

Energy, Entropy, and the Intimate Relations between the Two in Semi-Classical Gravity

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Apology

Άπολογία

in physics we feel, I think justifiably, that we have come to learn something about the world, often something concrete, sometimes something deep

this comes out most clearly, I think, when we keep in mind that physics does more than predict experimental outcomes based on clearly formulated mathematical models

it also teaches us about qualitative features of the world that we do not know how to model in anything like an adequate quantitative sense (turbulence, *e.g.*)

and it teaches us about broad and global features of the world, about its *possible* behaviors, that, it seems, we need general theorems to characterize (the relationships among topological, causal and affine structures in GR, *e.g.*, exemplified by the classic singularity theorems)

in philosophy, we try to understand what it is we've learned, and we realize that learning is only the first step in coming to understand

in reflecting on the state of our knowledge, we recognize that there always remain open questions about that knowledge. . .

- to clarify, elaborate and enrich the concepts we (are trying to!) use to formulate and represent the knowledge physics has given us
- to use that knowledge to put these concepts into cogent, fruitful and often new relation to one another, and sometimes to develop new concepts entirely
- to grasp what conceptual possibilities are opened up or closed off more generally by that knowledge

in the best of cases, this sets up a
self-sustaining feedback loop, a virtuous
epistemic circle

physics provides philosophy the knowledge to reflect on:

[W]e are met as cultivators of mathematics and physics. In our daily work we are led up to questions the same in kind with those of metaphysics; and we approach them, not trusting to the native penetrating power of our own minds, but trained by a long-continued adjustment of our modes of thought to the facts of external nature.

– James Clerk Maxwell
“Address to the Mathematical and Physical
Sections of the British Association” (1870)

philosophy provides physics the questions and conceptual possibilities whose investigation may not lead to definitive answers, but pleasantly sometimes opens up new avenues of research that lead us to more and deeper learning about the world:

[W]e must bear in mind that the scientific or science-producing value of the efforts made to answer these old standing questions is not to be measured by the prospect they afford us of ultimately obtaining a solution, but by their effect in stimulating men to a thorough investigation of nature. To propose a scientific question presupposes scientific knowledge, and the questions which exercise men's minds in the present state of science may very likely be such that a little more knowledge would shew us that no answer is possible. The scientific value of the question, How do bodies act on one another at a distance? is to be found in the stimulus it has given to investigations into the properties of the intervening medium.

James Clerk Maxwell

"Attraction" (*Encyclopædia Britannica*, edition IX, 1875)

in a field such as black hole thermodynamics, and semi-classical gravity more generally—where we have not only no empirical experience to test our theorizing, but, much more importantly (and worse), we have none to *guide* and *constrain* it. . .

where we have not been “trained by a long-continued adjustment of our modes of thought to the facts of external nature” . . .

I see no clear line to be drawn to demarcate physics from philosophy

And so my task here today, as this Socratic ἀπολογία suggests, is to play Socratic gad-fly

Outline

Remembrance of Things Past—Maxwell's Way:
Energy and Entropy in Classical Thermodynamics

In the Shadow of Old Energy and Entropy Conditions
Blooming in GR

Wall's Way—Sodom and Gomorrah

Remembrance of Things Past—Maxwell's Way: Energy and Entropy in Classical Thermodynamics

In the Shadow of Old Energy and Entropy Conditions
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as a philosopher and gad-fly one of my duties is to take a step back and remind us all from time to time of some basic ideas that we tend to lose sight of, so much part of the ground of our framing conceptual landscape are they

so I begin with a few basic—almost trite—observations about energy and entropy in classical physics (non-relativistic, non-quantum), to set the stage for the discussion of more advanced, abstract matters, far removed from our empirical access

in any event, one never goes wrong by invoking Fermi's Pleasure Principle to begin a talk

what energy and entropy share in common in classical physics:

1. their **universality**—every theory treating every kind of physical system attributes (or can attribute) both to them, except perhaps entropy for the simple mechanics of rigid bodies
2. and except—*most interestingly*—the Newtonian gravitational field (if one tries to attribute a “Bekenstein entropy” to a “Newtonian black hole”, it doesn’t work for many reasons, among them that the geometry of Newtonian spacetime does not admit event horizons, nor even trapping horizons)
3. their **fungibility**—all forms of each are, respectively, transformable into each other
4. both are **affine** quantities (no natural zero point)
5. each obeys a characteristic relation among all their possible forms, for energy an equality (**conservation**) and for entropy an inequality (**the Second Law**), and both of those relations have an intimate connection with time

what energy and entropy share in common in classical physics (cont.):

6. indeed, there are (we tend to forget) two entirely **separate conceptions** of energy in classical physics, the **generator of time translations** (“Hamiltonian”) and the **capacity to do work** (“Lagrangian”)
7. that they are *a priori independent* of each other is seen most clearly by the fact that, since work depends only on spatial interval, the latter makes not even an implicit reference to time¹

1. the work 1-form in Lagrangian mechanics is orthogonal to “accelerations”; one cannot represent “work”, much less dissipative processes, in Hamiltonian mechanics

what energy and entropy share in common in classical physics (cont.):

8. the two conceptions of energy are brought into contact with each other through the principle of conservation, which itself is balanced by the introduction of the idea of *heat* (neither a generator of time translations nor the capacity to do work) and thus, inevitably, the introduction of *entropy from energetic considerations*
9. thus there are in fact *three conceptions of energy* in classical physics, *united by the conservation principle and by the Second Law* (for the latter, most explicitly in the form of the Kelvin Postulate)
10. and, tantalizingly, *by the characterization of equilibrium*, as the maximizer of entropy and the minimizer of free energy (Kelvin's "Principle of Dissipated Energy")

Now when the appearance of one thing is strictly connected with the appearance of another, so that the amount which exists of the one thing depends on and can be calculated from the amount of the other which has disappeared, we conclude that the one has been formed at the expense of the other, and that they are both forms of the same thing.

– James Clerk Maxwell
Theory of Heat (1891, ch. IV, p. 93)

note how careful Maxwell's formulation is: it applies *both* to energy (conservation, an equality) *and* to entropy (Second Law, an inequality)

nonetheless, there are clear conceptual and physical differences, most fundamentally:

1. they are not jointly fungible
2. there is no such thing as an entropometer, but there are ergometers (“entropy mediates no physical couplings between systems”—it is a purely *modal* quantity)
3. energy conservation must be postulated, but plays essential roles in derivation of many important, general results
4. entropy increase can be derived, but plays no essential role in derivation of any important, general results (“the Second Law lacks fecundity”)
5. entropy constrains the possible future transformations of a system more severely than does energy conservation

Emden (1938, p. 909) (see rest of article as well):

As a student, I read with advantage a small book by F. Wald entitled 'The Mistress of the World and her Shadow'. These meant energy and entropy. In the course of advancing knowledge the two seem to me to have exchanged places. In the huge manufactory of natural processes, the principle of entropy occupies the position of manager, for it dictates the manner and method of the whole business, whilst the principle of energy merely does the bookkeeping, balancing credits and debits.

Remembrance of Things Past—Maxwell's Way: Energy and Entropy in Classical Thermodynamics

**In the Shadow of Old Energy and Entropy Conditions
Blooming in GR**

Wall's Way—Sodom and Gomorrah

many properties of energy carry over from classical physics, and are even deepened and become more subtle, for they are now united (along with momental quantities) in a single new quantity, *viz.*, stress-energy:

universality matter couples with “gravity” (“the universal force”, *viz.*, spacetime curvature) by way of stress-energy, attributed promiscuously to every kind of matter field

fungibility every form of stress-energy can be transformed (in principle) into every other kind

conservation one now has in general only “hyperlocal” conservation ($\nabla^n T_{na} = 0$), no integral conservation laws (“gravitational energy is non-local”)

zero point no longer an affine quantity (“ $T_{ab} = 0$ ” is unambiguous)

and energy conditions make their first appearance

- in classical GR, as in most theories, one has a great deal of freedom in what one takes as primitive and what as derived
- think of the geodesic principle and covariant conservation of stress-energy, inter-derivable
- this is not true of the classical energy conditions, neither pointilliste nor impressionist
- one *can't derive energy conditions* in classical GR
- they are always taken as primitive (Curiel 2017)
- perhaps they reach down to and get ahold of spacetime structure at a very deep level?

- this is **not so of entropy conditions** in classical GR²
- Bousso (1999a, 1999b), e.g., used the DEC in his original work to motivate his covariant entropy bound (“the total entropy flux S_L through any null hypersurface L satisfying some natural geometrical conditions must be such that $S_L \leq A/4$, where A is a spatial area canonically associated with L ”)
- Flanagan, Marolf, and Wald (2000) then proved it using the NEC

2. Putting aside for the moment what one means by entropy, here, what kind of entropy one is dealing with—just assume for the moment that “black hole area is something like an entropy”.

the relations between energy and entropy are neighborly, but not intimate:

1. relation between energy conditions and entropy conditions is “one way” only
2. there is no explicit unification of different types of entropy (in something like a GSL, *e.g.*), as there is for energy
3. as in classical thermodynamics—if one accepts the orthodox dogma (Wald’s Way), that there is no consistent thermodynamical theory of purely classical black holes—energy and entropy are not jointly fungible (throwing mass into a classical black hole doesn’t increase its entropy)
4. there is still no entropometer
5. and relation of *both* stress-energy and entropy to equilibrium (existence of timelike Killing field) is obscure at best

in any event, already energy here goes beyond the role it plays in non-relativistic physics. . .

to paraphrase Emden's marvelous remark, in the huge manufactory of natural processes, energy begins to occupy the position of, if not manager, at least assistant manager, for it constrains the manner and method of the whole business

Carlo Rovelli (personal communication, his emphases):

Entropy (and thermodynamics in general) is not yet understood (by anybody) even in CLASSICAL gravity. . . . What is well understood (since Tolman) is the thermodynamics of matter fields on a given geometry, which is to say: disregarding dissipation into gravitational degrees of freedom. But dissipation into gravitational degrees of freedom has no reason not to occur (which in the Black Hole case is the classical analog of your observation about gravity getting entangled with the Hawking radiation, for instance). So the classical thermodynamics of the gravitational field is a topic for which we do not have a science at all, yet.

I agree

Maybe we can make some progress in SCG.

Remembrance of Things Past—Maxwell's Way: Energy and Entropy in Classical Thermodynamics

In the Shadow of Old Energy and Entropy Conditions
Blooming in GR

Wall's Way—Sodom and Gomorrah

How did classical general relativity know that the horizon area would turn out to be a form of entropy, and that surface gravity is a temperature?

Ted Jacobson
“Thermodynamics of Spacetime:
The Einstein Equation of State”

Recall the SCEFE, the basis of SCG:

$$G_{ab} = 8\pi \langle \hat{T}_{ab} \rangle_{psi}$$

can it shed light on why classical GR already seems to know about the thermodynamics of gravity? if so, how?

Semi-classical gravity is remarkably powerful. It is far more powerful than it has any right to be. . . . Gravity knows stuff. It knows about its own quantum states, it knows about the quantum states of matter.

Raphael Bousso

talk at conference

“Energy Conditions in Quantum Field Theory”

(Leipzig, 13. Sep 2022)

pace Bousso, I do not want to assume any kind of holography (including AdS/CFT) to try to address these issues

SCG in its standard, flat-footed formulations seems to me already parlous enough to assume, in our current epistemic state

first hints of more intimate relations between energy and entropy,³ provisionally accepting BHT:

fungibility energy can now be directly transformed into entropy (“throw stuff into black hole”), and vice-versa (Hawking radiation); each is a direct measure of the other (“area and mass tell you each other”)

zero point they both have natural zero points, which is the same state (Schwarzschild $M = 0$)

equilibrium heuristic but compelling arguments that stationary black holes minimize free energy (“ $M - M_{\text{irr}}$ ”) and maximize entropy

entropometer! we can measure area of event horizon directly (not a modal quantity)—which is also, in this case an ergometer, as area gives you mass

3. Here, we take the black hole area to be uncontroversially (more or less) a thermodynamical entropy.

but the relations—including equalities—between entropy and energy even beyond those just stated become now a promiscuous, libertine, orgiastic debauch. . .

Sodom and Gomorrah, during the fun times!

entropy conditions⁴ take on the classical role of energy conditions:

1. GSL proves a singularity theorem, and rules out traversable wormholes, negative masses, other forms of faster-than-light travel between asymptotic regions, restarting inflation and CTCs: Wall (2013)
2. Bousso bound proves a singularity theorem: Bousso and Shahbazi-Moghaddam (2022)
3. quantum Penrose inequality using generalized entropy of bulk light sheets to constrain lower bound of ADM mass: Bousso, Shahbazi-Moghaddam, and Tomašević (2019)
4. quantum focusing implies singularity theorems, the GSL and boundary causality: Shahbazi-Moghaddam (2022)

(*N.b.*: the last only a conjecture, with supporting plausibility arguments and evidence from test cases)

4. It now becomes almost wholly unclear what is meant by entropy in any given application, and whether, in any event, they are jointly consistent.

*the principle of entropy increase (GSL) becomes
fecund!*

one now has, for the first time, not only derivations of energy conditions, but ones *based on entropy conditions*, and vice-versa

1. proving the (A)ANEC from the GSL: Wall (2010)
2. proving the ANEC from the QNEC: Bousso et al. (2016)
3. proving the NEC from the GSL: Parikh and Svesko (2017)

(we'll come to the QNEC—Quantum Null Energy Condition—shortly)

QNEC, Bousso et al. (2016)

1. any point p and null vector k^a define (at least locally) a null plane N
2. given any codimension-2 surface Σ that contains p and lies on N , consider the von Neumann entropy S_{out} of the quantum state of the ambient quantum fields restricted to one side of Σ
3. a second variation S''_{out} can be defined by deforming Σ along N , in a small neighborhood of p , by an area \mathcal{A}
4. QNEC: $\langle T_{kk}(p) \rangle \geq \frac{\hbar}{2\pi} \lim_{\mathcal{A} \rightarrow 0} S''_{\text{out}} / \mathcal{A}$.

one gets equivalences of entropic and energetic quantities:

1. Leichenauer, Levine, and Shahbazi-Moghaddam (2018): for null shape deformations as they appear in the QNEC, modulo a plausible, supported conjecture, second variations of the von Neumann entropy determine the full stress-energy tensor expectation value as an equality (and so, *à la* Jacobson 1995, one gets the EFE)
2. Wang (2020): (quasi-local) Bartnik-Bray inner mass exactly equals the (generally non-local) irreducible mass corresponding to the (generally non-local) outer entropy (Engelhardt and Wall 2018)

perhaps most striking:

1. set

$$K_\xi := \int_\Sigma \langle T_m{}^n \rangle \xi^m d\Sigma_n$$

2. then SCG First Law:

$$\delta M = \frac{\kappa}{8\pi} \delta A + \langle K_\xi \rangle$$

3. invoke First Law of quantum thermodynamics, *a.k.a.*, First Law of entanglement

$$\delta \langle K_\xi \rangle = T_H \delta S_{\text{ent}}$$

where $\delta S_{\text{ent}} = -\text{Tr} \rho \log \rho$ and $\rho = \frac{1}{Z} e^{-\beta_H K_\xi}$

4. $\Rightarrow \delta M = T_H \delta S_{\text{gen}}$ (where $S_{\text{gen}} = S_{\text{BH}} + S_{\text{ent}}$)

S_{gen} obeys *both* a First Law and a Second Law!!!

Da wird der Hund in der Pfanne verrückt.

Now when the appearance of one thing is strictly connected with the appearance of another, so that the amount which exists of the one thing depends on and can be calculated from the amount of the other which has disappeared, we conclude that the one has been formed at the expense of the other, and that they are both forms of the same thing.

– James Clerk Maxwell
The Theory of Heat (ch. iv, p. 93)

Are energy and entropy different aspects, different forms, of the same underlying entity? Should this be one of the unifications we seek now in physics?

two paths to the EFE: (1) entropy

further hints from Jacobson (1995):

1. assume form of entropy (proportional to area)
2. assume form of heat (matter flux of particular sort, = 0 when $T_{ab} = 0$)
3. assume Clausius relation (temperature from Unruh effect)
4. \Rightarrow Einstein field equation as consistency condition on thermodynamical relations (“equation of state”)

\Rightarrow *null energy condition as local law of entropy increase*

(see also Curiel, Finster, and Isidro 2020)

two paths to the EFE: (2) stress-energy

Theorem (Curiel 2019)

The only two covariant-index, symmetric, divergence-free, second-order concomitants of the metric with physical dimension of stress-energy (in geometrized units) are constant multiples of the Einstein tensor.

⇒ physical dimension of stress-energy (T_{ab} invariant under constant recalings of g_{ab}) determines coupling of curvature to matter

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