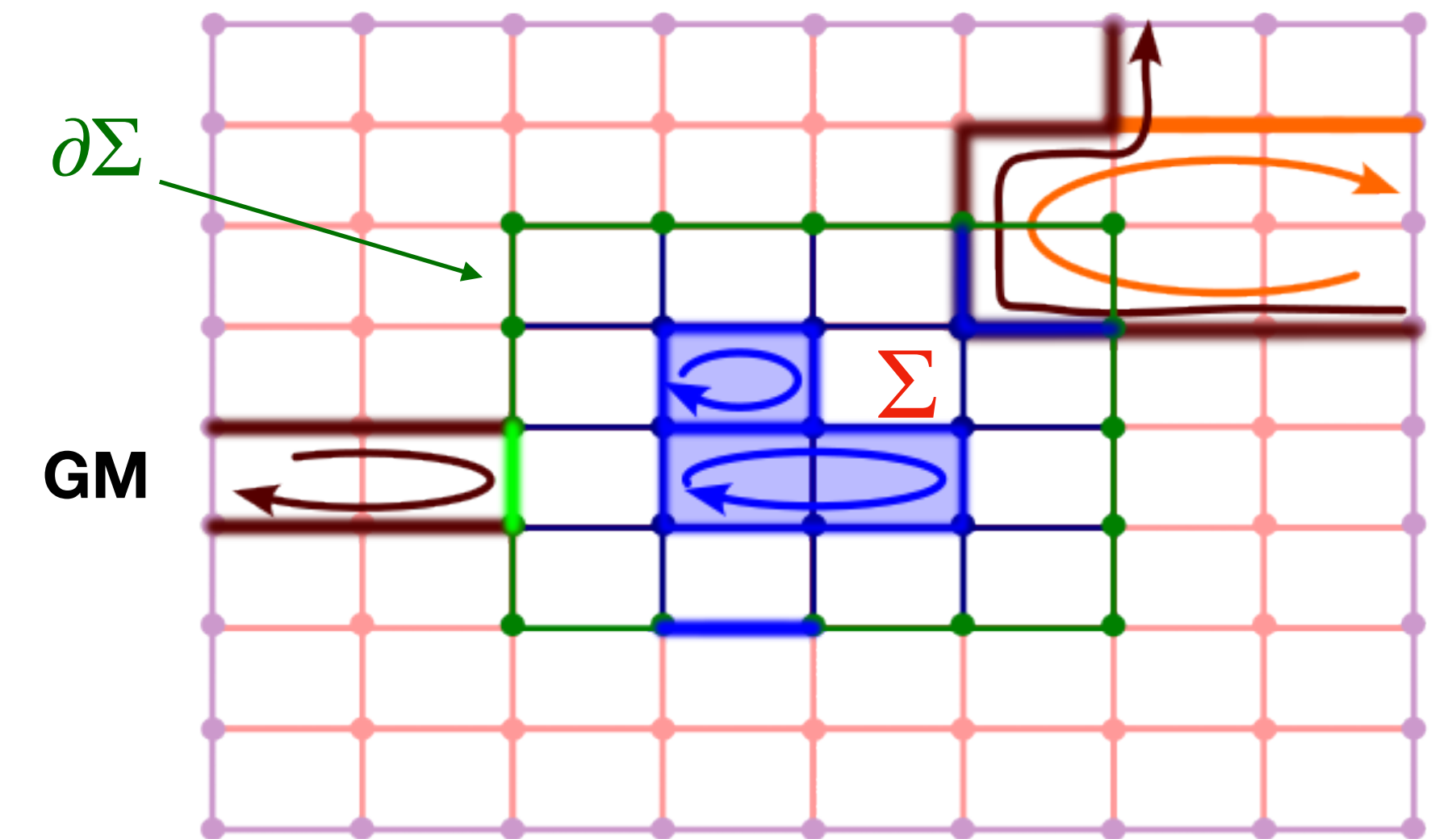
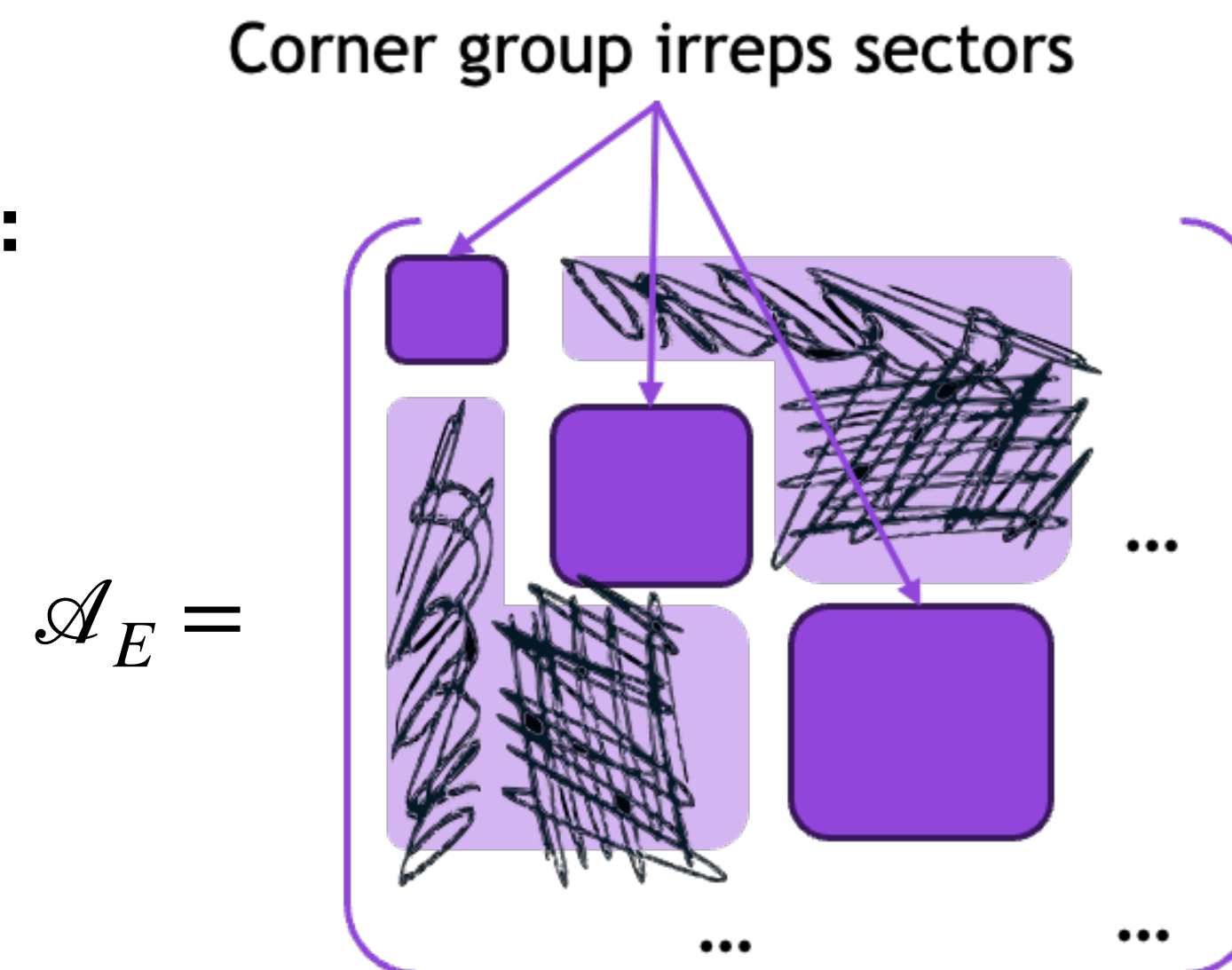
$\Rightarrow$ nonlocal **gauge-inv.** QRF for the electric corner group

if no access to it $\Rightarrow$ average over its orientations/corner group

$$\mathcal{A}_E = \mathbb{G}_{\partial\Sigma}(\mathcal{A}_{\text{ext}}) = \int_{G^{V_{\partial\Sigma}}} dg V_{\Phi}(g) \mathcal{A}_{\text{ext}} V_{\Phi}^{\dagger}(g)$$

electric corner charge superselection:

$\Rightarrow$ yields standard electric center algebra of entanglement entropy computations



$$S_{\Sigma}(\rho) = \sum_c p_c (S_{\text{vN}}(\rho_c) + \log d_c) - \sum_c p_c \ln p_c$$

[Casini, Huerta, Rosabal '13; Donnelly '12;
Buividovich, Polikarkov '08; Ghosh et al '15; ...]

Conclusions

- **Extension of SR covariance structures into quantum realm**
based on internal QRFs $\Rightarrow$ in terms of group structures really the same as in SR
- **Systematic method for changing QRF perspectives**
accommodates RFs in relative superposition
- **Gauge-inv. subsystems depend on choice of QRF (subsystem relativity)**
 $\Rightarrow$ correlations, thermal properties, dynamics, depend on frame
- **Two distinct notions of external frame independence**
 $\Rightarrow$ depending on whether G is a physical symmetry or redundancy
- **Peaceful coexistence of physical and gauge QRFs in lattice gauge theory**
 $\Rightarrow$ edge vs. Goldstone modes and recovery of standard entanglement entropy computations