THERMODYNAMICS OF SPACETIME: FROM UNIMODULAR GRAVITY TO QUANTUM GRAVITY PHENOMENOLOGY

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MOTIVATION

- Understand the interface between gravitational dynamics and thermodynamics: One step further
- Classical structure of gravity: GR vs UG
- On the search for quantum gravity...
 - Quantum phenomenological models to get insight into possible effects

- Look at the entropy: Quantum modifications of entropy
 - Same qualitative behavior in all the models
- Could we obtain phenomenological modified equations of motion from it?
 - General phenomenological quantum gravity dynamics
 - Quantum cosmology: Resolution of the singularity?

OUTLINE

- Einstein equations from thermodynamics: Basic concepts and UG vs GR
- Modified entropy
- > Quantum phenomenological equations of motion
- Interpretation of the modified dynamics and application to cosmology
- Discussion
- Further perspectives

EINSTEIN EQUATIONS FROM THERMODYNAMICS

- Local-causal horizon
- Gravitational dynamics equilibrium condition for maximal entropy



Quantum correlations across horizon + matter-energy crossing it

GEODESIC LOCAL CAUSAL DIAMONDS

- > In an arbitrary P choose any unit timelike vector n^{μ} and construct Riemann normal coordinates (RNC)
 - Metric expansion around P

$$g_{\mu\nu}(x) = \eta_{\mu\nu} - \frac{1}{3} R_{\mu\alpha\nu\beta} \left(P\right) x^{\alpha} x^{\beta} + O\left(x^3\right)$$

Area of the 2-sphere

$$\mathcal{A} = 4\pi l^2 - \frac{4\pi}{9} l^4 G_{00} \left(P \right) + O\left(l^5 \right)$$

Spacetime region causally determined by this ball is called geodesic local causal diamond (GLCD)

Conformal Killing vector —> conformal Killing horizon

Quantum correlations across horizon + matter-energy crossing it

- ► Entanglement entropy: $S = \eta \mathcal{A}$ → Bekenstein entropy $S_{BH} = \frac{k_B \mathcal{A}}{4l_P^2}$
- Entropy of matter: in terms of entanglement for GLCD explicitly evaluated for small perturbations from vacuum

Clausius entropy: Equal to leading order for conformal fields (without gravitational dynamics)

[A.A-S, M. Liska, PRD 102 (2020) 10, 104056]



Key (an minimal) assumptions: - <u>Einstein equivalence principle</u> - Entropy of horizon = <u>Bekenstein entropy</u>

[A.A-S, M. Liska, arxiv 2305.11756 [gr-qc]]



[T. Jacobson, PRL 116 (2016) 20, 201101]

• Different methods: Clausius entropy and entanglement entropy

[A.A-S, M. Liska, PRD 102 (2020) 10, 104056]



Unimodular gravity: Restricting full diffeomorphism invariance of GR

$$g = -1 \longrightarrow \sqrt{-g} = \omega_0$$

The simplest action $S = \int_{\Omega} \left[\frac{c^4}{16\pi G} \left(R - 2\bar{\Lambda} \right) + \mathcal{L}_{\text{matter}} \right] \omega_0 \mathrm{d}^4 x$

when varying it with respect to the metric

$$R_{\mu\nu} - \frac{1}{4}Rg_{\mu\nu} = \frac{8\pi G}{c^4} \left(T_{\mu\nu} - \frac{1}{4}Tg_{\mu\nu}\right)$$

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• Transverse diff invariance does not suffice to ensure $T_{\mu \ ;\nu}^{\ \nu}=0$



 Λ appears as an integration constant

• UG is fully equivalent to GR

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Weyl transverse gravity: transverse diffeomorphisms + Weyl invariance

- > Why obtaining Weyl transverse gravity from thermodynamics of spacetime?
 - Local energy conservation must be assumed separately
 - Λ appears as an integration constant
 - Only conformally invariant part couples to gravity > conformal invariance
 - GR and UG same classical dynamics

[A.A-S, M. Liska, PRD 102 (2020) 10, 104056]

MODIFIED ENTANGLEMENT ENTROPY

Leading order quantum gravity modifications to Bekenstein entropy

$$S_{BH,q} = \frac{k_B \mathcal{A}}{4l_P^2} + \mathcal{C}k_B \ln\left(\frac{\mathcal{A}}{\mathcal{A}_0}\right) + O\left(\frac{k_B l_P^2}{\mathcal{A}}\right)$$

- Predicted by different theories of quantum gravity and phenomenological approaches
- Calculations entanglement entropy —> local causal horizons

QUANTUM PHENOMENOLOGICAL EQUATIONS OF MOTION

EQUATIONS OF MOTION FROM ENTANGLEMENT EQUILIBRIUM

[T. Jacobson, PRL 116 (2016) 20, 201101]

Maximal vacuum entanglement hypothesis (MVEH):

"When the geometry and quantum fields are simultaneously varied from maximal symmetry, the entanglement entropy in a small geodesic ball is maximal at fixed volume"

Consider the GLCD carried out in a maximally symmetric spacetime (MSS), and we simultaneously vary geometry and state of quantum fields:

$$\delta S_e + \delta S_m = 0$$

• On one side: δS_e with the logarithmic correction

$$S_{e,q} = \eta \mathcal{A} + k_B \mathcal{C} \ln \frac{\mathcal{A}}{\mathcal{A}_0} + O\left(\frac{k_B l_P^2}{\mathcal{A}}\right)$$
$$\delta \mathcal{A}|_V$$
$$\delta S_{e,q} = S_{e,q} - S_{e,q}^{MSS} = \eta \delta \mathcal{A} + k_B \mathcal{C} \frac{\delta \mathcal{A}}{\mathcal{A}_{MSS}} - k_B \frac{\mathcal{C}}{2} \left(\frac{\delta \mathcal{A}}{\mathcal{A}_{MSS}}\right)^2 + O\left((\delta \mathcal{A})^3\right)$$

- On the other side: δS_m
 - Conformal Killing vector Conformal Killing horizon with T
 - Quantum fields in the thermal state ____ Minkowski vacuum (EEP)
 - Variation of entropy of quantum fields (Temperatures cancel out)

$$\delta S_m = \frac{2\pi k_B c}{\hbar} \frac{4\pi}{15} l^4 \left(\delta \langle T_{00} \rangle + \delta X \right) - 4\psi l^2 C^2 \frac{2\pi k_B l_P^2}{\hbar c^3} \frac{4\pi}{15} l^4 \left(\delta \langle T_{00} \rangle + \delta X \right) + O\left(l^5\right)$$

 \succ Considering the MVEH: $\delta S_{e,q} + \delta S_m = 0$

Requirement of recovering Einstein equations for $C \rightarrow 0$ fixes $\eta = k_B/4l_P^2$

$$\Rightarrow \begin{cases} \text{Due to EEP} \Rightarrow \text{every spacetime point } P \\ \text{GLCD in } P \Rightarrow \text{any arbitrary unit timelike vector } n^{\mu} \end{cases}$$

• We obtain:

$$S_{\mu\nu}(P) n^{\mu}n^{\nu} + \frac{Cl_{Pl}^{2}}{30\pi}S_{\alpha\beta}(P) n^{\alpha}n^{\beta}S_{\mu\nu}(P) n^{\mu}n^{\nu} - \Phi(P) = \frac{8\pi G}{c^{4}}\delta\langle T_{\mu\nu}(P)\rangle(P) n^{\mu}n^{\nu}$$

with: $S_{\mu\nu} = R_{\mu\nu} - Rg_{\mu\nu}/4$

We differentiate respect to n^{μ} getting a system of conditions — $\blacktriangleright \Phi$

Finally we obtain:

$$S_{\mu\nu} - \frac{Cl_{Pl}^2}{30\pi} S_{\mu\lambda} S^{\lambda}_{\ \nu} - \frac{8\pi G}{c^4} \delta \langle T_{\mu\nu} \rangle = -^{(0)} \Phi g_{\mu\nu}$$

Taking the trace
$$S_{\mu\nu} - \frac{Cl_{Pl}^2}{30\pi} S_{\mu\lambda} S^{\lambda}_{\ \nu} + \frac{Cl_{Pl}^2}{120\pi} \left(R_{\kappa\lambda} R^{\kappa\lambda} - \frac{1}{4} R^2 \right) g_{\mu\nu} = \frac{8\pi G}{c^4} \left(\delta \langle T_{\mu\nu} \rangle - \frac{1}{4} \delta \langle T \rangle g_{\mu\nu} \right)$$

(Classically Einstein equations recovered by imposing: $T_{\mu~;\nu}^{~~\nu} = 0 \longrightarrow \Lambda$) [A.A-S, M. Liska, JHEP 12 (2020) 196]

EQUATIONS OF MOTION FROM CLAUSIUS ENTROPY

Clausius entropy of matter crossing arbitrary bifurcate null surfaces in a curved spacetime



The entropy balance equation yields to a similar expression than previous procedure, and finally we obtain

$$S_{\mu\nu} - \frac{Cl_P^2}{18\pi} S_{\mu\lambda} S^{\lambda}_{\ \nu} + \frac{Cl_P^2}{72\pi} \left(R_{\kappa\lambda} R^{\kappa\lambda} - \frac{1}{4} R^2 \right) g_{\mu\nu} = \frac{8\pi G}{c^4} \left(T_{\mu\nu} - \frac{1}{4} T g_{\mu\nu} \right)$$

[A.A-S, M. Liska, JHEP 12 (2020) 196]

INTERPRETATION OF THE EQUATIONS

$$S_{\mu\nu} - D l_P^2 S_{\mu\lambda} S_{\nu}^{\lambda} + \frac{D l_P^2}{4} \left(R_{\kappa\lambda} R^{\kappa\lambda} - \frac{1}{4} R^2 \right) g_{\mu\nu} = \frac{8\pi G}{c^4} \left(T_{\mu\nu} - \frac{1}{4} T g_{\mu\nu} \right)$$

- Equivalence of the derivations Equivalence of entropies?
- Traceless equations Unimodular Gravity
- Nonlinear in second derivatives
- General equations
- Further understanding of the equations
 - Thermodynamics for Weyl transverse gravity

[A.A-S, L.J. Garay, M. Liska, PRD 106 (2022) 6, 064024 and CQG 40 (2023) 2, 025012]

• Extra terms due to higher-order corrections?

SKETCH OF A COSMOLOGICAL MODEL

[A.A-S, M. Liska, A. Vicente-Becerril, PLB 839 (2023) 137827]

Consider a homogeneous, isotropic cosmological model (FLRW model)

$$ds^{2} = -c^{2}dt^{2} + a^{2}(t)\left(\frac{dr^{2}}{1 - kr^{2}} + r^{2}\left(d\theta^{2} + \sin^{2}\theta d\phi^{2}\right)\right)$$
Matter content: Perfect fluid
Raychaudhuri equation: $\dot{H} - \frac{k}{a^{2}} - \frac{Dl_{P}^{2}}{c^{2}}\left(\dot{H} - \frac{k}{a^{2}}\right)^{2} = -4\pi G\left(\rho + \frac{p}{c^{2}}\right)$
Energy conservation
Perturbative Friedmann equation
$$H^{2} = \frac{8\pi G\rho}{3}\left(1 - \frac{2\pi D\rho}{\rho_{P}}\right) + \tilde{\Lambda}$$
• D>0 Avoid singularity
(similar to the bounce in LQC)

DISCUSSION

- Understanding thermodynamics of spacetime
 - Assumptions, principles and implications
 - Equivalence of entropies
 - Unimodular Gravity
- > Analyzing general low-energy quantum gravity effects
 - General quantum phenomenological dynamics from logarithmic corrections to entropy
 - Effects in a specific models: Cosmological model

FUTURE PERSPECTIVES

- Full structure of phenomenological equations
- Could we find the action which implies the equations of motion?
- Study of explicit solutions
- Constraints on parameter D
- Equivalence of entropies
- Structure of gravity: Weyl transverse gravity

Thank you very much for your attention!

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