

smooth setting,  $\rho$  a critical measure, ...

Noether-like theorem:

$$\Phi_\tau : M \rightarrow \mathcal{F}, \quad \tau \in (-\delta, \delta)$$

$$\text{symmetry} \quad \mathcal{L}(\Phi_\tau(x), y) = \mathcal{L}(x, \Phi_{-\tau}(y))$$

let  $\Omega \subset M$  be compact

$$\frac{d}{d\tau} \int_\Omega d\rho(x) \int_{M \setminus \Omega} d\rho(y) (\mathcal{L}(\Phi_\tau(x), y) - \mathcal{L}(x, \Phi_{-\tau}(y))) \Big|_{\tau=0} = 0$$

$$v(x) = \frac{d}{d\tau} \Phi_\tau(x) \Big|_{\tau=0} \in \Gamma(M, T\mathcal{F})$$

$$\text{symmetry} \quad (D_{1,\underline{v}} + D_{2,\underline{v}}) \mathcal{L}(x,y) = 0$$

$$\int_\Omega d\rho(x) \int_{M \setminus \Omega} d\rho(y) (D_{1,\underline{v}} - D_{2,\underline{v}}) \mathcal{L}(x,y) = 0,$$

corresponding to  $J_R^2$   
(not  $J_{R+\tau}^2$ )!

general class of conservation laws for  $k=1$ .

$\tilde{\mathcal{S}}_\rho = (F_\rho)_* (f_\rho \rho)$  family of critical measures

$$\int_\Omega d\rho(x) \int_{M \setminus \Omega} d\rho(y) (\partial_{1,\rho} - \partial_{2,\rho}) (f_\rho(x) \mathcal{L}(F_\rho(x), F_\rho(y)) f_\rho(y)) \Big|_{\rho=0}$$

$$= \int_\Omega \partial_\rho f_\rho(x) d\rho(x) \Big|_{\rho=0}$$

$\underline{v} = \frac{d}{d\rho} (f_\rho, F_\rho) \Big|_{\rho=0}$  is a solution of the  
linearized field eqns

$$\int_\Omega d\rho(x) \int_{M \setminus \Omega} d\rho(y) (\nabla_{1,\underline{v}} - \nabla_{2,\underline{v}}) \mathcal{L}(x,y) \quad \langle \underline{u}, \Delta \underline{v} \rangle \Big|_{\rho=0}$$

$$= \int_\Omega \nabla_{\underline{v}} \rho \quad \underline{\text{conserved one-form}}$$

alternative proof:

anti-symmetric

$$\int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) (\underline{D}_{1,\underline{v}} - \underline{D}_{2,\underline{v}}) \mathcal{L}(x,y)$$

$$= \int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) (\underline{D}_{1,\underline{v}} - \underline{D}_{2,\underline{v}}) \mathcal{L}(x,y)$$

$$+ \int_{\Omega} dg(x) \int_{\Omega} dg(y) (\underline{D}_{1,\underline{v}} - \underline{D}_{2,\underline{v}}) \mathcal{L}(x,y) \} = 0$$

$$= \int_{\Omega} dg(x) \int_M dg(y) (\underline{D}_{1,\underline{v}} - \underline{D}_{2,\underline{v}}) \mathcal{L}(x,y)$$

$$= \int_{\Omega} \underline{D}_{\underline{v}} \circ dg(x)$$

↑  
linearized field eqns

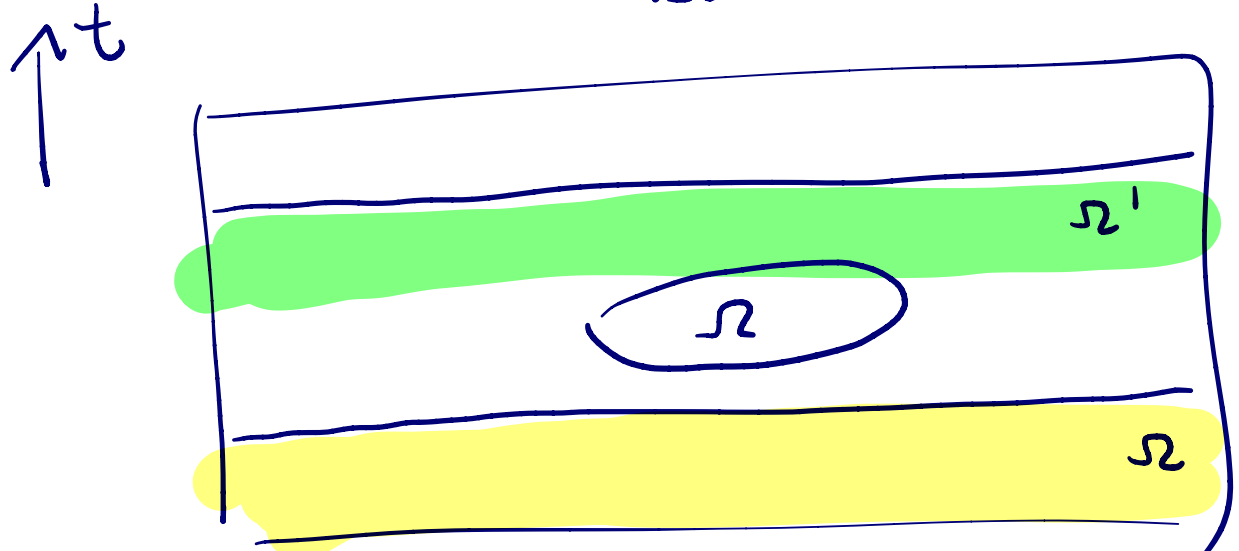
back to causal fermion systems,  $\dim \mathcal{H} < \infty$

$$e(x) = i [A, x], \quad A \text{ symmetric operator on } \mathcal{H}$$

$$\underline{v} = (0, e) \quad \text{commutator jet}$$

is a solution of the linearized field eqns without scalar component

$$\Rightarrow \int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) (\underline{D}_{1,\underline{v}} - \underline{D}_{2,\underline{v}}) \mathcal{L}(x,y) = 0$$



$$\int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) (D_{1,e} - D_{2,e}) \mathcal{L}(x,y)$$

$$= \int_{\Omega'} dg(x) \int_{M \setminus \Omega'} dg(y) (D_{1,e} - D_{2,e}) \mathcal{L}(x,y)$$

Consider the case  $A = |u\rangle\langle u|$ ,  $u \in \mathcal{H}$

$$\langle u|u \rangle_{\mathcal{H}}^{\Omega} := \int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) (D_{1,e(A)} - D_{2,e(A)}) \mathcal{L}(x,y)$$

Polarization gives

Commutator inner product

$$\langle u|v \rangle_{\mathcal{H}}^{\Omega} : \mathcal{H} \times \mathcal{H} \rightarrow \mathbb{C} \quad \text{sesquilinear form}$$

can be expressed in terms of physical wave functions  $\psi^u$  and  $\psi^v$

generalized to wave functions  $\psi, \phi \in C^0(M, SM)$

Back to causal variational principles

assume that  $M$  has a smooth manifold structure

$$\underline{v} = (\text{div } v, v) \quad \text{min solution, } v \in \Gamma(M, TM)$$

conserved one-form

$$\int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) (\nabla_{1,\underline{v}} - \nabla_{2,\underline{v}}) \mathcal{L}(x,y)$$

$$\stackrel{\text{Gauss}}{=} \int_{\partial\Omega} d\mu(v, x) \int_{M \setminus \Omega} dg(y) \mathcal{L}(x,y)$$

$$- \int_{\Omega} dg(x) \left( - \int_{\partial\Omega} d\mu(v, y) \mathcal{L}(x,y) \right) \approx$$

$$= \int_{\partial\Omega} d\mu(v, x) \left( \int_M \mathcal{L}(x,y) dg(y) \right)$$

$$= \rho \cdot \int_{\partial \Omega} d\mu(v, x)$$

explain  $d\mu$

$$dg = h(x) d^k x$$

volume form

$$\omega = h(x) \epsilon_{i_1 \dots i_k} dx^{i_1} \dots dx^{i_k} \quad k\text{-form}$$

$$\omega \lrcorner v = h(x) \epsilon_{i_1 \dots i_k} v^{i_1} dx^{i_2} \dots dx^{i_k} \quad (k-1)\text{-form}$$

$d\mu(v, x)$  is the corresponding volume measure