# Probing the mechanism behind the matter-antimatter asymmetry



### Julia Harz

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Working Seminar "Mathematical Physics", Universität Regensburg



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### What is our × Universe made of? NASA, ESA, and J. Lotz, M. Mountain, A. Koekemoer, and the HFF Team (STScI)



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### The Standard Model of Particle Physics



#### Standard particles

#### + their anti-particles

However, this is cannot be the final word yet ...



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# Why do we exist?



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5 م

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# Why is there more matter than anti-matter?



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### The baryon asymmetry

Equal amounts of matter and antimatter would have annihilated into radiation

As our Universe consists out of matter, a mechanism had to create the asymmetry





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### From the Big Bang to Today...



Credits: University of Cambridge / The Stephen Hawking Centre for Theoretical Cosmology



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### **Big Bang Nucleosynthesis (BBN)**

- 3 min after Big Bang
- BBN creates first light elements (D, He)

#### **Deuterium Bottleneck**

 $\rightarrow\,$  Nucleosynthesis starts with formation of D

$$p + n \rightarrow D + \gamma$$

Only if photo-dissociation ceases to be effective, chain of light elements can be formed

$$\eta_B^{-1} e^{-B_D/T} < 1 \qquad B_D = 0$$
$$T_{\rm nuc}^D \approx \frac{B_D}{\log \eta_B^{-1}} = 0$$

 $T \approx 0.1 \mathrm{MeV}$ 



$$\eta_B^{\rm obs} = (6.143 \pm 0.190) \times 10^{-10}$$

 $= 2.3 \mathrm{MeV}$ 



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### From the Big Bang to Today...



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### The Cosmic Microwave Background (CMB)

- 400.000 years after Big Bang
- measures temperature fluctuations from recombination









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### The Cosmic Microwave Background (CMB)

- 400.000 years after Big Bang
- measures temperature fluctuations from recombination



$$\Omega_b^{\rm obs} h^2 = 0.0224 \pm 0.0001$$





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12

### **Combination: BBN & CMB**



Excellent agreement even though measurements originate from two different epochs!



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### Ingredients for a baryon asymmetry

**Theoretical conditions** that have to be fulfilled (Sakharov conditions):





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### **B-L violation & Sphaleron Processes**

$$J^B_{\mu} = \frac{1}{3} \sum_{i} \left( \overline{q}_{Li} \gamma_{\mu} q_{Li} - \overline{u}^c_i \gamma_{\mu} u^c_i - \overline{d}^c_i \gamma_{\mu} d^c_i \right)$$
$$J^L_{\mu} = \sum_{i} \left( \overline{\ell}_{Li} \gamma_{\mu} \ell_{Li} - \overline{e}^c_i \gamma_{\mu} e^c_i \right)$$

 $\begin{array}{c} s_L \\ c_L \\ c_L \\ d_L \\ d_L \\ u_L \\ u_L \\ \nu_e \end{array} \begin{array}{c} s_L \\ b_L \\ b_L \\ b_L \\ \nu_{\tau} \\ \nu_{\mu} \end{array}$ 

 $\Delta L = \Delta B = 3$ 

Naively expected on classical level:  $\partial^{\mu}J^{B}_{\mu} = \partial^{\mu}J^{L}_{\mu} = 0$ 

Charge operators:

$$\hat{L} = \int d^3x J_0^L(x) \qquad \qquad \hat{B} = \int d^3x J_0^B(x)$$

#### Anomalies on quantum level due to chiral structure of SM:

$$\begin{split} \partial_{\mu}J^{B\mu} &= \partial_{\mu}J^{L\mu} = \frac{N_{f}}{32\pi}(g^{2}W^{a}_{\mu\nu}\tilde{W}^{a\mu\nu} - g'^{2}B_{\mu\nu}\tilde{B}_{\mu\nu}) = N_{f}\partial_{\mu}K^{\mu} \\ \Delta B &= \Delta L = N_{f}\int d^{4}x\partial_{\mu}K^{\mu} = N_{f}n_{\text{top}} \\ \\ \partial_{\mu}(J^{B\mu} - J^{L\mu}) = 0 & \text{B-L conserved} \\ \partial_{\mu}(J^{B\mu} + J^{L\mu}) \neq 0 & \text{B+L violated} \end{split}$$



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### **B-L violation & Sphaleron Processes**





$$\Delta B = \Delta L = N_f \int d^4 x \partial_\mu K^\mu = N_f n_{\rm top}$$

→ Sphalerons are active at temperatures above EWSB! (around 170 GeV)

$$\frac{\Gamma_{\rm SM}^b}{V} \sim \exp\left(-\frac{4\pi}{g_w}\frac{v_c}{T}\right) \qquad \begin{array}{l} \text{in broken phase} \\ \text{suppressed} \\ \\ \frac{\Gamma_{\rm SM}^s}{V} \sim \alpha_w^5 T^4 \qquad \begin{array}{l} \text{in symmetric} \\ \text{phase active} \end{array}$$



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### Ingredients for a baryon asymmetry

**Theoretical conditions** that have to be fulfilled (Sakharov conditions):





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### C and CP violation

Charge conservation implies:

$$\Gamma(X \to AB) = \Gamma(\overline{X} \to \overline{A} \ \overline{B})$$

**Requirement of charge violation:** 

$$\frac{dY_B}{dt} \approx \Gamma(X \to AB) - \Gamma(\overline{X} \to \overline{A} \ \overline{B})$$

Charge and parity conservation implies:

$$\Gamma(X \to q_L q_L) = \Gamma(\overline{X} \to \overline{q}_R \ \overline{q}_R)$$
$$\Gamma(X \to q_R q_R) = \Gamma(\overline{X} \to \overline{q}_L \ \overline{q}_L)$$

Requirement of charge and parity violation:

$$\frac{dY_B}{dt} \approx \left[ (\Gamma(\overline{X} \to \overline{q}_R \ \overline{q}_R) + \Gamma(\overline{X} \to \overline{q}_L \ \overline{q}_L)) - (\Gamma(X \to q_R q_R) + \Gamma(X \to q_L q_L)) \right]$$



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muon

tau

electron

### C and CP violation – Status of the SM?

$$CP | K^0 \rangle = - | \bar{K}^0 \rangle = - | s\bar{d} \rangle$$
$$CP | \bar{K}^0 \rangle = - | K^0 \rangle = - | d\bar{s} \rangle$$

 $CP |K_1^0\rangle = CP \left(\frac{1}{\sqrt{2}} \left(|K^0\rangle - |\bar{K}^0\rangle\right)\right) = + |K_1^0\rangle$ 

 $CP | K_2^0 \rangle = CP \left( \frac{1}{\sqrt{2}} \left( | K^0 \rangle + | \bar{K}^0 \rangle \right) \right) = - | K_2^0 \rangle$ 

**Expected decays:** 

$$K_1^0 
ightarrow 2\pi$$
 fast $K_2^0 
ightarrow 3\pi$  slow

**BUT:** In 1964, Cronin and Fitch realised that physical states are no pure CP eigenstates!

$$|K_{S}^{0}\rangle = \frac{1}{\sqrt{1+\left|\epsilon\right|^{2}}}\left(\left|K_{1}^{0}\rangle + \epsilon\left|\bar{K}_{2}^{0}\rangle\right)\right) \qquad |K_{L}^{0}\rangle = \frac{1}{\sqrt{1+\left|\epsilon\right|^{2}}}\left(\left|K_{2}^{0}\rangle + \epsilon\left|\bar{K}_{1}^{0}\rangle\right)\right)$$
Indirect CP violation

$$\frac{J_{CP}}{T_C^{12}} \approx 10^{-20} \quad \clubsuit \quad \mathcal{O}(10^{-10})$$

not enough CP violation within SM!



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### Ingredients for a baryon asymmetry

**Theoretical conditions** that have to be fulfilled (Sakharov conditions):





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### Departure from thermal equilibrium

$$\langle B \rangle_T = \text{Tr}[e^{-\beta H}B] = \text{Tr}[(CPT)(CPT)^{-1}e^{-\beta H}B]$$
  
=  $\text{Tr}[e^{-\beta H}(CPT)^{-1}B(CPT)] = -\langle B \rangle_T$ 

#### Departure from thermal equilibrium:

- First order phase transition (FOPT)
- Out-of-equilibrium decays

#### Strong FOPT during EWSB:

$$\frac{v_c}{T_c} \simeq \frac{3g^3v^2}{32\pi^2 m_h^2} \ge 1$$
$$m_h \le 32 \text{GeV}$$



#### → Higgs too heavy for first order phase transition



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### Why do we need new physics?

**Theoretical conditions** that have to be fulfilled (Sakharov conditions):

#### Standard Model?



#### There has to be new physics in order to explain our own existence!



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## Do we understand all known matter?



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# What is the nature of neutrinos?



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Neutrino **oscillations** require **massive** neutrinos, forbidden in the Standard Model.

#### How do neutrinos get their masses?



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### Neutrinos – what do we know?

• Neutrinos in the Standard Model are **massless** 

$$L_i \to \left(\begin{array}{c} \nu_i \\ \ell_i \end{array}\right) \qquad \qquad m_\nu = 0$$

• Neutrino **mixing** 

$$\left(\begin{array}{c}\nu_e\\\nu_\mu\\\nu_\tau\end{array}\right) = U_{PMNS} \left(\begin{array}{c}\nu_1\\\nu_2\\\nu_3\end{array}\right)$$

• Neutrino oscillations require massive neutrinos

$$P(\nu_i \to \nu_j) \propto \Delta m_{ij}^2 \qquad \frac{\Delta m_{12}^2 \sim 7.59 \times 10^{-5} \text{eV}^2}{\Delta m_{23}^2 \sim \Delta m_{31}^2 \sim 2.3 \times 10^{-3} \text{eV}^2}$$

• Normal vs. inverted hierarchy

#### How do neutrinos get their masses? What nature do neutrinos have? Are they their own anti-particles?



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The discovery of these oscillations shows that neutrinos have mass.





### **Right-handed Neutrinos as New Physics**



#### **Right-handed neutrinos could explain the neutrino masses**



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### Neutrino mass mechanism



 $\begin{array}{c} 1/2\\ y_{\nu}L\epsilon H\overline{\nu}_R \supset m_D\nu_L\overline{\nu}_R\\ -1/2 \quad 0 \end{array}$ 

$$m_{\nu} = \frac{m_D^2}{m_M}$$

hypercharge

*T*I

 $m_M \overline{\nu}_R \nu_R^c$ 

Majorana mass

tiny Yukawa couplings

 $m_{\nu}/\Lambda_{EW} \le 10^{-12}$ 

→ lepton number no accidental symmetry anymore

→ higher dimensional operator

 $\frac{1/2}{m_M} \frac{1/2}{\overline{
u}_L} \nu_L^c$ 

not at tree-level within the SM possible

#### → Lepton number violation (LNV)





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### SeeSaw mechanism

 $\mathcal{L} \supset Y_{\nu} L \overline{\nu}_R \phi + M_M \overline{\nu}_R \nu_R^c$  $= M_D \nu_L \overline{\nu}_R + M_M \overline{\nu}_R \nu_R^c$ 

$$\mathbf{M}_{\nu} = \begin{pmatrix} \delta m_{\nu}^{1loop} & M_D \\ M_D^T & M_M \end{pmatrix}$$

#### Diagonalisation with $M_M >> M_D$ :

$\nu \simeq \nu_L + \theta \nu_R^c$	$m_{\nu} \simeq \frac{M_D^2}{M_M}$	mainly <b>active</b> SU(2) <sub>L</sub> doublet states with <b>light</b> masses	Seesaw type I
$N \simeq \nu_R + \theta^T \nu_L^c$	$m_N \simeq M_M$	mainly <b>sterile</b> singlet states with <b>heavy</b> masses	

#### "naturally" small neutrino masses



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## Baryogenesis / Leptogenesis.



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### Basic principle of standard baryogenesis





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### Basic principle of standard leptogenesis





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### Basic principle of standard leptogenesis





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### Leptogenesis.



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### **Overview of leptogenesis models**





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### Leptogenesis & the neutrino mass mechanism



#### Combined analysis of both regimes and comparison with existing literature (Klaric et al. 2021)



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# High-scale leptogenesis I

- Generation of lepton asymmetry via heavy neutrino decays with sources of CP violation
- **Competition** with lepton number violating (LNV) **washout** processes
- **Conversion** to a baryon asymmetry via **sphaleron** processes



$$\varepsilon_1 \equiv \frac{\Gamma(N_1 \to LH) - \Gamma(N_1 \to \overline{L} \,\overline{H})}{\Gamma(N_1 \to LH) + \Gamma(N_1 \to \overline{L} \,\overline{H})} \approx -\frac{3}{8\pi} \frac{1}{(hh^{\dagger})_{11}} \sum_{i=2,3} \mathrm{Im}(hh^{\dagger})_{li}^2 \frac{M_1}{M_i}$$

Fukugita et al. (1986)





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# High-scale leptogenesis II

- Generation of lepton asymmetry via heavy neutrino decays with sources of CP violation
- **Competition** with lepton number violating (LNV) **washout** processes
- Conversion to a baryon asymmetry via sphaleron processes





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Probing the mechanism behind the matter-antimatter asymmetry Fukugita et al. (1986)



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# High-scale leptogenesis III

- **Generation** of lepton asymmetry via **heavy neutrino decays** with sources of **CP violation**
- **Competition** with lepton number violating (LNV) **washout** processes
- Conversion to a baryon asymmetry via sphaleron processes

$$B = \sum_{i} (2\mu_{q_i} + \mu_{u_i} + \mu_{d_i})$$

- ELH Yukawa:
- DQH Yukawa: 0
- UQH Yukawa:
- QQQL sphalerons:
- No electric charge:

$$0 = \mu_E + \mu_L + \mu_H$$

$$0 = \mu_D + \mu_Q + \mu_H$$

$$0 = \mu_U + \mu_Q - \mu_H$$

 $0 = 3u_0 + u_1$ 

$$0 = N_{\mu\nu}(\mu_{\rho} - 2\mu_{\mu} + \mu_{\rho} - \mu_{\mu} + \mu_{F}) -$$

$$L = \sum_{i} (2\mu_{\ell_i} + \mu_{e_i})$$

$$0 = N_{\rm gen}(\mu_Q - 2\mu_U + \mu_D - \mu_L + \mu_E) - 2N_{\rm Higgs}\mu_H$$

$$B = \frac{8N_f + 4m}{22N_f + 13m}(B - L)$$

Fukugita et al. (1986)





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# High-scale leptogenesis IV

- **Generation** of lepton asymmetry via **heavy neutrino decays** with sources of **CP violation**
- **Competition** with lepton number violating (LNV) **washout** processes
- Conversion to a baryon asymmetry via sphaleron processes





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## High-scale leptogenesis – models

#### Plethora of examples:

- Extension of seesaw type-I by new scalars with same quantum numbers as SM fermions → e.g. long-lived scalars, R-hadrons, heavy sterile neutrinos e.g. Fong et al. (2013)
- **Z' models**  $\rightarrow$  same-sign di-lepton final states e.g. Chun (2005)
- Left-right symmetric models  $\rightarrow$  falsification by low mass  $W_{R}^{}$  e.g. Dev. et al. (2015)
- Soft leptogenesis → type-I: charged LFV e.g. Adhikari et al. (2015)
   → type-II: same-sign di-lepton resonance, same-sign tetra-leptons
   e.g. Chun et al. (2006)

See review "Probing Leptogenesis", JH, Chun et al. (arxiv:hep-ph/1711.02865)

Generic problem: new physics at high scales



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# Basic principle of leptogenesis





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# Basic principle of leptogenesis



# Strategy: Search for washout processes with the potential to falsify models!



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#### **Lepton Number Violation**

#### LNV occurs only at odd mass dimension beyond dim-4:



#### See surveys of all LNV operators up to dim-11 e.g. in

Babu, Leung (2001), Gouvea, Jenkins (2008), Graf, JH, Deppisch, Huang (2018)



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## **Probing LNV washout interactions**





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# Falsifying baryogenesis with LHC & 0vββ decay

Observation of any LNV washout process at the LHC would falsify high-scale baryogenesis

Deppisch, JH, Hirsch (2014)





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### Probing Leptogenesis at the LHC

#### Washout processes could be observable at the LHC



#### Observation of any washout process at LHC would falsify high scale baryogenesis!

(scale of asymmetry generation *above* M<sub>x</sub>)



Deppisch, JH, Hirsch, Phys. Rev. Lett. (2014) Deppisch, JH, Hirsch, Päs, Int. J. Mod. Phys. A (2015)



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47

### Probing Leptogenesis at the LHC

#### Washout processes could be observable at the LHC



Deppisch, JH, Hirsch, Phys. Rev. Lett. (2014) Deppisch, JH, Hirsch, Päs, Int. J. Mod. Phys. A (2015)



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# Falsifying baryogenesis with LHC & 0vββ decay

Observation of any LNV washout process at the LHC would falsify high-scale baryogenesis

Deppisch, JH, Hirsch (2014)



Observation of neutrinoless double beta decay with new physics from > dim-5 LNV operators would falsify high-scale baryogenesis

Deppisch, Graf, JH, Huang (2018) Deppisch, JH, Huang, Hirsch, Päs (2015)





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#### Half life of Neutrinoless Double Beta Decay





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### Probing Leptogenesis with $0\nu\beta\beta$



$$T_{1/2}^{-1} = G_{0\nu} |\mathcal{M}|^2 |\epsilon_{\alpha}^{\beta}|^2$$

Observation would fix the **effective coupling** for one operator

$\mathcal{O}$	Operator
$1^{H^2}$	$L^i L^j H^k H^l \overline{H}^t H_t \epsilon_{ik} \epsilon_{jl}$
2	$L^i L^j L^k e^c H^l \epsilon_{ij} \epsilon_{kl}$
$3_a$	$L^i L^j Q^k d^c H^l \epsilon_{ij} \epsilon_{kl}$
$3_b$	$L^i L^j Q^k d^c H^l \epsilon_{ik} \epsilon_{jl}$
$4_a$	$L^i L^j \overline{Q}_i \bar{u^c} H^k \epsilon_{jk}$
$4_b^{\dagger}$	$L^i L^j \overline{Q}_k u^{\overline{c}} H^k \epsilon_{ij}$
8	$L^i \bar{e^c} \bar{u^c} d^c H^j \epsilon_{ij}$

$$\frac{G_F \epsilon_7}{\sqrt{2}} = \frac{g^3 v}{2\Lambda_7^3}$$

effective coupling can be related to the scale of the operator

$\mathcal{O}_D$	$\Lambda_D^0$ [GeV]
$\mathcal{O}_5$	$9.1 \times 10^{13}$
$\mathcal{O}_7$	$2.6  imes 10^4$
$\mathcal{O}_9$	$2.1 \times 10^3$
$\mathcal{O}_{11}$	$1.0  imes 10^3$



$$\Lambda_7 \left(\frac{\Lambda_7}{c_7' \Lambda_{Pl}}\right)^{\frac{1}{5}} \lambda_7 < T < \Lambda_7$$

Limit above which the washout is highly effective can be calculated in dependence of the operator scale

Deppisch, Graf, JH, Huang (2018) Deppisch, JH, Huang, Hirsch, Päs (2015)

 $\frac{\Gamma_W}{H} > 1$ 



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Emmy Noether

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#### Probing Leptogenesis with $0\nu\beta\beta$



#### Potential to falsify baryogenesis models!



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Probing the mechanism behind the matter-antimatter asymmetry Deppisch, Graf, JH, Huang (2018) Deppisch, JH, Huang, Hirsch, Päs (2015)



# Falsifying baryogenesis with LHC & 0vββ decay

Observation of any LNV washout process at the LHC would falsify high-scale baryogenesis

Deppisch, JH, Hirsch (2014)

Observation of neutrinoless double beta decay with new physics from > dim-5 LNV operators would falsify high-scale baryogenesis

Deppisch, Graf, JH, Huang (2018) Deppisch, JH, Huang, Hirsch, Päs (2015)

#### Caveats might apply, e.g.:

- Flavor specific leptogenesis
- Dark U(1) symmetries

Aristizabal Sierra, Fong, Nardi, Peinado (2014) Frandsen, Hagedorn, Huang, Molinaro, Päs (2018)

#### How do LHC and 0vββ decay compare wrt their strength to falsify leptogenesis? Does TeV-scale LNV really render standard thermal leptogenesis invalid?



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## A simplified model study of TeV scale LNV

#### UV realization of dim-9 operator:

$$L_{LNV}^{eff} = \frac{C_1}{\Lambda^5} \bar{Q} \tau^+ d\bar{Q} \tau^+ d\bar{L} L^C + \text{h.c.}$$



JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



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**TeV-scale LNV** 

"washout" interactions

## A simplified model study of TeV scale LNV

#### UV realization of dim-9 operator:

#### TeV-scale LNV "washout" interactions

Integrating out heavy d.o.f. leads to dim-9 LNV operator:

$$L_{LNV}^{eff} = \frac{C_1}{\Lambda^5} \bar{Q} \tau^+ d\bar{Q} \tau^+ d\bar{L}L^C + \text{h.c.}$$



Right-handed neutrino interactions ("standard thermal LG"):  $\mathcal{L} \supset y_{\nu} \bar{L}HN - \frac{m_N}{2} \bar{N}^c N + \text{h.c.}$  high-scale source of lepton asymmetry

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



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## A simplified model study of TeV scale LNV

#### UV realization of dim-9 operator:

TeV-scale LNV "washout" interactions

Integrating out heavy d.o.f. leads to dim-9 LNV operator:

$$L_{LNV}^{eff} = \frac{C_1}{\Lambda^5} \bar{Q} \tau^+ d\bar{Q} \tau^+ d\bar{L}L^C + \text{h.c.}$$



Right-handed neutrino interactions ("standard thermal LG"):

$$\mathcal{L} \supset y_{\nu} \bar{L} H N - \frac{m_N}{2} \bar{N}^c N + \text{h.c}$$

high-scale source of lepton asymmetry

#### Can TeV-scale LNV destroy the generated asymmetry from standard thermal LG?

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



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$$\frac{dY_N}{dz} = -\left(D + S\right)\left(Y_N - Y_N^{eq}\right)$$
$$\frac{dY_{B-L}}{dz} = -\epsilon D\left(Y_N - Y_N^{eq}\right) - WY_{B-L}$$

 $\Delta L = 1$ 







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JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



 $\mathcal{O}(m_S) \approx \mathcal{O}(m_F) \approx \mathcal{O}(\text{TeV})$ 



### Low-scale LNV destroys lepton asymmetry previously generated by standard LG scenario.

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)





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Lepton asymmetry from standard thermal LG is washed out below observable value.

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)





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Lepton asymmetry from standard thermal LG is washed out below observable value.

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)





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#### **Reach at colliders**



Case	Mass hierarchy		Process	
C1	$m_S < m_F$	$pp \to e^{\pm}F,$	$F \to e^{\pm}S^{\mp},$	$S^{\mp} \to jj$
C2	$m_S = m_F$	$pp \to e^{\pm}F,$	$F \rightarrow e^{\pm} j j$	
C3	$m_S > m_F$	$pp \to S^{\pm},$	$S^{\pm} \rightarrow e^{\pm}F,$	$F \to e^{\pm} j j$

**Signal generation:** Madgraph + Pythia 8 + Delphes

#### Background:

- SM processes with same-sign leptons (e.g. jjWW)
- Charge misidentification
- Jet-fake leptons from heavy flavour decays

#### S/B discrimination:

neural network

Emmy Noether-Programm DFG Protogrammercal

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JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



#### Reach at 0vßß decay experiments



$$\frac{C_{\text{eff}}}{2\Lambda^5} \left( \bar{q}_L \tau^+ q_R \ \bar{q}_L \tau^+ q_R + \bar{q}_R \tau^+ q_L \ \bar{q}_R \tau^+ q_L \right) \bar{e}_L e_R^c + \text{h.c.}$$

$$\frac{1}{\sqrt{2\Lambda^5}} \frac{C_{\text{eff}} \Lambda_H^2 F_\pi^2}{\pi^- \pi^- \bar{e}_L e_R^c} + \text{h.c.}$$

$$\frac{1}{T_{1/2}} = |M_0|^2 \left[ G_{0\nu} \times (1 \,\mathrm{TeV})^2 \right] \left( \frac{\Lambda_H}{\mathrm{TeV}} \right)^4 \left( \frac{1}{144} \right) \\ \times \left( \frac{v}{\mathrm{TeV}} \right)^8 \left( \frac{1}{\cos \theta_C} \right)^4 \left[ \frac{C_{\mathrm{eff}}^2}{(\Lambda/\mathrm{TeV})^{10}} \right]$$

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



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### Combined results: Leptogenesis, LHC & 0vββ decay



# Comprehensive analysis confirms EFT results and demonstrates interesting interplay between collider and 0vßß reach.

JH, Ramsey-Musolf, Shen, Urrutia-Quiroga (2021)



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# Baryogenesis.



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### **Electroweak baryogenesis**



#### Unfortunately, Higgs boson is too heavy for EWBG!



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## **Electroweak baryogenesis with New Physics**

# Are there new degrees of freedom that modify the scalar potential and lead to a SFOPT for successful baryogenesis?

- Prime example: MSSM with a light stop
  - Lattice calculations set limit of <155 GeV
  - Is the necessary light stop excluded?

Delphine et al. (1996), Carena et al. (1996, 1998, 2003, 2009), Espinosa et al. (1996), Huber et al. (1999), Profumo (2007), Curtin (2012), Liebler (2015) and more....

#### • General extended scalar sectors, e.g.

- 2HDM with extra bottom Yukawa coupling Modak et al. (2020)
- B-LSSM (B-L symmetric MSSM) Yang et al. (2019)
- New gauge singlets and vector-like leptons Bell et al. (2019)

#### **General difficulties:**

- Constraints from EDMs
- Higgs physics sets stringent constraints







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# **Overview of baryogenesis models**





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### High-scale baryogenesis & nn oscillations



HIBEAM/NNBAR program is a proposed two-stage experiment at the European Spallation Source (ESS) **to search for baryon number violation**.

Future sensitivity at ESS:

$$\tau_{n\overline{n}} \ge 10^{10} s$$

Naive estimate:

$$\tau_{n-\bar{n}} \approx \frac{\Lambda_{\rm NP}^5}{\Lambda_{\rm QCD}^6}$$

 $\Lambda_{\rm NP} > 10^6 {\rm GeV}$ 



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#### Free vs bound neutron-antineutron oscillations



#### **Free oscillation**





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69

#### What do we learn from nn oscillations?

- Baryon number violating
- Possible washout process (BNV)
- Possible asymmetry generation mechanism (BNV, CPV, OoE)



Observation of neutron-antineutron oscillations at  $arLambda_{
m NP}$ 

 $\frac{\Gamma_W(T, \Lambda_{NP})}{H(T)} > 1$ 

Identify scale T above which the washout rate is large enough to wipe out a previously generated asymmetry.





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#### Probing baryogenesis with nn oscillations



Fridell, JH, Hati (2021)





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Emmy Noether-

DFG R

#### Probing baryogenesis with nn oscillations



Fridell, JH, Hati (2021)





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Programm

DFG :
# Probing baryogenesis with nn oscillations



n observation of neutron-antineutron oscillations would imply a strong asymmetry washout until a scale reachable at (future) colliders!

Fridell, JH, Hati (2021)



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## **Possible topologies**





Grojean et al. (2019)





- Left-right symmetric model
- SO(10) GUT
- Post-sphaleron set-up

NOW:

- simplified model set-up considering asymmetry generation (CPV source!)
- confronting with current and future experimental results

Mohapatra, Marshak (1980) Babu, Mohapatra, Nasri (2006) Baldes, Bell, Volkas (2011) Babu, Mohapatra (2012) E. Herrmann (2014)



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# Simplified model of Topology II

$$\mathcal{L}_{II}^{\text{eff}} \supset f_{ij}^{dd} X_{dd} \bar{d}_i^c \bar{d}_j^c + \frac{f_{ij}^{ud}}{\sqrt{2}} X_{ud} (\bar{u}_i^c \bar{d}_j^c + \bar{u}_j^c \bar{d}_i^c) + \lambda \xi X_{dd} X_{ud} X_{ud} + \text{h.c.}$$



 Non-SUSY SO(10) unification requires TeV-scale X<sub>ud</sub> and GUT-scale X<sub>dd</sub> / v<sub>B-I</sub>

$$m_{X_{dd}} > m_{X_{ud}} > m_d$$

Field	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$Q = T_{3L} + Y$	$B$
$X_{dd}$	$\overline{6}$ or 3	1	$+\frac{2}{3}$	$+\frac{2}{3}$	$-\frac{2}{3}$
$X_{ud}$	$\overline{6}$ or $3$	1	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$

 $x_1$ 

Topology - II







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#### Baryogenesis

$$\frac{dY_{X_{dd}}}{dz} = -\left(D+S\right)\left(Y_{X_{dd}}-^{\mathrm{eq}}_{X_{dd}}\right)$$
$$\frac{dY_B}{dz} = -\epsilon D\left(Y_{X_{dd}}-Y^{\mathrm{eq}}_{X_{dd}}\right) - W Y_B$$





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### **Results: low-scale scenario**

$$f_{ud} = f_{dd}$$
$$m_{X_{dd}} = 3m_{X_{ud}}$$



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### Results: high-scale scenario



 $m_{X_{dd}} = 10^{14} \text{ GeV}$ 





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## Conclusions

- Matter-antimatter asymmetry is one of the most pressing open questions of (astro)particle physics
- Many models, many ideas **how** can we **identify** the one that is realised in nature?
- Some models connected to **light** new physics → **testable**!
- Some models connected to new physics at **high-scale** → **difficult to test**!
- "Falsification" of high-scale models instead:
- Observation of LNV can imply strong lepton (and baryon!) asymmetry washout
- Observation of neutron-antineutron oscillations directly constrain BNV interactions; generally imply a strong washout down to LHC scales (first generation)

#### Great future ahead to (hopefully) nail down the mechanism behind BAU!



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### Thank you for your attention!



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