Paracausal deformations of globally hyperbolic spacetimes and their application in AQFT

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Università degli Studi di Trento

17/12/2021

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Questions:

Can we relate classical field theories defined on different curved backgrounds? (PDE problem)

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Møller operators.

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- Can we relate classical field theories defined on different curved backgrounds? (PDE problem)
- What about the quantum field theories? (algebras and states) Answers:
 - Møller operators.
 - Paracausal deformations of globally hyperbolic metrics;

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Preliminaries: Lorentzian geometry and globally hyperbolic spacetimes.

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

- Preliminaries: Lorentzian geometry and globally hyperbolic spacetimes.
- Paracausal deformations of globally hyperbolic metrics.
- Convex combinations of normally hyperbolic operators and the Møller operators.

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Based on a recent paper with V.Moretti and S.Murro: Paracausal deformations of Lorentzian metric and geometric Møller isomorphisms in algebraic quantum field theory. (2021).

Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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■ Lorentzian manifold (n + 1)-dimensional manifold (M, g) with $g \in \Gamma(T^*M \otimes_s T^*M)$ and signature (-, +, ..., +);

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- Causal sets J[±](A) = A∪ points in M reachable by future/past directed smooth causal curves.

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Spacetime: smooth, connected, oriented, time-oriented *n* + 1 dimensional Lorentzian manifold.

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■ Globally hyperbolic spacetime: $g \in \mathcal{GH}(M)$ if no closed timelike curves + compact diamonds $J^+(P) \cap J^-(Q)$ for all $P, Q \in M$.

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Theorem (Bernal-Sànchez)

 $(\mathsf{M}, g) \text{ globally hyperbolic } \Longrightarrow \exists$

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 (M, g) globally hyperbolic $\implies \exists$

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- an isometry $h = -\beta^2 dt^2 + h_t$, h_t family of Riemannian metrics on the slices.

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• $\mathcal{M}(M)$ Lorentzian metrics.

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Preorder relation on $\mathcal{M}(M)$

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2 $g_{\chi} = ((1 - \chi)g^{\sharp} + \chi g'^{\sharp})^{\flat}$ and $g_{\lambda} = (1 - \lambda)g + \lambda g' \in \mathcal{M}(M)$;

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Proofs of these statments are done pointwise: exercises about quadratic forms.

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 g_{χ} will be important later!

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Hyp: $g \in \mathcal{T}(M)$ $g' \in \mathcal{GH}(M)$; $\chi, \lambda \in C^{\infty}(M, [0, 1])$

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Paracausal deformations of globally hyperbolic spacetimes and their application in AQFT

Daniele Volpe



Hyp: $g \in \mathcal{T}(M)$ $g' \in \mathcal{GH}(M)$; $\chi, \lambda \in C^{\infty}(M, [0, 1])$ **1** Cauchy hypersurfaces of g' are Cauchy for g;

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Hyp: g ∈ *T*(M) g' ∈ *GH*(M); χ, λ ∈ C[∞](M, [0, 1])
1 Cauchy hypersurfaces of g' are Cauchy for g;
2 g, g_λ, g_χ ∈ *GH*(M).
Sketch of the proof:

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Sketch of the proof: $v \in T_PM$ g'spacelike $\implies v$ is g spacelike. Therefore Σ is g' spacelike $\implies \Sigma$ is g spacelike. g timelike curves are g' timelike, so they intersect Σ once $\implies \Sigma$ is g Cauchy.

The paracausal relation I

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Definition (Paracausal relation)

$$\begin{array}{l} g,g' \in \mathcal{GH}_{\mathsf{M}} \\ g \text{ is paracausally related to } g \left(g \simeq g'\right) \text{ if there is a finite sequence} \\ g = g,g',\ldots,g_N = g' \in \mathcal{GH}_{\mathsf{M}} \text{ such that, for } k = 0,\ldots,N-1, \\ (\text{i}) \ g_k \leq g_{k+1} \text{ or } g_{k+1} \leq g_k \\ (\text{ii}) \ \text{ (a) if } g_k \leq g_{k+1}, \text{ then } V_p^{g_k+} \subset V_p^{g_{k+1}+} \text{ for all } p \in \mathsf{M}, \\ (\text{b) if } g_{k+1} \leq g_k, \text{ then } V_p^{g_{k+1}+} \subset V_p^{g_p+} \text{ for all } p \in \mathsf{M}. \end{array}$$

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The paracausal relation

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The paracausal relation



 $g_{k-1} \preceq g_k$

 $g_k \leq g_{k-1}$

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v

 $V_p^{g_k+}$

At each step the future cones of one metric are included in the future cones of the other metric!

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$$M = \mathbb{R}^{n+1}$$

$$\eta_0 = -dt \otimes dt + \sum_{i=1}^{n+1} dx^i \otimes dx^i$$

$$\eta_1 = -d\tau \otimes d\tau + \sum_{i=1}^{n+1} dy^i \otimes dy^i$$

$$\tau = x_1, \ t = y_1, \ y_k = x_k \text{ if } k > 1.$$

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The partially ordered (P.O.) sequence relating $\eta_0 \simeq \eta_1$:



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Minkowski cylinders $R \times S^1$ with opposite time orientations:

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Minkowski cylinders $R \times S^1$ with opposite time orientations:



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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT Minkowski cylinders $R \times S^1$ with opposite time orientations:



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One metric in the sequence $\notin \mathcal{GH}(M)$:

Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT Minkowski cylinders $R \times S^1$ with opposite time orientations:



One metric in the sequence $\notin \mathcal{GH}(M)$:



Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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Paracausal relation: sufficient conditions

$$g \simeq g'$$

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Paracausal relation: sufficient conditions

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$$g \simeq g'$$

$$I \quad V_P^{g^+} \cap V_P^{g'^+} \neq \emptyset;$$

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Paracausal relation: sufficient conditions

$$\begin{split} g &\simeq g' \\ & \blacksquare \ V_P^{g+} \cap V_P^{g'+} \neq \varnothing; \end{split}$$

2 \exists common Cauchy temp. function with $t^{-1}(s)$ compact;

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 $\exists \exists g$ -Cauchy temp. function with

• $t^{-1}(s)$ compact and g' spacelike;

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- $t^{-1}(s)$ compact and g' spacelike;
- dt is g' past directed.

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Paracausal relation: sufficient conditions

$$V_P^{g+} \cap V_P^{g'+} \neq \emptyset;$$

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2 \exists common Cauchy temp. function with $t^{-1}(s)$ compact;

3 ∃ g-Cauchy temp. function with

- $t^{-1}(s)$ compact and g' spacelike;
- dt is g' past directed.

Equivalent characterization

$$\begin{array}{l} g \simeq g' \iff \exists \text{ a sequence } \{g_i\} \subset \mathcal{GH}(\mathsf{M}) \text{ such that} \\ V_P^{g_i+} \cap V_P^{g_{i+1}+} \neq \varnothing \ \forall P \in \mathsf{M}. \end{array}$$

Normally hyperbolic operators

Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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■ A spacetime (M, g).

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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- A spacetime (M, g).
- A complex hemitian vector bundle E equipped with metric compatible connection ∇.

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• The space of its smooth sections $\Gamma(E)$.

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Normally hyperbolic operator:

A linear second order differential operator $N:\Gamma(E)\to\Gamma(E)$ with $\sigma_N(\xi)=-g^\sharp(\xi,\xi)\,Id_E.$

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• Solutions: $Ker_{sc}(N) := \{ \mathfrak{f} \in \Gamma_{sc}^{g}(E) \mid N\mathfrak{f} = 0 \}$

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- **Solutions**: Ker_{sc}(N) := { $\mathfrak{f} \in \Gamma_{sc}^{g}(E) \mid N\mathfrak{f} = 0$ }
- Symplectic form $\sigma_g^{\mathsf{N}} : \operatorname{Ker}_{sc}(\mathsf{N}) \times \operatorname{Ker}_{sc}(\mathsf{N}) \to \mathbb{C}$

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Easiest example of n.h.o: Klein-Gordon operator $K = \Box_g + m^2$ on the trivial bundle.

Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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1 If the metric tensor g is **globally hyperbolic**

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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If the metric tensor g is globally hyperbolic \implies the Cauchy problem for N is well-posed.

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2 Moreover the solution "propagates with finite speed".

Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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 $(1)+(2) \implies$ normally hyperbolic operators on globally hyperbolic spacetimes are **Green hyperbolic**.

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Definition (Green hyperbolic operators)

There exist advanced Green operator and retarded Green operator $G^\pm\colon\Gamma_{\textit{pc/fc}}(E)\to\Gamma(E)$

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$$(G^{\pm}\mathfrak{f}) \subset J^{\pm}(\operatorname{supp}\mathfrak{f})$$
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$$\blacksquare \ \mathsf{G}^{\pm} \circ \mathsf{N}\, \mathfrak{f} = \mathsf{N} \circ \mathsf{G}^{\pm} \mathfrak{f} = \mathfrak{f} \text{ for all } \mathfrak{f} \in \mathsf{\Gamma}_{\mathit{pc/fc}}(\mathsf{E}) \ ,$$

• supp $(G^{\pm}\mathfrak{f}) \subset J^{\pm}(\operatorname{supp}\mathfrak{f})$ for all $\mathfrak{f} \in \Gamma_{pc/fc}(\mathsf{E})$;

The kernel is characterized by the causal propagator

$$\mathsf{G} := \mathsf{G}^+|_{\mathsf{\Gamma}_c(\mathsf{E})} - \mathsf{G}^-|_{\mathsf{\Gamma}_c(\mathsf{E})} : \mathsf{\Gamma}_c(\mathsf{E}) \to \mathsf{\Gamma}(\mathsf{E})$$

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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Let us fix a differentiable manifold M and two different globally hyperbolic metrics g and g'.

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

Daniele Volpe

Let us fix a differentiable manifold M and two different globally hyperbolic metrics g and g'.

Question

What's the relation between the solution spaces of N and N', normally hyperbolic respectively w.r.t g and g'?

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Gluing the operators

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Gluing the operators

• Let $\chi \in C^{\infty}(M, [0, 1])$ with $\chi = 0$ before t_0 and $\chi = 1$ after t_1 $(t_1 > t_0)$.What about the Cauchy problem for $N_{\chi} = (1 - \chi)N + \chi N'$?

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Principal symbol:

$$\sigma_2(\mathsf{N}_{\chi},\xi) = -(1-\chi)g_0^{\sharp}(\xi,\xi)\mathsf{Id}_{\mathsf{E}} - \chi g_1^{\sharp}(\xi,\xi)\mathsf{Id}_{\mathsf{E}} = -g_{\chi}^{\sharp}(\xi,\xi)\mathsf{Id}_{\mathsf{E}} \,.$$

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

Daniele Volpe

Let us fix a differentiable manifold M and two different globally hyperbolic metrics g and g'.

Question

What's the relation between the solution spaces of N and N', normally hyperbolic respectively w.r.t g and g'?

Gluing the operators

• Let $\chi \in C^{\infty}(M, [0, 1])$ with $\chi = 0$ before t_0 and $\chi = 1$ after t_1 $(t_1 > t_0)$.What about the Cauchy problem for $N_{\chi} = (1 - \chi)N + \chi N'$?

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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Møller maps

$$\begin{array}{l} g \leq g' \implies \exists \text{ isomorphisms:} \\ \mathsf{R}_{+} = \mathsf{Id} - \mathsf{G}^{+}_{\mathsf{N}_{\chi}}(\mathsf{N}_{\chi} - \mathsf{N}) : \mathit{Ker}^{g}_{\mathit{sc}}(\mathsf{N}) \rightarrow \mathit{Ker}^{g_{\chi}}_{\mathit{sc}}(\mathsf{N}_{\chi}) \end{array}$$

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$$\mathsf{R} = \mathsf{R}_{-} \circ \mathsf{R}_{+} : \mathit{Ker}^{g}_{\mathit{sc}}(\mathsf{N}) \to \mathit{Ker}^{g'}_{\mathit{sc}}(\mathsf{N}')$$

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Symplectic forms preserved:

 $\sigma_{g'}^{\mathsf{N}'}(\mathsf{R}\Psi,\mathsf{R}\Phi)=\sigma_g^{\mathsf{N}}(\Psi,\Phi)\quad\text{for every }\Psi,\Phi\in\mathsf{Ker}_{\mathit{sc}}^g(\mathsf{N}).$

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

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Paracausally related metrics: $g' \simeq g$;

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- Paracausally related metrics: $g' \simeq g$;
- **P.O.** sequence of metrics: $g_0 := g, g_1, \ldots, g_N := g' \in \mathcal{GH}_M$;

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■ Natural n.o. operators: $N_0 := N, N_1, \dots, N_N := N' : \Gamma(E) \rightarrow \Gamma(E)$;

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General Møller operator: $R = R_0 \cdots R_{N-1}$.

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Møller operator:

 $\mathsf{R}: \Gamma(\mathsf{E}) \to \Gamma(\mathsf{E})$

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Møller operator:

$$\mathsf{R}: \Gamma(\mathsf{E}) \to \Gamma(\mathsf{E})$$

• New definition of the adjoint operator:

$$\int_{\mathsf{M}} \langle \mathfrak{h}(x) \,|\, (\mathsf{T}\mathfrak{f})(x) > \operatorname{vol}_{g'}(x) = \int_{\mathsf{M}} \langle \left(\mathsf{T}^{\dagger_{gg'}}\mathfrak{h}\right)(x) \,|\, \mathfrak{f}(x) > \operatorname{vol}_{g}(x)$$

 $\mathfrak{f} \in \mathsf{Dom}(\mathsf{T}), \mathfrak{h} \in \Gamma_c(\mathsf{E})$

Adjoint Møller operator:

$$\mathsf{R} \to \mathsf{R}^{\dagger_{gg'}} : \Gamma_c(\mathsf{E}) \to \Gamma_c(\mathsf{E})$$

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Causal propagators

The Møller operator intertwines the causal propagators.

$$RG_N R^{\dagger_{gg'}} = G_{N'}$$

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Quantization: the Møller *-isomorphism

deformations of globally hyperbolic spacetimes and their applications in AQFT

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CCR *-algebras: $\mathcal{A} = \frac{\bigoplus_{n} Ker_{sc}^{g}(N)^{\otimes n}}{\Psi_{\psi} \otimes \Phi_{\phi} - \Phi_{\phi} \otimes \Psi_{\psi} - \sigma_{\nu}^{N}(\psi, \phi) \mathsf{ld}};$

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Quantization: the Møller *-isomorphism

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$$\mathcal{R}: \mathcal{A} \to \mathcal{A}';$$

States: $\omega : \mathcal{A} \to \mathbb{C}$ such that

$$\omega(\mathsf{Id}) = 1$$
 $\omega(a^*a) \ge 0;$

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Paracausal deformations of globally hyperbolic spacetimes and their application: in AQFT

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$$\omega' = \omega \circ \mathcal{R};$$

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- The singularity structure of the states is preserved!
- The isomorphism preserves Hadamard states

$$WF(\omega_2) = \{(x, k_x; y, -k_y) \in T^*\mathsf{M}^2 \setminus \{0\} | (x, k_x) \sim_{\parallel} (y, k_y), k_x \succ 0\}.$$

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Paracausal deformations of globally hyperbolic spacetimes and their applications in AQFT

Daniele Volpe

Conclusions

Free classical and quantum field theories on curved backgrounds are structurally comparable when the background metrics are paracausally related.

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For future research we plan to:

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For future research we plan to:

 study the paracausal relation of globally hyperbolic metrics in more detail;

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Thanks for the attention!