Quantum entropies in Quantum Field Theory

Albert Much

University of Leipzig, Institute for Theoretical Physics, D-04081 Leipzig, Germany

Causal Fermion Systems 2025, New Perspectives in Mathematics and Physics, Regensburg 2025

Outline for this Talk

Beckenstein Bound

► Generalize to QFT (Casini) ⇒ Relative Entropy

Crash Course Modular Theory

- Relative Entropy defined via Modular operator
- ► Applications to noncommutative QFT

The Classical Bekenstein Bound

What happens if we throw a hot cup of (filter-)coffee into a black hole?



Black Hole Entropy

For a stationary, non-charged, non-rotating black hole, the entropy is proportional to the event horizon area:

$$S_{BH} = \alpha M^2$$
.

Dropping Matter into a Black Hole (Poor man's derivation)

Consider a system outside the black hole with mass $m \ll M$ and entropy S. The total initial entropy is:

$$S^- = S_{BH} + S$$

Dropping m into the black hole (neglecting energy losses) gives:

$$S_{BH+m} = \alpha (M+m)^2 \approx \alpha M^2 + 2\alpha Mm = S_{BH} + 2\alpha Mm$$

Entropy Increase Condition

$$S_{BH+m}-S^-\geq 0$$

Thus,

$$S < 2\alpha Mm$$

The Bekenstein Bound

If BH is large enough to swallow the system (radius R), and identifying E=m, one finds:

Bekenstein Bound [Beckenstein 1981]

For a system with radius R and energy E:

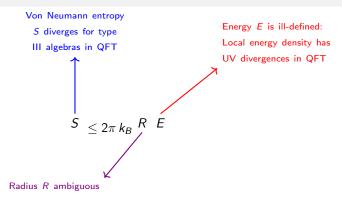
$$S \leq 2\pi k_B R E$$

where k_B is the Boltzmann constant.

- Universal entropy bound on physical systems
- Derived from black hole entropy formula

Quantum Field Theory Formulation of the Beckenstein Bound

Bekenstein Bound: Formulation in QFT



Global vs local: "For global states, E is well-defined, but R is not; for localized states R is well-defined, vacuum entanglement (pair creation) makes both S and E UV-sensitive and ill-defined without careful renormalization."

QFT Challenges: Divergences

Problem [Casini 08]

Naive application of von Neumann entropy in QFT leads to divergences

- Physical origin: UV divergences from vacuum fluctuations
- ► Solution: Subtract vacuum contributions

Regularized Entropy

$$S_V = S(\rho_V) - S(\rho_V^0) = -\operatorname{tr}(\rho_V \log \rho_V) + \operatorname{tr}(\rho_V^0 \log \rho_V^0)$$

where:

- V: spatial region of interest
- $ightharpoonup
 ho_V$: reduced density matrix of the excited state associated with V
- \triangleright $S(\rho_V^0)$: von Neumann entropy of vacuum state

Local Hamiltonian and Vacuum State

Vacuum Density Matrix

$$\rho_V^0 = \frac{e^{-K}}{\operatorname{tr} e^{-K}}$$

where K is the local Hamiltonian (self-adjoint, dimensionless).

Example

For half-space $V = \{x^1 > 0\}$, K is the boost operator:

$$K = L_{01} = 2\pi \int_{x^1 > 0} x^1 T_{00}(0, x) dx$$

▶ T_{00} : Hamiltonian density operator, physical units: $[R \cdot E] = [L_{01}]$

Regularized Energy

Problem

Expectation value $\langle K \rangle_{\rho_V}$ diverges due to vacuum fluctuations

Solution

Take for right hand side, difference between excited and vacuum states:

$$K_V = \operatorname{tr}(K\rho_V) - \operatorname{tr}(K\rho_V^0)$$

QFT Bekenstein Bound

$$S_V \leq K_V$$

Relative Entropy

Definition

$$S_{rel,V} = \operatorname{tr}
ho_V(\log(
ho_V) - \log(
ho_V^0)) \geq 0$$

Theorem (Positivity of Relative Entropy)

The relative entropy $S_{rel,V}$ is always non-negative.

Quantum relative entropy $S(\rho \parallel \sigma)$ has two key interpretations:

- 1. State Distinguishability, where reference state σ
- 2. Quantifies unexpected information content when encountering ρ while expecting σ

Key Result

The positivity of relative entropy implies the Bekenstein bound!

Full Circle

Black hole entropy \rightarrow Bekenstein bound \rightarrow QFT formulation \rightarrow Relative entropy positivity

Mathematical Framework to study relative Entropy

Tomita-Takesaki-Modular theory [1967, 1970]

Tomita-Takesaki Theory & Relative Entropy

Tomita-Takesaki Theory: Foundations

Basic Setup:

- \blacktriangleright $(\mathcal{M}, \mathcal{H}, |\Omega\rangle)$: von Neumann algebra \mathcal{M} on Hilbert space \mathcal{H}
- $ightharpoonup |\Omega\rangle$: cyclic and separating vector for $\mathcal M$

Tomita Operator *S*:

- ▶ Definition: $SA|\Omega\rangle = A^*|\Omega\rangle$ for all $A \in \mathcal{M}$
- S is closable and densely defined

Polar Decomposition:

$$S = J\Delta^{1/2}$$

- ▶ J: modular conjugation (anti-linear, unitary, $J^2 = 1$)
- ▶ JMJ = M' (commutant relation)
- Δ: modular operator (positive, self-adjoint)

Modular Group and KMS States

Modular Automorphism Group:

- $ightharpoonup \sigma_t(A) = \Delta^{it} A \Delta^{-it} \text{ for } t \in \mathbb{R}$
- ▶ One-parameter group of automorphisms: $\sigma_t \circ \sigma_s = \sigma_{t+s}$
- ▶ Covariance: $\sigma_t(\mathcal{M}) = \mathcal{M}$

KMS State Definition:

▶ A state ω is KMS at $\beta > 0$ for $\{\sigma_t\}$ if, for all $A, B \in \mathcal{M}$, the map $t \mapsto \omega(A \sigma_t(B))$ extends analytically to $0 < \text{Im } z < \beta$ and satisfies

$$\omega(A\,\sigma_t(B)) = \omega(\sigma_{t+i\beta}(B)\,A).$$

Connection to Modular Theory:

• State $\omega(\cdot) = \langle \Omega | \cdot | \Omega \rangle$ is KMS for $\{ \sigma_t \}$ at $\beta = 1$

Relative Modular Theory

Setup:

- ightharpoonup Two cyclic and separating states: ω and ω'
- ► Corresponding vectors: $|\Omega\rangle$ and $|\Omega'\rangle$

Relative Tomita Operator:

- ▶ Definition: $S_{\omega',\omega} A |\Omega'\rangle = A^* |\Omega\rangle$ for $A \in \mathcal{M}$
- ▶ Polar decomposition: $S_{\omega',\omega} = J_{\omega',\omega} \Delta_{\omega',\omega}^{1/2}$

Properties:

- ▶ $J_{\omega',\omega}$: relative modular conjugation
- $ightharpoonup \Delta_{\omega',\omega}$: relative modular operator
- $I_{\omega',\omega}\mathcal{M}J_{\omega',\omega}=\mathcal{M}'$

Araki-Uhlmann Relative Entropy

Definition:

$$S_{
m rel}(\omega',\omega) = -\langle \Omega | \logig(\Delta_{\omega',\omega}ig) | \Omega
angle$$

Special Case - Unitarily Equivalent States:

- ▶ If $\omega(U \cdot U^{-1}) = \omega'(\cdot)$ for some unitary U
- Araki [1977]-Uhlmann [1977] formula:

$$S_{\mathrm{rel}}(\omega',\omega) = i \frac{d}{dt} \langle \Omega \big| U \Delta^{it} U^* \Delta^{-it} \Omega \rangle \big|_{t=0}$$

Applications of relative Entropy in QFT

Bisognano-Wichmann Theorem & KMS

Physical Context:

- ▶ Rindler wedge: $W_R := \{x = (x_0, x_1, ..., x_n) \in \mathbb{R}^d : x_1 > |x_0|\}$
- Wedge algebra: local observables in the right wedge

Bisognano-Wichmann Theorem:

- ▶ Modular operator: $\Delta = e^{2\pi L_{01}}$
- ▶ Boost generator: $L_{01} = \int d^{d-1}x x^1 T_{00}(x)$

KMS Connection:

- lacktriangle Vacuum state $|\Omega
 angle$ is KMS for wedge algebra at $eta=2\pi$
- Modular group σ_t implemented by boosts: $\sigma_t(A) = U(\Lambda(-2\pi t))AU(\Lambda(2\pi t))$
- Unruh effect

Explicit Relative Entropy Calculation

Setup:

- ightharpoonup Reference state: vacuum ω
- ► Excited state: ω' generated by $U|Ω⟩ = e^{iφ(f)}|Ω⟩$
- lacktriangle Operators localized in right wedge \mathcal{W}_R

Calculation:

$$S_{0}(\omega', \omega) = i \frac{d}{dt} \langle \Omega | U \Delta^{it} U^{*} \Delta^{-it} \Omega \rangle \Big|_{t=0}$$

$$= i \frac{d}{dt} \langle \Omega | e^{i\phi(f)} e^{2\pi i t L_{01}} e^{-i\phi(f)} e^{-2\pi i t L_{01}} \Omega \rangle \Big|_{t=0}$$

$$= 2\pi \int_{x^{0}=0, x^{1}>0} x^{1} T_{00} dx$$

 T_{00} is the energy density smeared with f.

Recent Works on S_{rel} : Incomplete Chronological Overview

2019

- ▶ Longo, R. (2019) "Entropy of Coherent Excitations"
 - Explicit formula for vacuum relative entropy of coherent states on wedge algebras

- ► Hollands, S. (2019) "Relative entropy for coherent states in chiral CFT"
 - ▶ Result: Relative entropy between vacuum and exponentiated stress tensor equals *c* times Schwarzian action of diffeomorphisms

- ► Hollands, S.; Ishibashi, A. (2019) "News vs Information"
 - ► Relative entropy in linearized quantum gravity around black holes
 - ► Main result: $\frac{d}{dt}(S + A/4) = 2\pi F$ (entropy-area-flux relation)

Recent Works on S_{rel}

- Casini H., Grillo S., Pontello, D. (2019) "Relative entropy for coherent states from Araki formula"
 - Computes vacuum—coherent relative entropy in the Rindler wedge via Araki's formula, matching the canonical result.

2020

► Faulkner T., Hollands S., Swingle B., Wang Y. (2020)-"Approximate recovery and relative entropy I. general von Neumann subalgebras"

Recent Works on Relative Entropy: 2021

2021

- ▶ Edoardo D'Angelo (2021)- "Entropy for spherically symmetric, dynamical black holes from the relative entropy between coherent states of a scalar quantum field"
- ► Kurpicz, F.; Pinamonti, N.; Verch, R. (2021) "Temperature and entropy-area relation of quantum matter near spherically symmetric outer trapping horizons"
- ► Ciolli, F.; Longo, R.; Ranallo, A.; Ruzzi, G. (2021) "Relative entropy and curved spacetimes"
 - QNEC inequality for coherent states in curved spacetime

Some Works on Relative Entropy: 2023-2025

2022

- ▶ Galanda, S.; AM; Verch, R. (2023) "Relative Entropy of Fermion Excitation States on the CAR Algebra"
 - ► Counterpart to CCR results, uses self-dual CAR algebra

2024

- Fröb, M.; AM; Papadopoulos, K. (2024) "Relative Entropy in de Sitter is a Noether Charge"
 - Connection of relative entropy to Noether charge of modular flow translations

► Fröb, M.; Sangaletti, L. (2024) - "Petz-Rényi relative entropy in QFT from modular theory"

- ► Finster F., Jonsson R., Lottner M., Murro S., AM, (2024)-Notions of Fermionic Entropies of a Causal Fermion System
 - Defines fermionic von Neumann, entanglement, and relative entropies for causal fermion systems via reduced one-particle density operator
 - connects to modular-theoretic computations of relative entropy

2025

▶ Hollands, S.; Longo, R. (2025) - "A New Proof of the QNEC"

- ► Finster, F.; AM (2025) "The Relative Fermionic Entropy in Two-Dimensional Rindler Spacetime"
 - Fermionic relative entropy using modular theory and density operators
 - Application to non-unitary excitations in Rindler spacetime

Applications of relative Entropy in noncommutative QFT

Relative Entropy in QG - Intro NCQFT

[DFR01] QFT in a NC Minkowski-spacetime is represented on $\mathcal{V}\otimes\mathscr{F}_s(\mathcal{H})$, where \mathcal{V} is the representation space of X

$$[X_{\mu}, X_{\nu}] = i\theta_{\mu\nu},$$

where $\mu, \nu = 0, \dots, 3$ and θ is a skew-symmetric matrix

$$heta_{\mu
u} = egin{pmatrix} 0 & \Theta & 0 & 0 \ -\Theta & 0 & 0 & 0 \ 0 & 0 & 0 & \Theta' \ 0 & 0 & -\Theta' & 0 \end{pmatrix},$$

with $\Theta, \Theta' \in \mathbb{R}$ and $\Theta \geq 0$.

Application to NCQFT

[GL07] proved that ϕ_{\otimes} can be written on the Fock space $\mathscr{F}_s(\mathscr{H})$ by existence of a unitary map \mathcal{U} from $\mathcal{U}: \mathcal{V} \otimes \mathscr{F}_s(\mathscr{H}) \to \mathscr{F}_s(\mathscr{H})$

$$egin{aligned} \phi_{ heta}(f) &:= \int d^4x \, f(x) \, \phi_{ heta}(x) \ &= \int rac{d^3 \mathsf{p}}{\omega_\mathsf{p}} \, \left(f^-(p) \, e^{-ip heta P} a(p) + f^+(p) \, e^{ip heta P} a^*(p)
ight), \end{aligned}$$

with $\omega_p=+\sqrt{\vec{p}^2+m^2}$ and $p=(\omega_p,\vec{p})$ and $f\in\mathscr{S}(\mathbb{R}^4)$ and P is the momentum operator.

Furthermore, the authors proved that ϕ_{θ} (and $\phi_{-\theta}$) is a wedge local field.

Solution

Theorem

The deformed relative entropy $S_{\theta}(\omega',\omega)$ is up to first order in Θ explicitly given by

$$S_{\theta}(\omega',\omega) = i \frac{d}{dt} \langle \Omega | e^{i\phi_{\theta}(f_{\Theta'})} e^{2\pi i t L_{\mathbf{0}\mathbf{1}}} e^{-i\phi_{\theta}(f_{\Theta'})} e^{-2\pi i t L_{\mathbf{0}\mathbf{1}}} \Omega \rangle \Big|_{t=0}$$
$$= S_{0}(\omega',\omega) + \frac{8\pi}{3} \Theta \left(\int d\mu(k) \omega_{k} |f_{\Theta'}^{+}(k)|^{2} \right)^{2},$$

where $S_0(\omega', \omega)$ is the undeformed relative entropy.

The deformed version of the Beckenstein bound

$$\boxed{S \leq 2\pi \, R \, E + \frac{8\pi}{3} \, \Theta \, m^2.}$$

By coefficient comparison of the Beckenstein bound but not neglecting the m^2 term

$$S \leq 2\pi R E + 4\pi G m^2,$$

we identify Θ with the Planck-length squared I_P^2 , i.e.

$$\Theta = \frac{3}{2}G = \frac{3}{2}I_P^2.$$

Thank you for your Attention!



Armin Uhlmann 1971



Huzihiro Araki 2009