

$\approx 1989$

$$P(x,y) = - \sum_a \Psi_a(x) \overline{\Psi_a(y)}$$

kernel of the fermionic projector  
observations

(1) contains all information on fermionic wave functions

$$P(x,y) = P^{\text{sea}}(x,y) - \sum_a \Psi_a(x) \overline{\Psi_a(y)}$$

$\uparrow$  particle states

$$+ \sum_b \Phi_b(y) \overline{\Phi_b(x)}$$

$\underbrace{\phantom{\sum_b}}$  antiparticle states

(2) contains all information on bosonic fields  $A$

$$P^{\text{sea}}(x,y) = \int \frac{d^4 k}{(2\pi)^4} (k + m) \delta(k^2 - m^2) \Theta(-k) e^{-ik(x-y)}$$

$$(i\not\cancel{D} - m) P^{\text{sea}}(x,y) = 0$$

$$(i\not\cancel{D} + \not{A} - m) \tilde{P}^{\text{sea}}(x,y) = 0$$

$P^{\text{sea}}(x,y)$  determines  $A$

(3) contains the causal structure

$P(x,y)$  is singular if  $(x-y)^2 = 0$

disregard:

Dirac eqn

Maxwell & Einstein eqns

Minkowski space

instead formulate eqns directly with  $P(x,y)$

Let  $M$  be a discrete set

$$x \in M : (S_x, \langle \cdot | \cdot \rangle_x) \text{ spin space}$$

$$P(x, y) : S_y \rightarrow S_x$$

$$\Downarrow \sum_a |\Psi_a(x)\rangle\langle\Psi_{a(y)}|$$

$$\Psi_a : M \rightarrow SM := \bigcup_{x \in M} S_x$$
$$x \mapsto \Psi_a(x) \in S_x$$

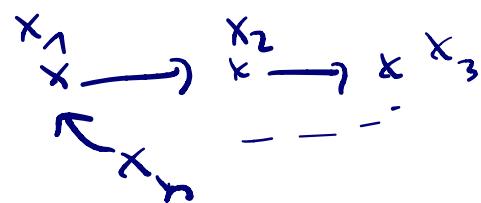
principle of the fermionic projector ( $\approx 1950$ )

guiding principles

- variational approach, action principle
- the Dirac sea vacuum should be a critical point
- Maxwell + Einstein eqns should come out  
in a specific limiting case
- Dirac eqn should hold in a specific limiting case

n-point action

$$x_1, \dots, x_n \in M$$



closed chain  $P(x_1, x_n) P(x_n, x_{n-1}) \dots P(x_3, x_2) P(x_2, x_1) : S_{x_n} \hookrightarrow$

$$A_{x_1, \dots, x_n} :=$$

$\mathcal{L}(A_{x_1, \dots, x_n})$  Lagrangian

$$S = \sum_{x_1, \dots, x_n \in M} \mathcal{L}(A_{x_1, \dots, x_n})$$

gauge phases

$$P^{\text{vac}}(x,y) e^{-i \int_x^y A_j(y-x)^j} + \dots \\ = P^{\text{vac}}(x,y) e^{i(N(y)-N(x))}$$

Lagrangian is gauge invariant.

$h=1$ :  $\mathcal{L}(P(x,x))$  too simple

$h > 2$ : too rigid



Remark 6.2.5 in PFP-book 2006

two-point actions

$$\mathcal{L}(A_{xy}), \quad A_{xy} = P(x,y) P(y,x)$$

$$S = \sum_{xy \in M} \mathcal{L}(A_{xy}) \quad \text{closed chain}$$

first attempt: choose  $\mathcal{L}$  as a

→ polynomial in traces of powers of  $A_{xy}$

$$\text{e.g. } \text{Tr}(A_{xy}^2) - c(\text{Tr}(A_{xy}))^2$$

1996      }  
1997      } preprints in German

basic difficulty: eigenvalues of  $A_{xy}$   
involve phase factors

~1999: take absolute values of  $\lambda_i^{xy}$

- two advantages : - phases of  $\lambda_i^{xy}$  drop out  
 -  $\mathcal{S} \geq 0$  ; minimization possible

$$\mathcal{L}(x,y) = \sum_{i,j} (|\lambda_i^{xy}|^p - |\lambda_j^{xy}|^p)^2, \quad p \geq 1$$

2002-03 : state stability gives  $p=1$   
 details in book of 2006.

since then : constraint  
 generalise to  $\mathcal{S}$  measure an  
 local correlation operator

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