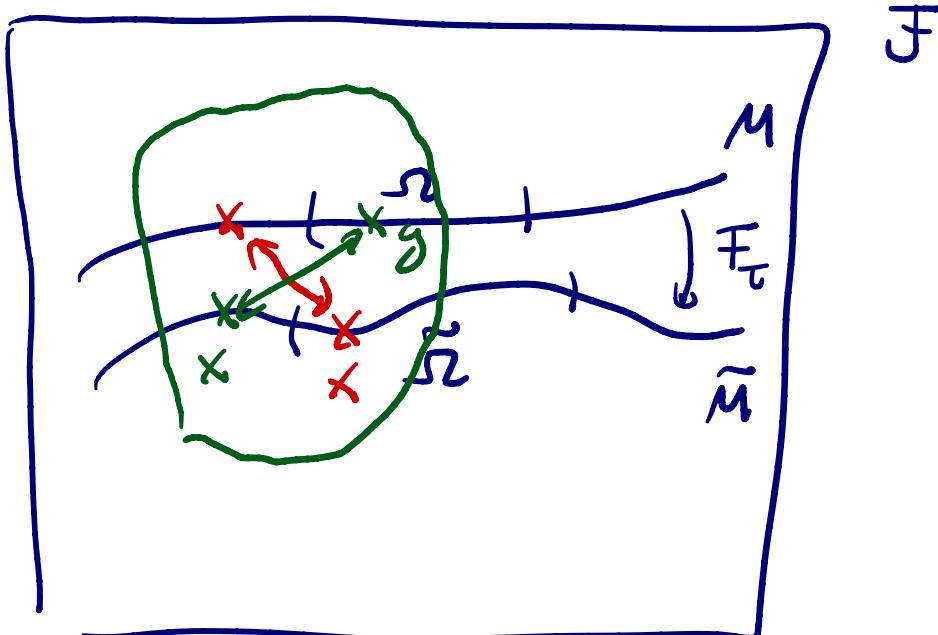


Let $\left\{ \begin{array}{l} g \text{ (vacuum)} \\ \tilde{g} \text{ (interacting system)} \end{array} \right\}$ be measures on \mathcal{F}

Let $\Omega \subset M := supp g$ be compact
 $\tilde{\Omega} \subset \tilde{M} := supp \tilde{g}$

nonlinear surface layer integral

$$\mu^{\tilde{\Omega}, \Omega} (\tilde{g}, g) := \underbrace{\left(\int_{\tilde{\Omega}} d\tilde{g}(x) \int_{M \setminus \Omega} dg(y) - \int_{\tilde{M} \setminus \tilde{\Omega}} d\tilde{g}(x) \int_{\Omega} dg(y) \right)}_{\mathcal{F}} \mathcal{L}(x, y)$$



assume that

$$\tilde{g} = (F_\tau)_* (f_\tau g), \quad F_\tau : M \rightarrow \tilde{M}$$

$$\tilde{\Omega} = F_\tau(\Omega) \quad \int_{\tilde{\Omega}} d\tilde{g}(x) g(x) = \int_{\Omega} dg(x) f_\tau(x) g(F_\tau(x))$$

$$\Rightarrow \mu^{\tilde{\Omega}, \Omega} (\tilde{g}, g) = \left(\int_{\Omega} dg(x) \int_{M \setminus \Omega} dg(y) - \int_{M \setminus \Omega} dg(x) \int_{\Omega} dg(y) \right) f_\tau(x) \mathcal{L}(F_\tau(x), y)$$

linear in τ :

conserved one-form

$$\mu_g^\Omega(\underline{v})$$

quadratic in τ :

surface layer inner product

$$(\underline{v}, \underline{v})_g + \dots$$

$$\underline{v} = \frac{d}{d\tau} (f_\tau, F_\tau) |_{\tau=0}$$

conservation law

identify M and \tilde{M}

$$F: M \rightarrow \tilde{M}$$

$$\underline{\Phi} := F^{-1}, \tilde{M} \rightarrow M$$

$$\text{choose } \tilde{\Omega} = F(\Omega)$$

goal: choose F such that $\mu^{\tilde{\Omega}, \Omega} (\tilde{f}, g) = 0 \quad \forall \Omega \subset M$
compact

$\underline{\Phi}_* \tilde{f}$ measure on M

$$\begin{aligned} \Rightarrow \mu^{\tilde{\Omega}, \Omega} (\tilde{f}, g) &= \int_{\Omega} d(\underline{\Phi}_* \tilde{f}) \otimes \int_{M \setminus \Omega} dg(y) \mathcal{L}(F(x), y) \\ &\quad - \int_{\Omega} dg(x) \int_{M \setminus \Omega} d(\underline{\Phi}_* \tilde{f}) \mathcal{L}(x, F(y)) \end{aligned}$$

Using a symmetry argument,

$$\begin{aligned} &= \int_{\Omega} d(\underline{\Phi}_* \tilde{f}) \otimes \int_M dg(y) \mathcal{L}(F(x), y) \\ &\quad - \int_{\Omega} dg(x) \int_M d(\underline{\Phi}_* \tilde{f}) \mathcal{L}(x, F(y)) \\ &= \int_{\Omega} d(\underline{\Phi}_* \tilde{f}) \int_M dg(y) \mathcal{L}(F(x), y) = (\underline{\Phi}_* \tilde{f})(\Omega) \\ &\quad - \int_{\Omega} dg(x) \int_{\tilde{M}} d\tilde{f}(y) \mathcal{L}(x, y) = \omega(\Omega) \end{aligned}$$

Introduce

$$d\omega(x) = \left(\int_{\tilde{M}} \mathcal{L}(x, y) d\tilde{f}(y) \right) dg(x) \quad \text{measure on } M$$

$$d\tilde{\omega}(x) = \left(\int_M \mathcal{L}(x, y) dg(y) \right) d\tilde{f}(x) \quad \text{measure on } \tilde{M}$$

correlation measures

$$= ((\underline{\Phi}_* \tilde{\omega}) - \omega)(\Omega) = 0 \quad \text{if } \Omega \in M \text{ compact}$$

Thus we need to arrange that

$$\underline{\Phi}_* \tilde{\omega} = \omega$$

General problem:

given ω on M
 $\tilde{\omega}$ on \tilde{M}

Is there $\underline{\Phi}: \tilde{M} \rightarrow M$ s.t. $\underline{\Phi}_* \tilde{\omega} = \omega$?

Answer: Yes, under general assumptions

Greene and Shishamana rather abstract

Moser's theorem infinitesimally,

then integrate

Causal fermion systems

$$(\mathcal{H}, F, g), \quad (\tilde{\mathcal{H}}, \tilde{F}, \tilde{g})$$

identify \mathcal{H} with $\tilde{\mathcal{H}}$

choose $V: \mathcal{H} \rightarrow \tilde{\mathcal{H}}$ unitary

↑ not canonical

freedom: $V \rightarrow VU$ with $U \in U(\mathcal{H})$

integrate over U

partition function

$$Z^{\text{sr}}(\beta, \tilde{f}) = \int_{\mathcal{U} \subset U(\mathcal{H})} e^{\beta \tilde{f} \circ \tilde{\omega}_{\mathcal{U}, \mathcal{U}}} (\tilde{f}, U_g) \frac{dU}{\text{haar measure}}$$

gives rise to
quantum state

