

Spontaneous wavefunction collapse: grounded in the familiar

Lajos Diósi

Wigner Research Center for Physics



Eötvös Loránd University



*

Supports by: National Research, Development and Innovation Office for
“Frontline” Research Excellence Program (Grant No. KKP133827).

9 Oct 2025

Abstract

A simple introduction to the concept of spontaneous collapse is possible directly on the familiar grounds of standard measurement, collapse, and their mathematics. Reasons for the assumption of spontaneous collapse will be explained, together with certain contexts, open issues, and perspectives.

Abstract

Quantum Measurement 1932-

Quantum Monitoring 1988-

Spontaneous Quantum Measurements/Monitorings: Why?

Is the moon there ... ?

Spontaneous Quantum Monitoring

Ghirardi-Rimini-Weber(-Bell) 1986-

Continuous Spontaneous Localization 1990-

Gravity-Related Spontaneous Collapse 1987-

Penrose vs D, D versus Penrose

Summary and More Things to Talk About

Quantum Measurement 1932- von Neumann

Observable $\hat{x}$, measurement precision (unsharpness) σ .

Irreversible, alters energy/momentum of the measured system.

Selective measurement – **Collapse** – nonlinear, stochastic

$$\psi \longrightarrow \text{DEVICE} \longrightarrow \psi|_x = \mathcal{N} \exp\left(\frac{(x - \hat{x})^2}{4\sigma^2}\right) \psi$$

Random outcome x , probability:

$$p(x) = \frac{1}{\sqrt{2\pi\sigma^2}} \psi^\dagger \exp\left(\frac{(x - \hat{x})^2}{2\sigma^2}\right) \psi$$

Non-selective measurement – **Decoherence** – linear, deterministic

$$\psi\psi^\dagger \longrightarrow \text{DEVICE} \longrightarrow \int p(x) \psi|_x \psi|_x^\dagger dx$$

$$\hat{\rho} \longrightarrow \text{DEVICE} \longrightarrow \mathcal{M}_{CP} \hat{\rho}$$

Quantum Monitoring 1988- D,Belavkin,Wiseman-Milburn

Monitoring = time-continuous measurement

E.g.: Measurements at unsharpness $\sigma \rightarrow \infty$, repeated at rate $\lambda \rightarrow \infty$, constant $\gamma = \lambda/\sigma^2$ is the strength of monitoring.

Irreversible, alters energy/momentum of the monitored system.

Selective monitoring – **Dynamical Collapse** –

Nonlinear Stochastic Schrödinger Eq. (**NLSSE**) for $\Psi|_{\{x\}}$:

$$\frac{d\Psi|_{\{x\}}}{dt} = -\frac{i}{\hbar}\hat{H}\Psi|_{\{x\}} - \frac{\gamma}{8}(\hat{x} - \langle\hat{x}\rangle)^2\Psi|_{\{x\}} + \frac{\sqrt{\gamma}}{2}(\hat{x} - \langle\hat{x}\rangle)\Psi|_{\{x\}}w_t$$

Random outcome (signal) $x_t = \langle\hat{x}\rangle_t + w_t/\sqrt{\gamma}$

Non-selective monitoring – **Dynamical Decoherence** –

Linear, deterministic Master Eq. (**ME**) for $\hat{\rho} = \langle\Psi|_{\{x\}}\Psi|_{\{x\}}^\dagger\rangle_{stoch}$:

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma}{8}[\hat{x}, [\hat{x}, \hat{\rho}]]$$

Spontaneous Measurements/Monitorings: Why?

Standard Quantum Measurements/Monitorings

- ▶ are the unique mechanism of classical data emergence
- ▶ assume measuring devices, which is
 - ▶ necessary and confirmed in micro-systems
 - ▶ annoying in macro-systems.

That's why we choose the simplest idea, and assume Spontaneous Quantum Measurements/Monitorings

- ▶ weak and ignorable in micro-systems
- ▶ amplified and dominant in macro-systems

Spontaneous Measurements/Monitorings retain the math of standard M's/M's, explain the emergence of classicality in macroscopic d.o.f., at the price:
tiny irreversibility and non-conservation of energy/momentum.

Is the moon there ... ?

Is the moon there when nobody looks? [Mermin 1985]

Two quotations:

Pascual Jordan: "Observations not only disturb what has to be measured, they produce it....We compel [the electron] to assume a definite position.... We ourselves produce the results of measurements."

Abraham Pais: "I recall that during one walk Einstein suddenly stopped, turned to me and asked whether I really believed that the moon exists only when I look at it."

[back](#)

Spontaneous Quantum Monitoring

Concept: Spontaneous (deviceless) Measurements universally present every time and everywhere, parametrized so that their effect is ignorable for microscopic degrees of freedom, is amplifying on the mezoscales, and becomes dominant on macroscales.

Choice of monitored observables, of parametrization, do matter!

model	monitored observables	effective rate of monitoring	spatial resolution
<u>GRW</u> 1986	particle positions $\hat{x}_\alpha$	$\lambda = 10^{-17}/\text{sec}$ collapse rate	$\sigma = 10^{-5}\text{cm}$ localization length
<u>DP</u> 1987	mass density field $\hat{\mu}(x)$	$G/\hbar$	$\sigma > 10^{-9}\text{cm}$ short length cutoff
<u>CSL</u> 1990	mass density field $\hat{\mu}(x)$	$\gamma \propto \lambda \sigma^3 / m_0^2$ $\lambda < 10^{-12}/\text{sec}$	$\sigma = 10^{-5}\text{cm}$ localization length

Ghirardi-Rimini-Weber(-Bell) 1986-

Position $\hat{x}_\alpha$ of each elementary particle is spontaneously measured at rate $\lambda = 10^{-17}/s$ and precision $\sigma = 10^{-5}cm$:

$$\psi \longrightarrow \mathcal{N} \exp \left(\frac{(x_\alpha - \hat{x}_\alpha)^2}{4\sigma^2} \right) \psi$$

Random outcome x_α , probability:

$$p_\alpha(x_\alpha) = \frac{1}{\sqrt{2\pi\sigma^2}} \psi^\dagger \exp \left(\frac{(x_\alpha - \hat{x}_\alpha)^2}{2\sigma^2} \right) \psi$$

No effect on microscopic d.o.f. but on massive d.o.f., e.g., the c.o.m. $\hat{x}_{cm}$ of $N = 10^{23}$ particles, when the effective collapse rate becomes $10^8/s$.

Rigid body c.o.m. Dynamic Decoherence ME:

$$\frac{d\hat{\rho}_{cm}}{dt} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}_{cm}] - \frac{N\lambda}{8\sigma^2} [\hat{x}_{cm}, [\hat{x}_{cm}, \hat{\rho}_{cm}]], \quad (\text{if } \Delta x_{cm} \ll \sigma)$$

$$\tau_{GRW}^{-1} = \frac{1}{8} N\lambda (\Delta x_{cm}/\sigma)^2$$

back

Continuous Spontaneous Localization 1990- G-R-Pearle

Smeared spatial mass density distribution:

$$\hat{\mu}_\sigma(x) = \sum_\alpha m_\alpha G_\sigma(x - \hat{x}_\alpha)$$

spontaneously monitored at (effective) rate $\gamma_{CSL} = (2\sqrt{\pi}\sigma)^3\lambda$.

Selective monitoring – Dynamical Collapse – NLSSE:

$$\begin{aligned} \frac{d\Psi}{dt} = & -\frac{i}{\hbar}\hat{H}\Psi - \frac{\gamma_{CSL}}{2m_0^2} \int (\hat{\mu}_\sigma(x) - \langle \hat{\mu}_\sigma(x) \rangle)^2 d^3x \Psi \\ & + \frac{\sqrt{\gamma_{CSL}}}{m_0} \int (\hat{\mu}_\sigma(x) - \langle \hat{\mu}_\sigma(x) \rangle) \Psi w_t(x) d^3x \end{aligned}$$

Measured signal: $\mu_t(x) = \langle \hat{\mu}_\sigma(x) \rangle_t + \frac{1}{2}m_0\gamma_{CSL}^{-1/2}w_t(x)$

$$\langle w_t(x)w_\tau(y) \rangle_{stoch} = \delta(t - \tau)\delta(x - y)$$

Nonselective monitoring – Dynamical Decoherence – ME:

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] - \frac{\gamma_{CSL}}{2m_0^2} \int [\hat{\mu}_\sigma(x), [\hat{\mu}_\sigma(x), \hat{\rho}]] d^3x$$

[back](#)

Gravity-Related Spontaneous Collapse 1987- D, Penrose

Smeared mass density $\hat{\mu}_\sigma(x)$ is spontaneously monitored by $1/r$ -correlated spontaneous measurements at (eff.) rate $G/\hbar$.

Selective monitoring – Dynamical Collapse – NLSSE:

$$\begin{aligned} \frac{d\Psi}{dt} = & -\frac{i}{\hbar} \hat{H} \Psi - \frac{G}{2\hbar} \int [\hat{\mu}_\sigma(x) - \langle \hat{\mu}_\sigma(x) \rangle] [\hat{\mu}_\sigma(y) - \langle \hat{\mu}_\sigma(y) \rangle] \frac{d^3x d^3y}{|x-y|} \Psi \\ & + \sqrt{G/\hbar} \int (\hat{\mu}_\sigma(x) - \langle \hat{\mu}_\sigma(x) \rangle) \Psi w_t(x) d^3x \end{aligned}$$

Measured signal: $\mu_t(x) = \langle \hat{\mu}_\sigma(x) \rangle_t + \frac{1}{2} \sqrt{\hbar/G} w_t(x)$

$$\langle w_t(x) w_\tau(y) \rangle_{stoch} = \delta(t - \tau) / |x - y|$$

Nonselective monitoring – Dynamical Decoherence – ME:

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}] - \frac{G}{2\hbar} \int \int [\hat{\mu}_\sigma(x), [\hat{\mu}_\sigma(y), \hat{\rho}]] \frac{d^3x d^3y}{|x-y|}$$

[back](#)

Penrose vs D, D versus Penrose

Penrose

D

Incompatibility between Quantum Theory and
General Relativity

Newtonian Gravity

Ultimate uncertainty of
time-translation Killing vector

Newton potential Φ

Measure of uncertainty:

no GR covariant form

simple non-relativistic form:

$$\int |\nabla \delta \Phi|^2 dV$$

no dynamical eqs.

NLSE and ME

no “price”

energy/mom nonconservation

Decay rate of macroscopic superposition $|M; \text{Left}\rangle + |M; \text{Right}\rangle$:
postulated

derived from ME

$$\tau_{DP}^{-1} = \Delta E_G / \hbar$$

$$\Delta E_G = 2U_G(LR) - U_G(LL) - U_G(RR)$$

Summary and More Things to Talk About

Nature arranges extremely weak (otherwise standard) quantum measurements of particle positions (GRW) or of mass config's (DP, CSL) everywhere and -time in the Universe, while measuring DEVICES are kept hidden from our eyes and physics. The ensuing 'spontaneous' collapse/decoherence yields a unified theory of the micro- and macroworld, without resorting to the device-related measurement postulate.

Further things to talk about:

- ▶ Experiments: two ways to go
 - ▶ Gran Sasso Experiment* (correct and silly reactions)
- ▶ Relativity? Non-Markovianity? Dissipativity?
- ▶ Is collapse (NLSE) testable, or just decoherence (ME) is?
- ▶ DP, Tilloy-D, Oppenheim: A healthier semiclassical gravity
- ▶ Hybrid classical-quantum ME = equivalent formalism
- ▶ ...

* 2021 Donadi-Piscicchia-Curceanu-D-Laubenstein-Bassi