

Canonical Quantum Gravity

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TOC:

- Classical Hamiltonian formulation of GR
- Canonical quantisation principles
- Loop quantum gravity (LQG) realisation

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Initial value formulation

- Canonical approach to classical & quantum GR has long tradition
- Classical: Initial value formulation, numerical integration of Einstein equations, BH merger simulations ... [Arnowitt, Deser, Misner, ...]
- Quantum: QG as a constrained QFT [Bergmann, DeWitt, Dirac, Komar, Wheeler, ...]
- **Key assumption:** Globally hyperbolic spacetimes (M, g)
- $\Rightarrow M \cong \mathbb{R} \times \sigma$ [Geroch, Sanchez & Bilal]
- M can be foliated by leaves $t \mapsto \Sigma_t \cong \sigma$ ("3+1" split: space+time)
- Lapse, shift parametrisation of foliation

$$ds^2 = -[l dt]^2 + q_{ab}[dx^a + s^a dt][dx^b + s^b dt]$$

- Legendre transform for constrained systems (Dirac algorithm): $(g, \dot{g}) \mapsto (q, p)$ (3-metric on σ , conj. momentum), Poisson brackets $\{., .\}$. Similar for matter.
- plus: spatial diffeomorphism and Hamiltonian constraints V_a ; $a = 1, 2, 3$, S (temporal-spatial, temporal-temporal components of Einstein eqns)

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Classical Hypersurface Deformation Algebroid (HDA)

- Let $Q := \sqrt{\det(q)}$, smeared constraints:
 $V[u] = \int_{\sigma} d^3z u^a V_a; \quad S[f] = \int_{\sigma} d^3z f S$

- Classical HDA $\mathfrak{h}$ [Hojman, Kuchar, Teitelboim]

$$\{V[u], V[v]\} = -V[[u, v]], \quad \{V[u], S[f]\} = -S[u[f]],$$

$$\{S[f], S[g]\} = -V[q^{-1} (M dN - N dM)]$$

- Observations:
 - encodes local spacetime diffeomorphism covariance
 - $\mathfrak{h}$ not a Lie algebra due to q^{-1} : “open algebra, algebroid”
- Anomaly-free representation of $\mathfrak{h}$ major challenge in quantum reduction (QR) approach to CQG

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Classical reduction approach

- **Totally constrained** Hamiltonian $C(T) := V[s] + S[I]$
- V, S generator of **gauge transf.** rather than physical motion
- Observables, physical motion? (“problem of time”)
- Relational (Dirac) observables $\Leftrightarrow$ Gauge fixing
- Split geometry, matter canonical pairs into 2 sets $(q, p) = ((Q, P), (X, Y))$ called “**true**” (observable) and “**gauge**” (not observable)
- Solve constr. $C = 0$ for $Y = -h(X; Q, P)$, impose gauge cond. $G := X - k = 0$
- Solve stability condition $\frac{d}{dt}G = 0 \Leftrightarrow \{C(T), X\} = \partial_t k$ for $T = T_*$
- Reduced (true, physical,...) **Hamiltonian** for $\mathcal{O} = \mathcal{O}(Q, P)$ [Hanson, Regge, Teitelboim]

$$\{H, \mathcal{O}\} := \{C(T), \mathcal{O}\}_{X=k, Y=-h, T=T_*} \Rightarrow H = \int d^3z \dot{k} \cdot h(X = k, Q, P)$$

- How practical depends on matter content (e.g. scalar fields)

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Generalities

- **Non-perturbative:** PT applied to QG: non-renormalisable [Goroff, Sagnotti 85]
- **Background independence:** no metric singled out by Einstein eqns.
- **Continuum theory:** no ad hoc discrete input
- **Standard QFT:** construct QFT of gravitational field in 3+1
- **Algebraic QFT (AQFT) methods:** Algebras $\mathfrak{A}$, exp. val. states ω
- **Standard AQFT:** Fix classical background g_0 , construct $\mathfrak{A}_{g_0}(O)$
 If O, O' **causally disjoint** wrt g_0 then $[\mathfrak{A}_{g_0}(O), \mathfrak{A}_{g_0}(O')] = 0$
- BI violation, in QG $g = g_0$ not fixed, even operator, **fluctuating causality**
- **Algebraic QG (AQG):**
 Suppose 1. $\omega(g^2) - \omega(g)^2$ “small” (semiclassical), 2. O, O' **causally disjoint** wrt $\omega(g)$ then show $\omega([\mathfrak{A}(O), \mathfrak{A}(O')])$ small
- **Difference:**
 AQFT: g_0 distinguished, used to construct $\mathfrak{A}_{g_0}(O)$
 AQG: $g_0 \not\mathfrak{A}, \mathfrak{A}(O)$ BI, **quantum locality** rel. to state only
- How to construct $\mathfrak{A}(O)$ in CQG? Like in usual QFT in CST:
 Solve Heisenberg picture EOM wrt physical H from “time zero” fields

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